

What can we learn from one-halo clustering?

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Introduction

Subhalo Abundance Matching (SHAM) has a number of advantages over Halo Model + HOD/CLF:

- No parameterization required (except for scatter)

(e.g., Vale & Ostriker 04; Kravtsov+04; Conroy+06; Guo & White 10; Reddick+13)

- Halo Assembly Bias $b_h(r | M, z_f, \lambda, \dots)$

(Tinker+05; Gao & White 2005; Wechsler+06; Zentner 07; Dalal+08; Sunayama+16; Yao+17; Villareal+17)

- Halo Occupation Bias $\langle N_s | M \rangle \rightarrow \langle N_s | M, z_f \rangle$

(vdB+05; Zentner+05; Giocoli+10; Yao+15; Jiang & vdB 16)

- Radial segregation of subhalos/satellite galaxies

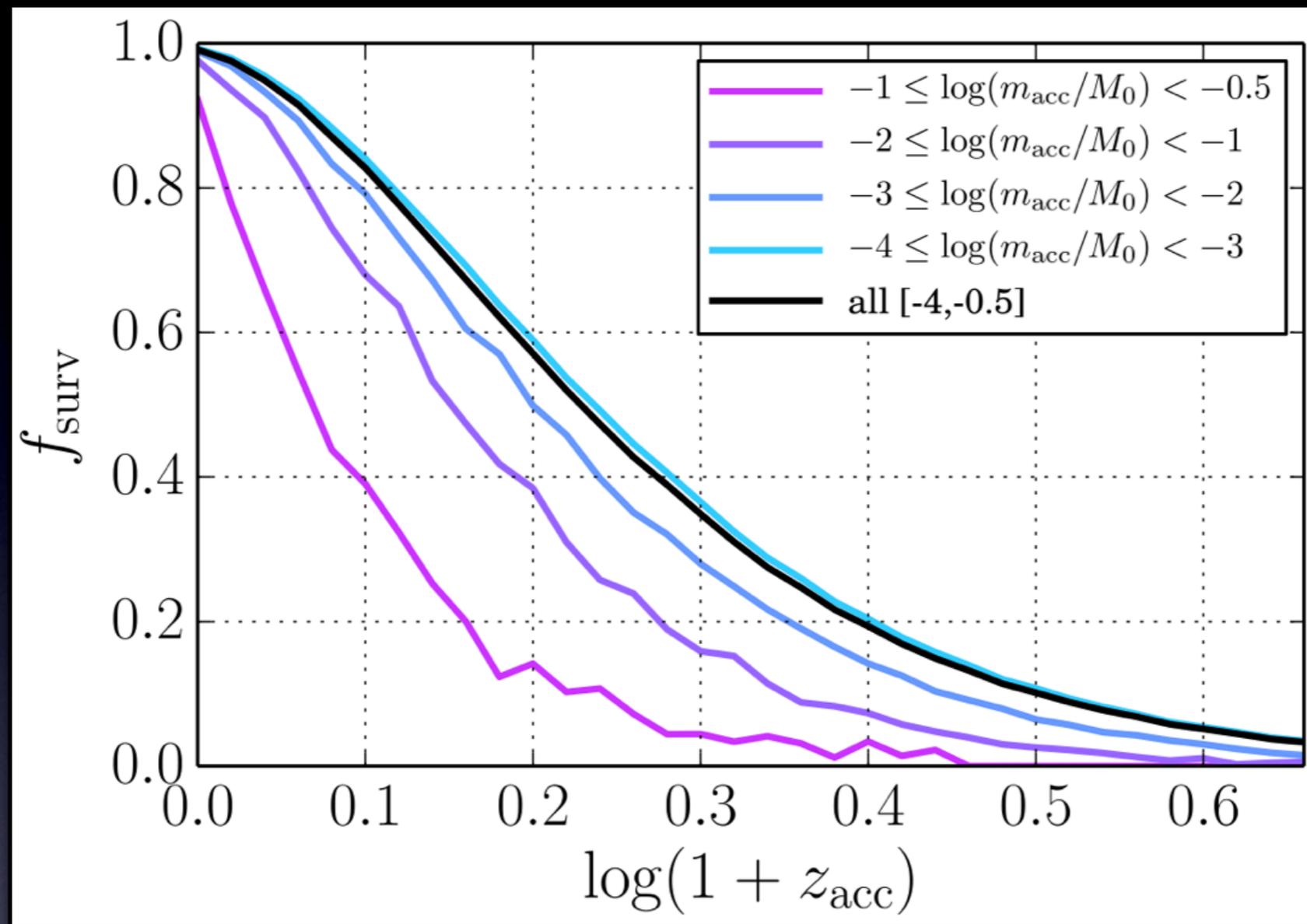
(Gao+04; Nagai & Kravtsov 05; Faltenbacher & Diemand 06; vdB+17)

