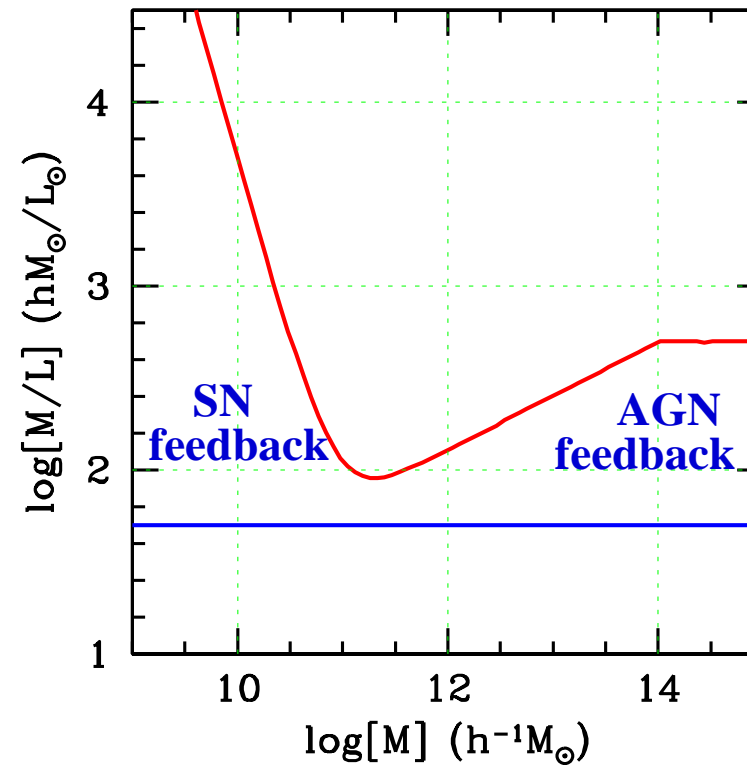
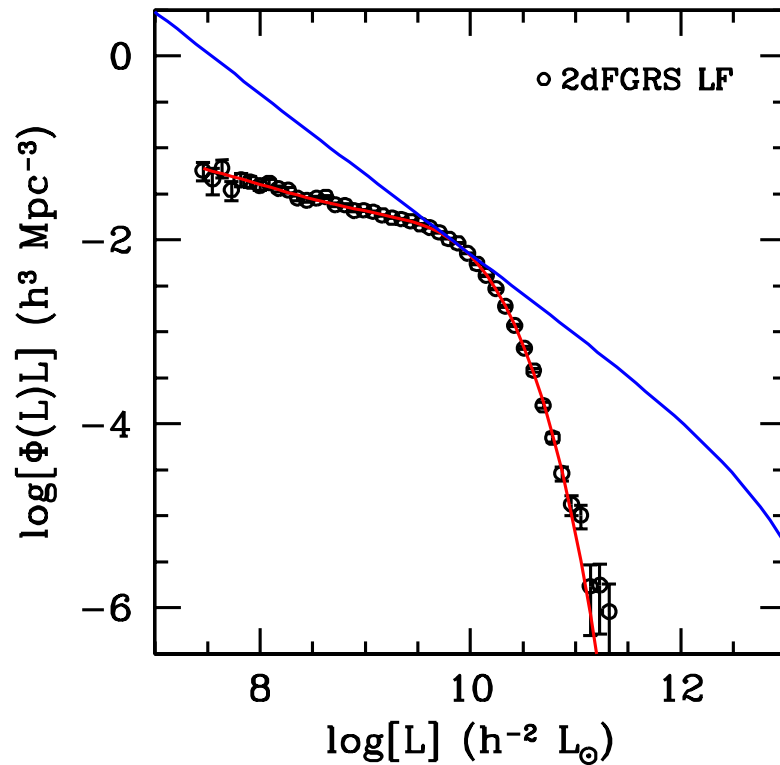


Killing Dwarfs with Hot Pancakes



Frank C. van den Bosch (MPIA)
with Houjun Mo, Xiaohu Yang & Neal Katz

The Paradigm...



- The **halo mass function** is much steeper than **luminosity function**
- Galaxy Formation has to become **inefficient** in low mass haloes
- In the **standard paradigm** this is interpreted as due to **SN feedback**
- Since **binding energy** is lower in less massive haloes, and each SN produces a constant energy of $\sim 10^{51}$ erg, **SN feedback** is indeed expected to effect less massive haloes more strongly.

and why it is wrong

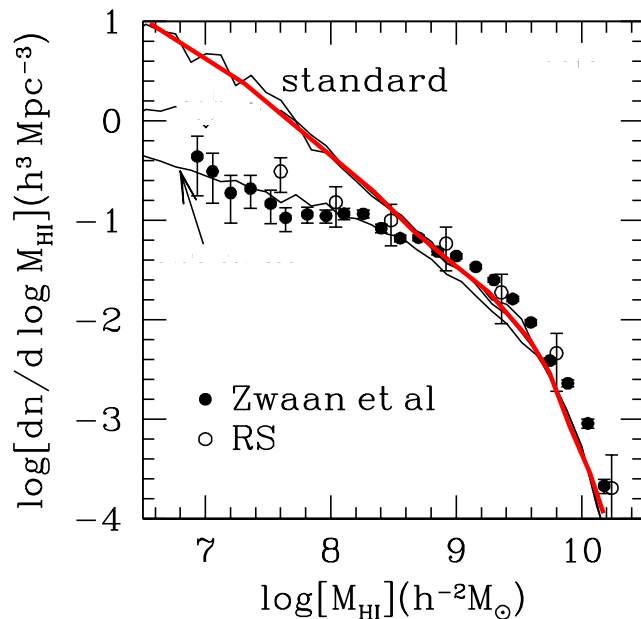
- To keep the gas **hot**, you need to **continuously** inject it with energy
As soon as energy ejection comes to halt, gas will **cool**.
- In order to have **SN**, you need to form stars
- In order to form stars, you need high column densities of HI