- Non-Poissonian nature of $P(N_s | M)$

(Boylan-Kolchin+10; Busha+11; Wu+13; Mao+15; Jiang & vdB 16)

But: SHAM is only as accurate as the simulation used....

Subhalo Disruption in Bolshoi



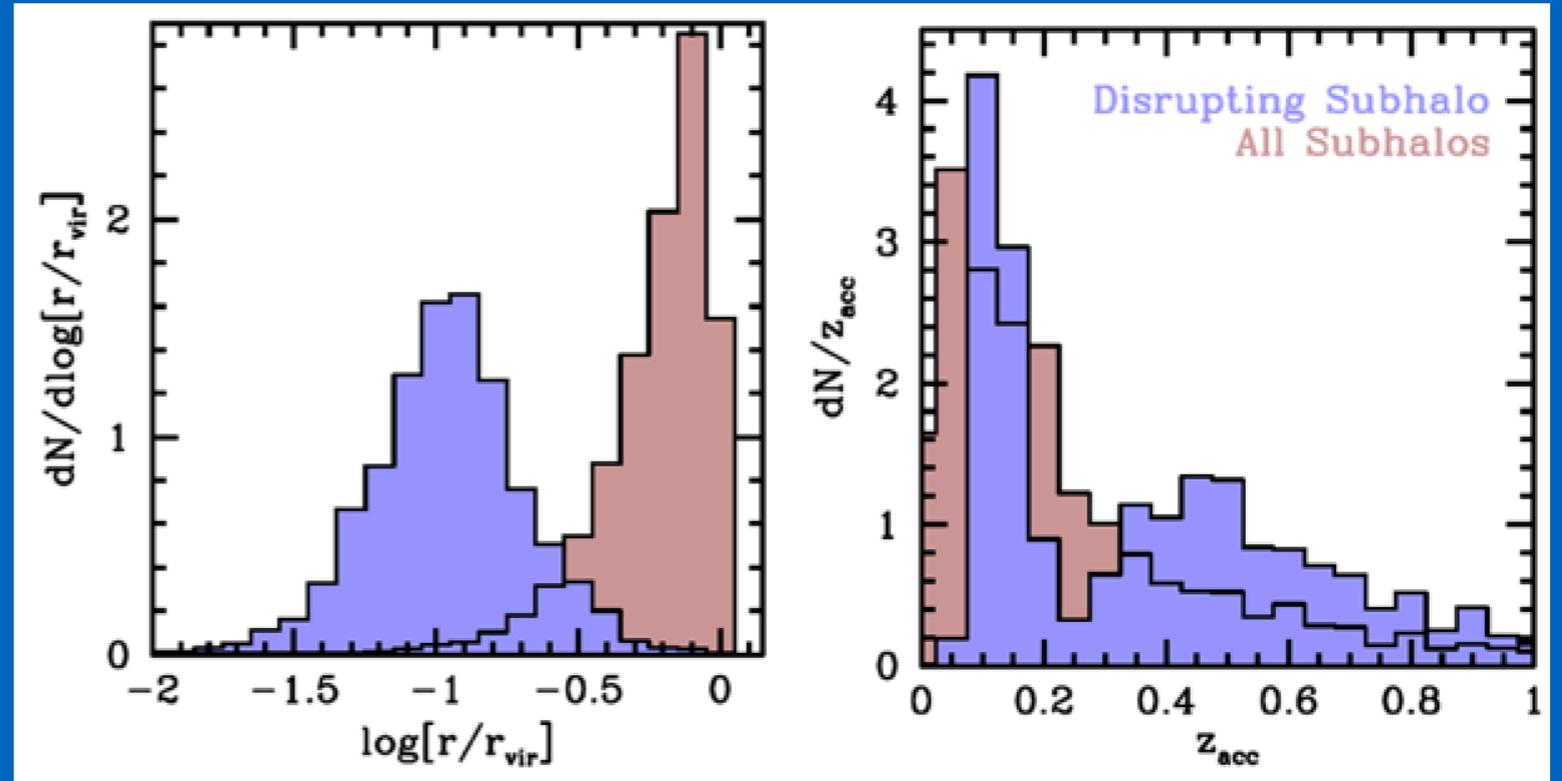
Jiang & vdB, 2016

- Fractional Disruption Rate ≈ 13 percent per Gyr
- Only ~ 35 percent of subhaloes accreted at $z=1$ survive to $z=0$
- Is tidal disruption real or numerical artifact?
If real, what are the physical conditions for disruption?

Subhalo Disruption

Disruption Demographics

- Most subhalo disruption occurs near $r \approx 0.1 r_{\text{vir},h}$
- Significant fraction of subhalos disrupt during first peri-centric passage



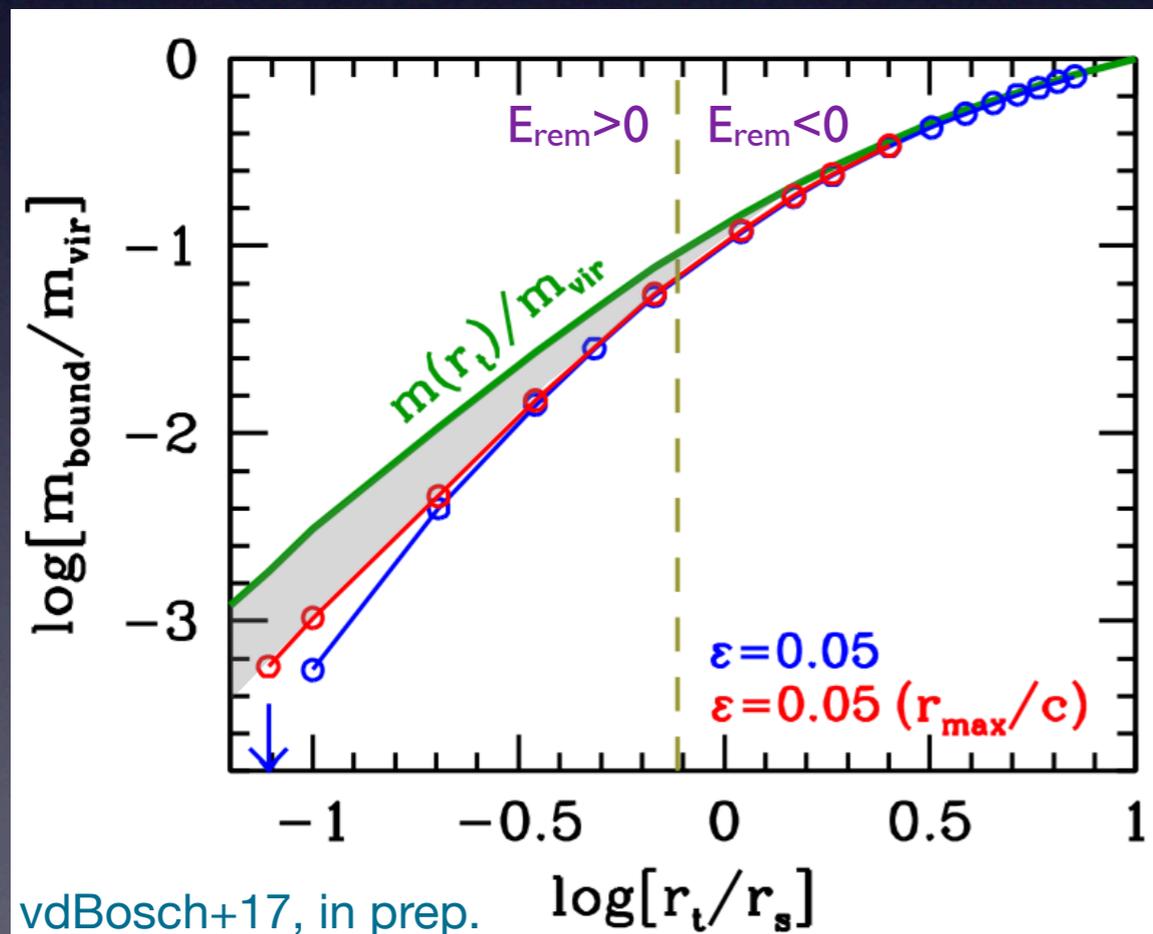
vdBosch 2017

Disruption Mechanisms

- Tidal Stripping
- Tidal Heating
 - Pericentric Passage
 - Subhalo-Subhalo Encounter
- Numerical overmerging

Does Stripping cause Disruption? NO!