$$\Sigma_{\text{gas}} > \Sigma_{\text{crit}} = \frac{\sigma_{\text{gas}} \kappa}{\pi G Q_{\text{crit}}}$$

(Kennicutt 1989)

One expects that (disk) galaxies have a cold gas mass equal to

$$M_{\text{gas}} \simeq 2\pi \int \Sigma_{\text{crit}}(R) R dR$$



Starting from **halo mass function**, and assuming each halo has cold gas mass equal to M_{gas} , severely overpredicts the **HI mass function!!!**

SN feedback can be tuned to fit galaxy **luminosity function**, but it can not **simultaneously** match the **HI mass function**

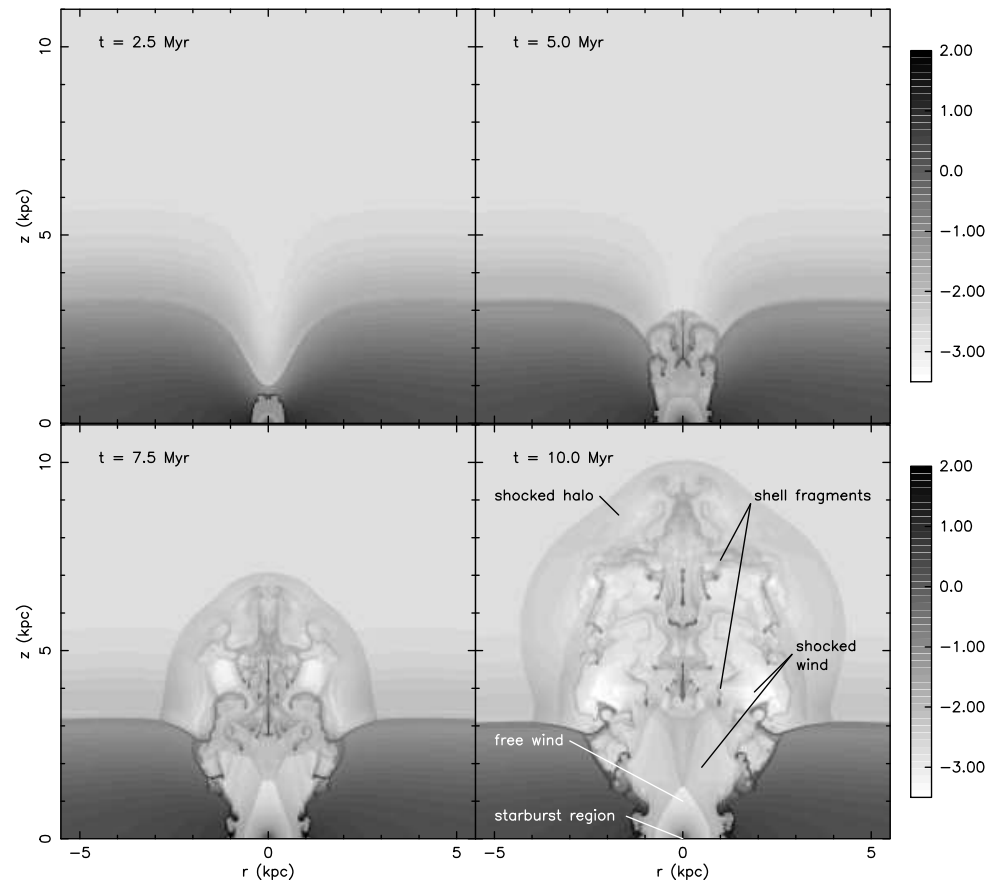
(Data: RS=Rosenberg & Schneider 2002)

and wrong

In low mass haloes, **SN feedback** is required to blow away (or keep hot) more than **90** percent of all baryons

Detailed, hydrodynamical simulations, however, show that **mass ejection efficiency** is very low, contrary to the ejection efficiency of **metals** and **energy**

(Mac-Low & Ferrara 1999; Strickland & Stevens 2000)



Mass loading efficiency of **galactic winds** is low because of onset of **Rayleigh-Taylor** and **Kelvin-Helmholtz** instabilities

Wind “bursts open”, hot (metal-rich) gas escapes but momentum of cold “loaded” gas stalls.