- Instantaneous stripping of NFW halo can leave remnant with $E_b > 0$.
(Hayashi et al. 2003)
- For **isotropic** halo, this happens when $r_{\text{trunc}} < 0.77 r_s$ [$M(r_{\text{trunc}}) < 0.08 M_{\text{vir}}$]
- Subhalos spontaneously disintegrate once $r_{\text{tid}} < 0.77 r_s$?
This assumption is made in several models of subhalo evolution
(e.g., Zentner & Bullock 2003; Taylor & Babul 2004; Klypin et al. 2015)



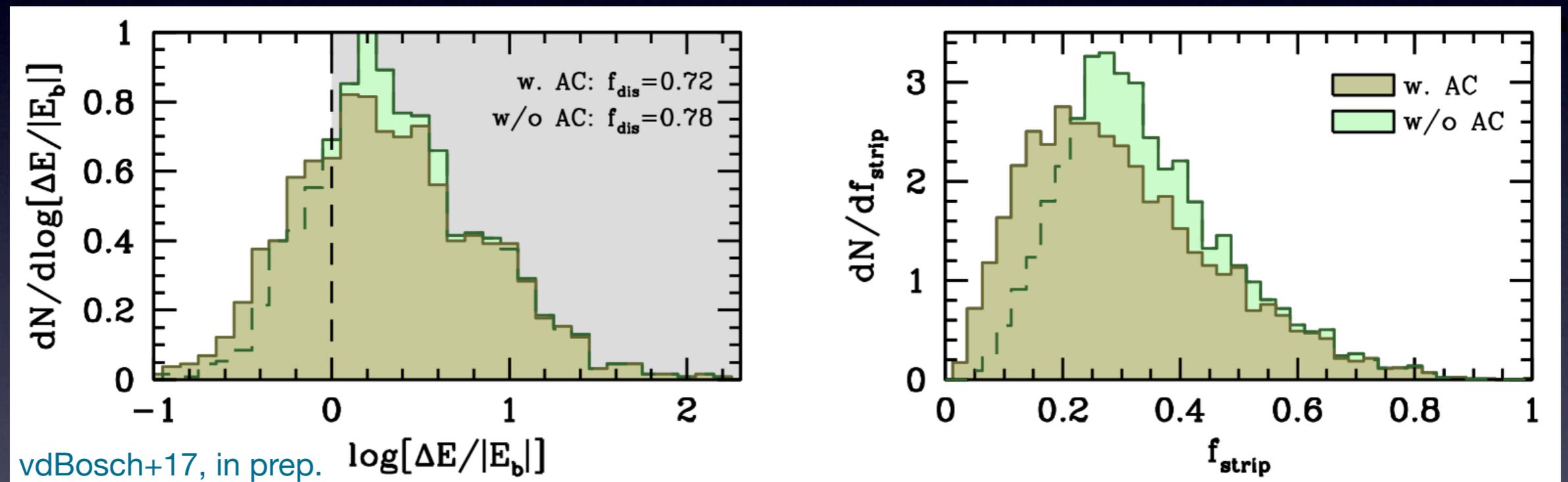
However:

particles have broad distribution of binding energies, and majority of particles remain **bound**.

Simulations confirm that remnant rapidly **re-virializes** to a **bound** system with somewhat smaller, but non-zero mass.

What about Tidal (Impulsive) Heating?

- For each subhalo in Bolshoi, compute orbital energy & circularity *at accretion*.
- Compute tidal heating, $\Delta E/|E_b|$, by integrating impulse approximation along subhalo's orbit (one period) using detailed model of Gnedin, Hernquist & Ostriker (1999).
- Apply same method to Monte-Carlo realizations of NFW subhalos to compute ΔE_i and E_i for each individual DM particle. Determine $f_{\text{strip}} = f(\Delta E_i/E_i > 1)$

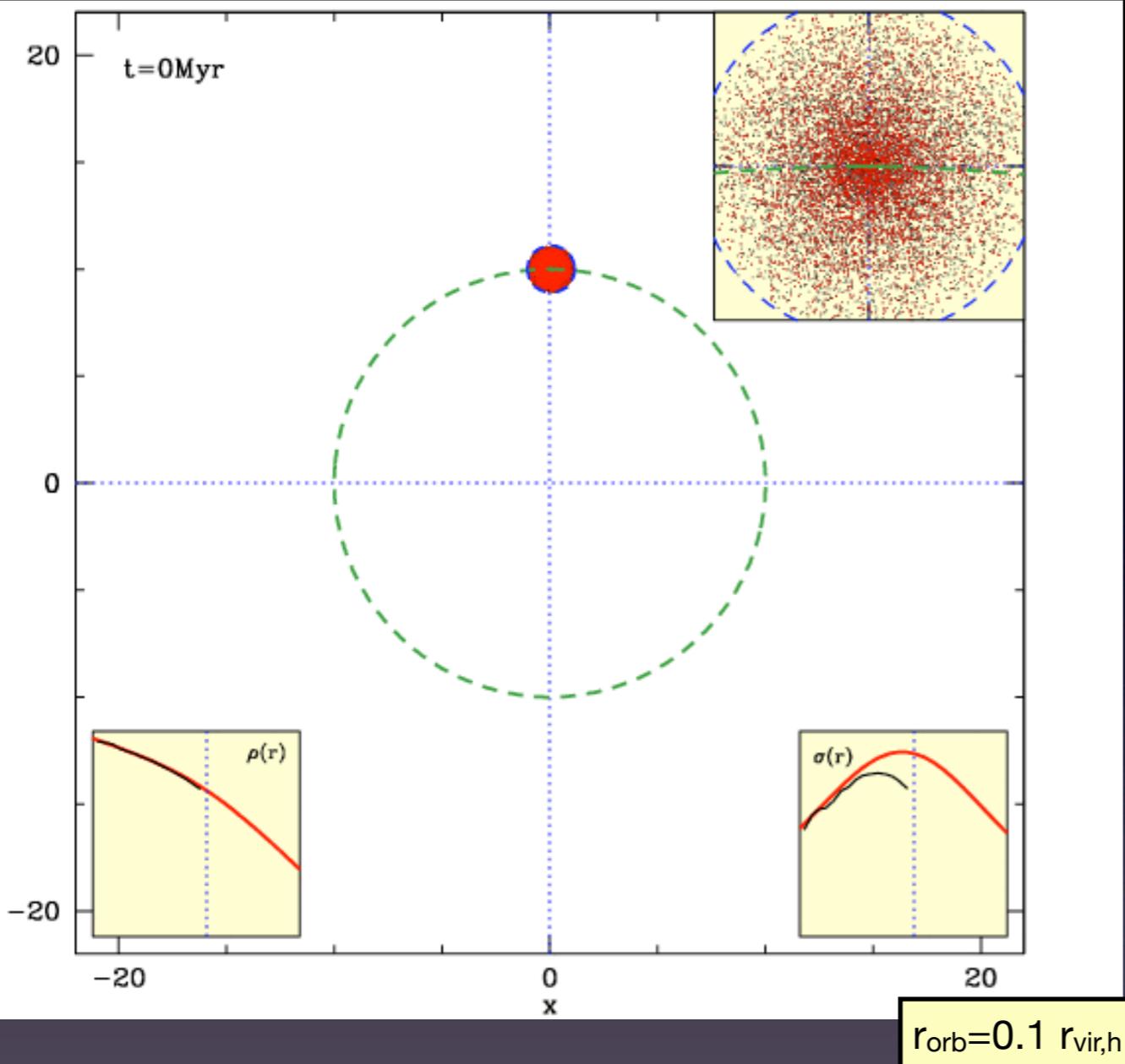


Energy input exceeds subhalo binding energy for **~80 percent** of all subhalos.

Yet, on average only **~25 percent** of subhalo particles become unbound.

Even when $\Delta E/|E_b| = 100$ as much as **20 percent** of subhalo remains bound!!!

Idealized Numerical Simulations



Simulate N-body NFW subhalo orbiting inside the **static** potential of a NFW host halo.

We consider both **circular** orbits (static tidal field; no impulsive heating) as well as **eccentric** orbits.