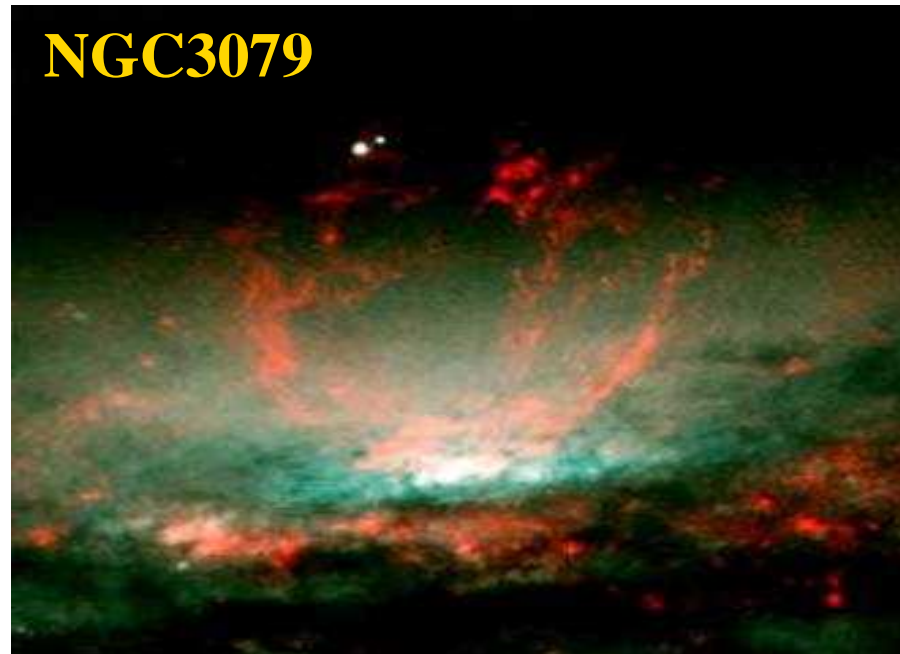
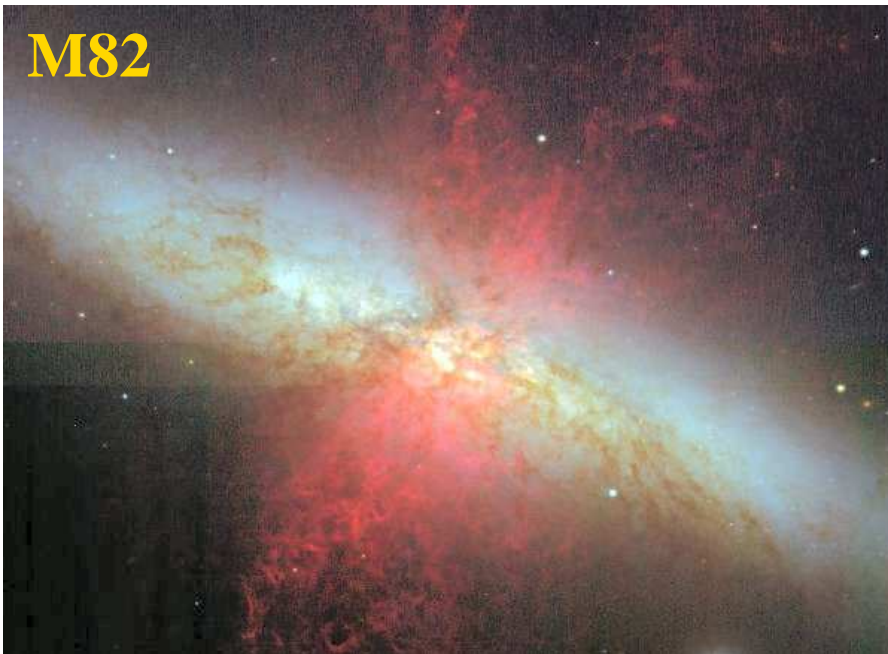
(From Strickland & Stevens 2000)

and wrong

Even in the **most extreme** cases (**starbursting dwarfs**), the observed **mass outflow rate** is about twice the **star formation rate**

(Martin 1999; Heckman et al. 2000)

One can **at most** blow out ~ 66 percent of all baryons!



Furthermore, Semi-Analytical Models require **efficiencies** with which SN energy is **thermalized** that are unrealistic or even unphysical

(e.g., Benson et al. 2003)

In case you didn't get it

SN feedback is NOT the main physical mechanism
that explains the faint-end slope of the LF

In case you didn't get it

SN feedback is NOT the main physical mechanism that explains the faint-end slope of the LF

QUESTION: but then what is the physical mechanism?

In case you didn't get it

SN feedback is NOT the main physical mechanism that explains the faint-end slope of the LF

QUESTION: but then what is the physical mechanism?

ANSWER:



The Cosmic Baryon Budget

The **WMAP** measurements of the **CMB** have constrained the baryonic and total matter densities to be

$$\Omega_b h^2 = 0.024 \pm 0.001 \quad \Omega_m h^2 = 0.14 \pm 0.02$$

(Spergel et al. 2003)

This implies a **Universal baryon fraction** of $f_{\text{bar}} = \Omega_b / \Omega_m = 0.17$

On large scales, we expect the baryons and dark matter to be well **mixed**

Within individual haloes, **shocks** and **cooling** are expected to have partially **segregated** the baryons from the dark matter.

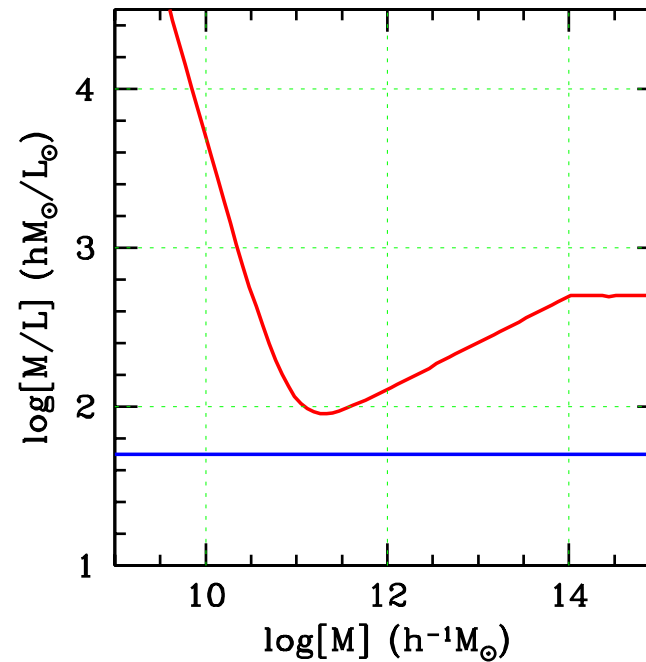
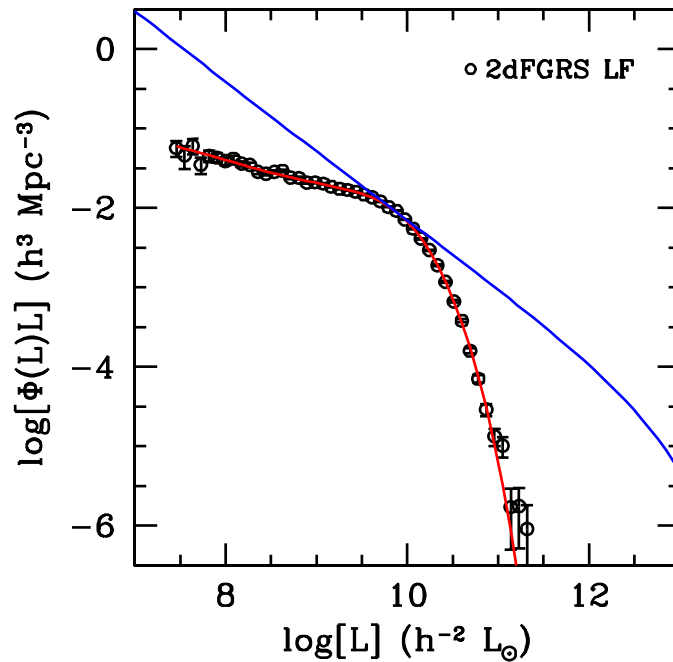
However, in the absence of **feedback** (galactic winds) one still expects that each halo has a baryonic mass fraction of $f_{\text{bar}} \simeq 0.17$

Baryons come in the form of:

- stars plus their remnants
- cold gas; both atomic & molecular
- hot gas (plasma)

How are the baryons distributed among these different components?

The stellar mass fraction



With an average **total** mass-to-light ratio of $M/L_B \sim 200$ and an average **stellar** mass-to-light ratio of $M_*/L_B \sim 5$ we obtain

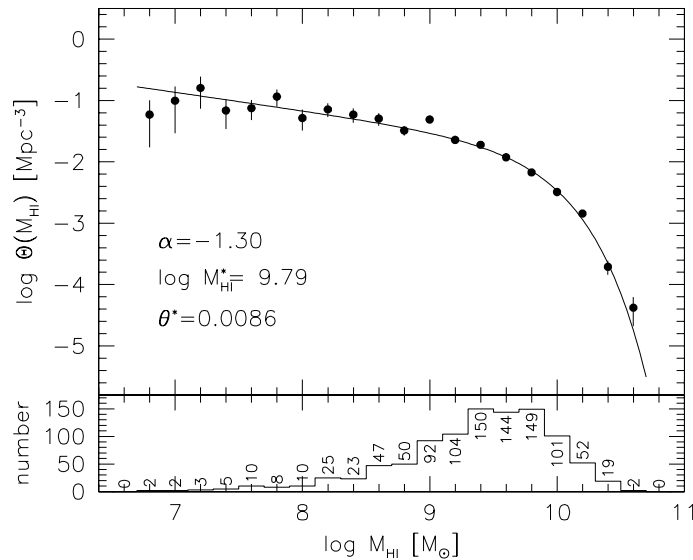
$$\left\langle \frac{M_*}{M} \right\rangle = \left\langle \frac{M_*}{L_B} \right\rangle \cdot \left\langle \frac{L_B}{M} \right\rangle \simeq \frac{5}{200} = 0.025$$

Only ~ 15 percent of all baryons are in stars

NOTE: a similar result is obtained when integrating the LFs of **disks** and **spheroids** and multiplying by their respective stellar mass-to-light ratios

(Fukugita, Hogan & Peebles 1998)

Cold & Hot Gas



The matter density in **neutral Hydrogen** can be obtained from integrating the HI mass function.

Multiply by **1.32** to account for **Helium**

Multiply by **1.8** to account for **Molecules**

(Data from Zwaan et al. 2003)

Only ~ 3 percent of the baryons are in cold gas

The matter density in **hot gas** can be obtained from **X-ray** measurements

Using a constant, measured ratio between the mass of X-ray emitting gas and the gravitational mass of clusters, and integrating over the cluster mass function, yields matter density contributed by hot cluster gas.

Adding contribution from groups, using a similar method, finally yields

About ~ 37 percent of the baryons are in hot, X-ray emitting gas

(Fukigita, Hogan & Peebles 1998)

Where are the baryons?

In summary, observationally we have accounted for the following baryons:

| | |
|----------|-----|
| Stars | 15% |
| Cold gas | 3% |
| Hot gas | 37% |
| TOTAL | 55% |

About 45% of expected baryons have not been accounted for observationally

Question: Where are the missing baryons?

- Stellar Remnants & Machos
- Primordial Black Holes
- Warm/Hot Intergalactic Medium (WHIM)

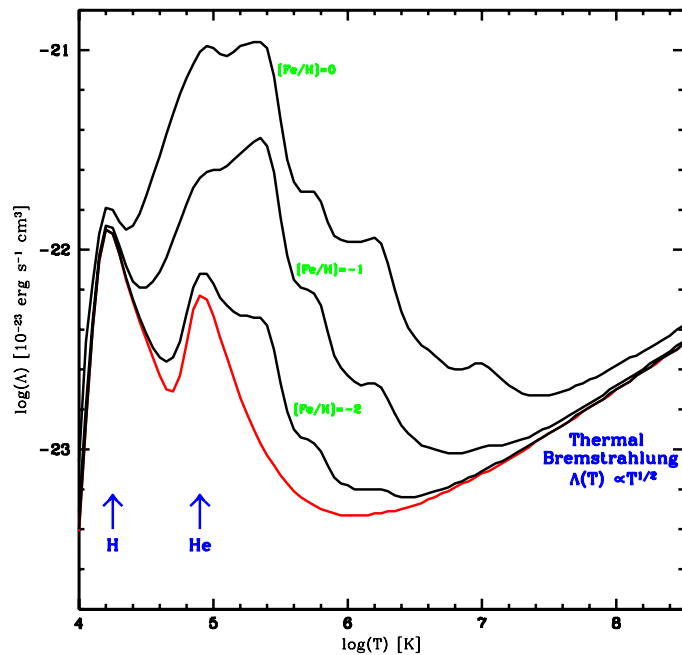
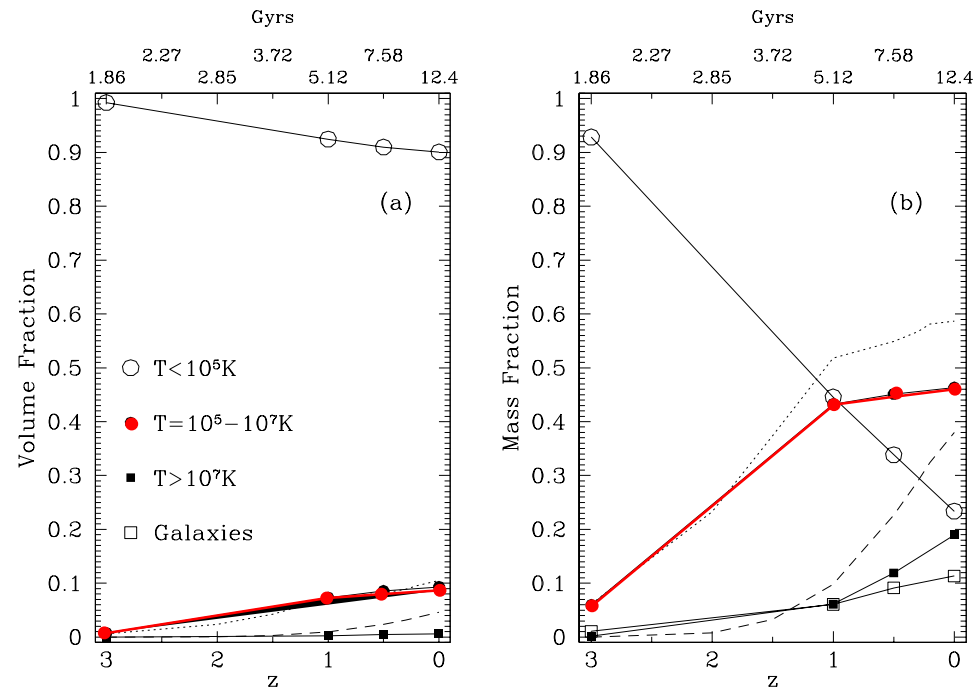
Although the former two options are serious possibilities, we focus on the case of a significant **WHIM**

The Warm/Hot Intergalactic Medium

Cosmological hydrodynamical simulations predict that about **45%** of the baryons are in **WHIM** at $z = 0$

This gas is located in **pancakes & filaments** and heated due to **shocks**

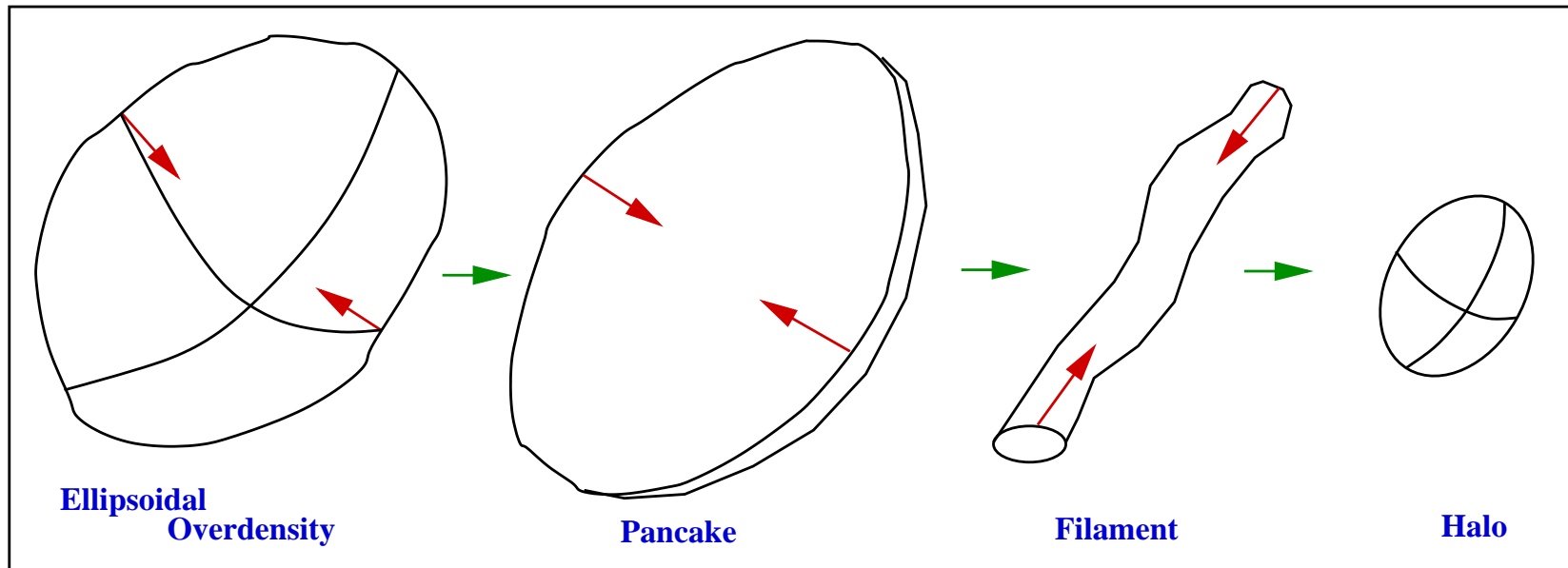
(From Cen & Ostriker 1999)



As long as the hot gas has a temperature $10^5 \text{K} \lesssim T \lesssim 10^7 \text{K}$ and is of sufficiently low density and metallicity, it will not be visible in **X-rays**

(Cooling Functions from Sutherland & Dopita 1993)

A Recipe for Hot Pancakes



- Structures form by **ellipsoidal collapse**
- Dark Matter is **collisionless** → orbit crossing & virialization
- Baryonic Matter is **collisional** → shocking & thermalization

Pancakes & Filaments are expected to be hot

Unless the **cooling time** in the pancakes is too short, dark matter haloes that form within pancakes, form in a **preheated medium**.

Pancake Physics

The **internal energy density** of a baryonic gas is:

$$U_{\text{int}} = \frac{3}{2} \frac{\rho_{\text{gas}} k_B}{\mu} T$$

The **kinetic energy density** of system that collapses with velocity V_{coll} is:

$$U_{\text{kin}} = \frac{1}{2} \rho_{\text{gas}} V_{\text{coll}}^2$$

If we consider the **strong shock jump condition** ($V_{\text{after}} = V_{\text{coll}}/4$), and assume that kinetic energy is converted into internal energy, then

$$T = \frac{3\mu V_{\text{coll}}^2}{16k_B} \propto M_p^{2/3}$$

with M_p the pancake mass. The **cooling time** of this shock-heated gas is:

$$\tau_{\text{cool}} = \frac{U}{|dU/dt|} \simeq \frac{3}{2} \mu \frac{k_B T}{\rho_{\text{gas}} \Lambda(T, Z)}$$

The **hot pancake** will stay **hot** as long as the cooling time, τ_{cool} , is longer than the Hubble time, $\tau_{\text{Hubble}} \propto H^{-1}(z)$

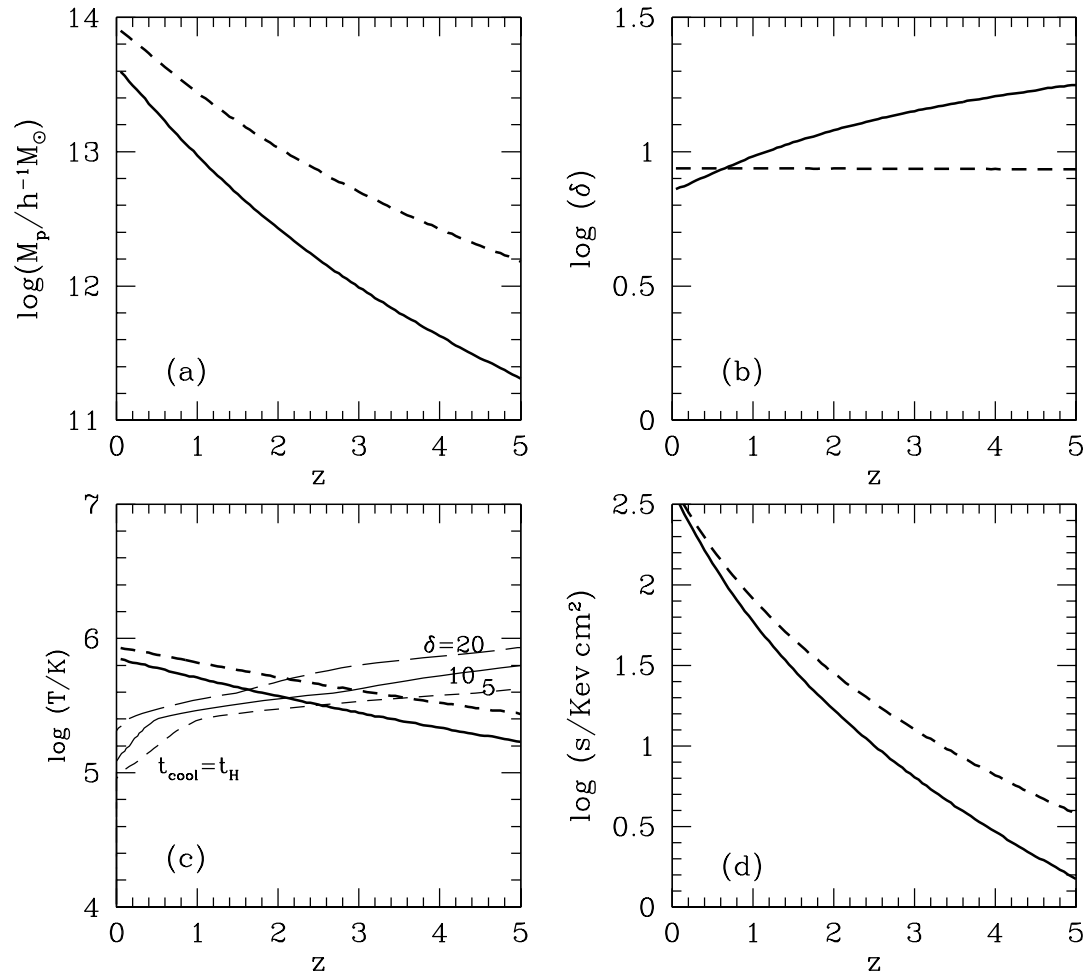
For the **density** of the gas in the pancakes we can write

$$\rho_{\text{gas}} = \delta_p \Omega_b \frac{3H_0^2}{8\pi G} (1+z)^3$$

with $\delta_p \simeq 10$ the overdensity in the pancake.

More Pancake Physics

What remains is to compute the pancake mass M_p . From **ellipsoidal collapse model** one can compute mass of pancakes that form around present-day low mass haloes, as function of redshift z



Hot Pancakes, with $M_p \gtrsim 5 \times 10^{12} h^{-1} M_\odot$, exist for $z \lesssim 2$

The Impact on Galaxy Formation

The **specific entropy** of the shock heated gas in the pancakes is

$$s = \frac{T}{n_e^{2/3}} = 17 \text{KeVcm}^2 \left(\frac{\Omega_b h^2}{0.024} \right)^{-2/3} \left(\frac{T}{10^{5.5} \text{K}} \right) \left(\frac{1+\delta}{10} \right)^{-2/3} \left(\frac{1+z}{3} \right)^{-2}$$

Because of its **high entropy**, the gas is not bound to low mass haloes.

Approximately, one finds that

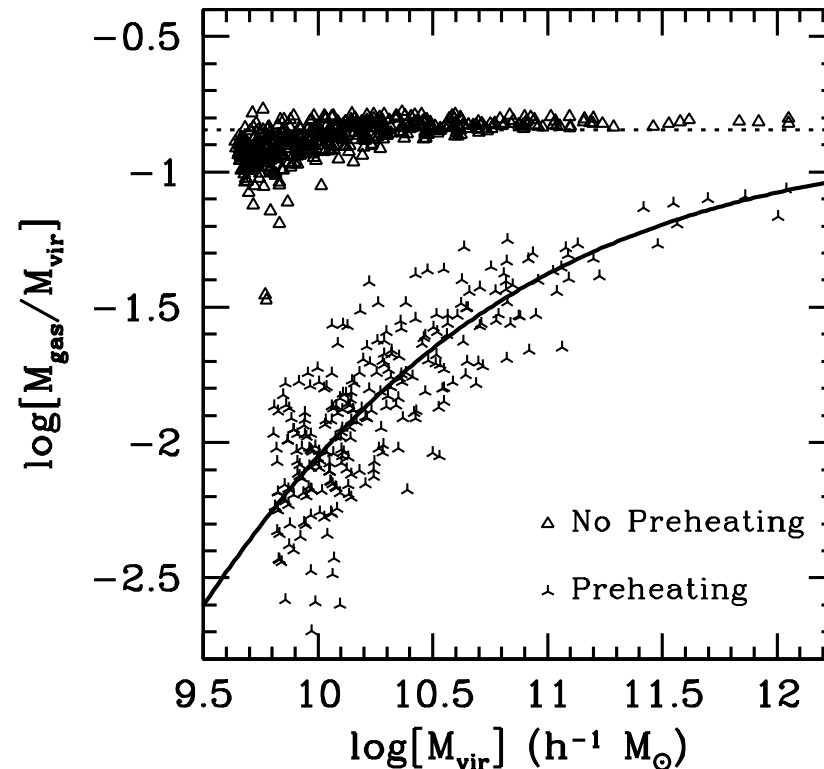
$$\frac{M_{\text{bar}}}{M_{\text{vir}}} \simeq 0.17 \left[1 + \left(\frac{M}{5 \times 10^{11} h^{-1} M_{\odot}} \right)^{-1} \right]$$

Halo formation in **preheated medium** results in a strong reduction of **baryonic mass fraction** that can cool.

In particular, galaxy formation in low mass haloes is **severely inhibited**:

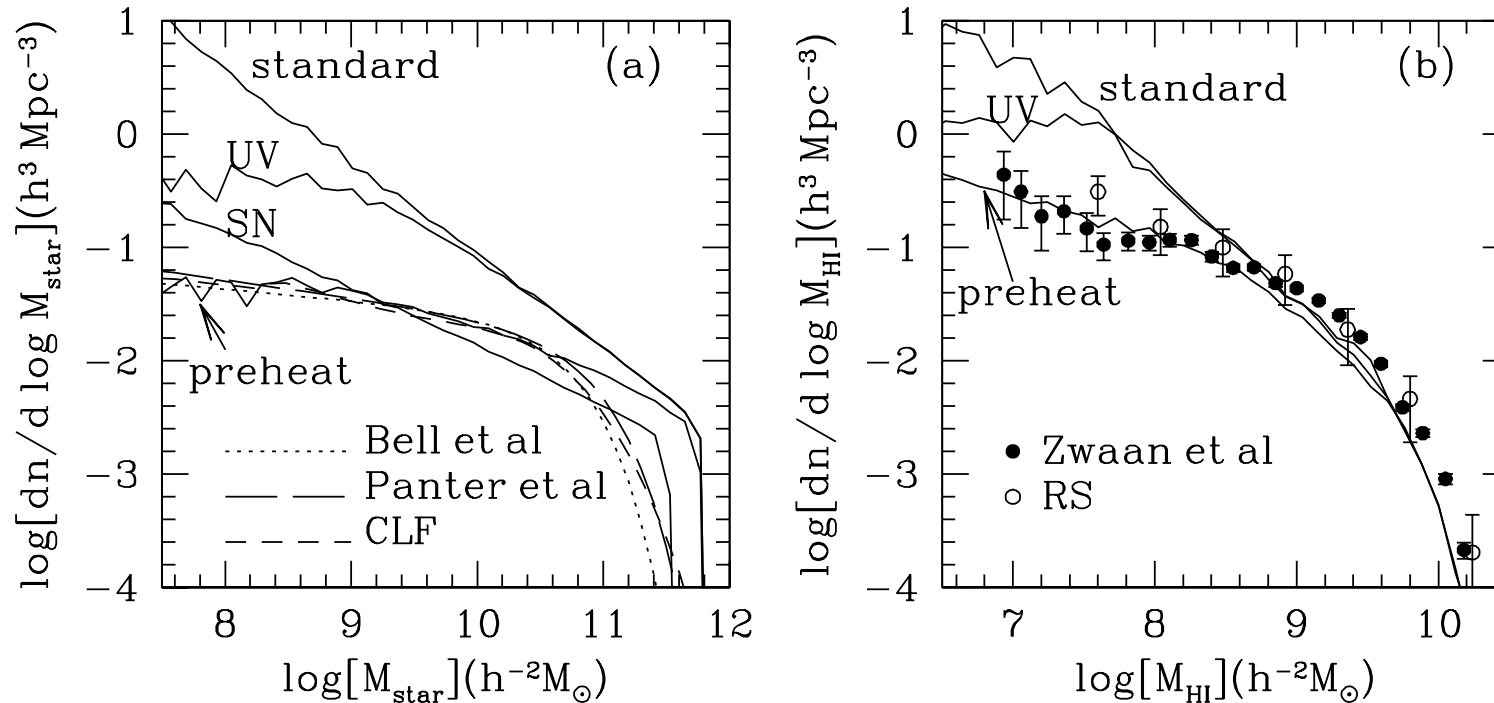
Hot Pancakes kill Dwarfs

(From van den Bosch, Abel & Hernquist 2003)



The Punchline

The **baryonic mass fractions** predicted for haloes embedded in **shock heated pancakes** are exactly what is required to fit **both** the **stellar** mass function and the **HI** mass function:



Here $M_{\text{star}} = 2\pi \int [\Sigma_{\text{disk}}(R) - \Sigma_{\text{crit}}] R dR$

and $M_{\text{gas}} = M_{\text{bar}} - M_{\text{star}}$.

The disk surface density $\Sigma_{\text{disk}}(R)$ is computed using the standard MMW model for disk formation

Mo, Mao & White 1998

NOTE: No feedback is included here!!!

Conclusions

- Galaxy Formation in extremely **inefficient** in low mass haloes
- Supernova feedback can **not** simultaneously explain the **stellar mass function** and the **HI mass function**
- Below a redshift $z \simeq 2$, about half the baryons are in a **WHIM**. This is consistent with simulations, observations, and analytical predictions
- Most low mass haloes form inside preheated pancakes
- The predicted amount of preheating is exactly what is needed to explain the **stellar mass function** and the **HI mass function**

Supernova feedback is only efficient at $z \gtrsim 2$

At lower redshifts hot pancakes rule.

For details, see Mo, Yang, van den Bosch & Katz, 2005, MNRAS, 363, 1155