Goals:

- Determine under what conditions **physical** disruption occurs.
- Determine under what conditions **numerical** disruption occurs.

We have run 1000+ such simulations to determine the resolution (N, ϵ) required to properly resolve the subhalo's mass evolution, $f_{\text{bound}}(t)$.

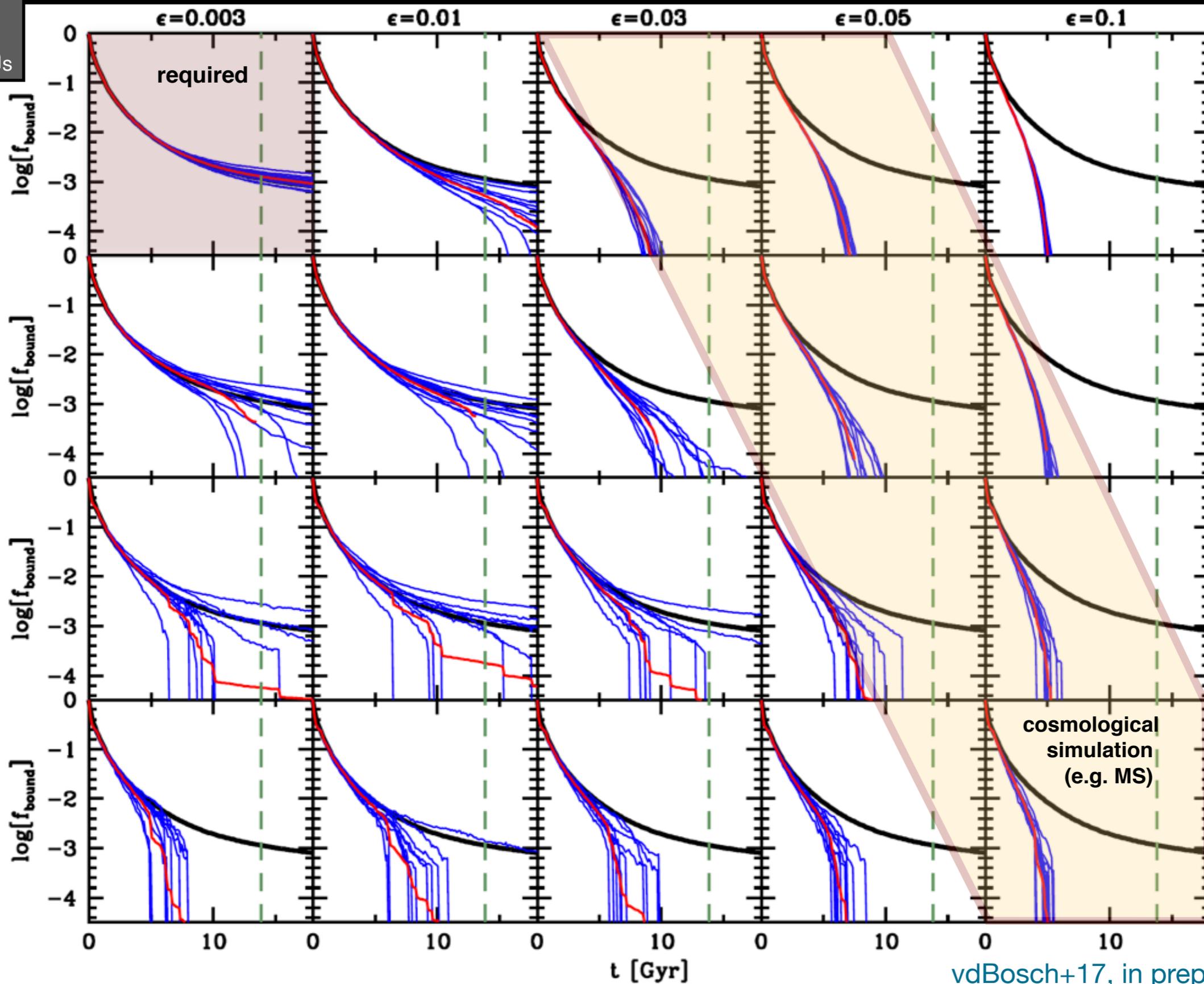
Towards Numerical Convergence

$r_{\text{orb}}=0.1$

$C_h=5$

$C_s=10$

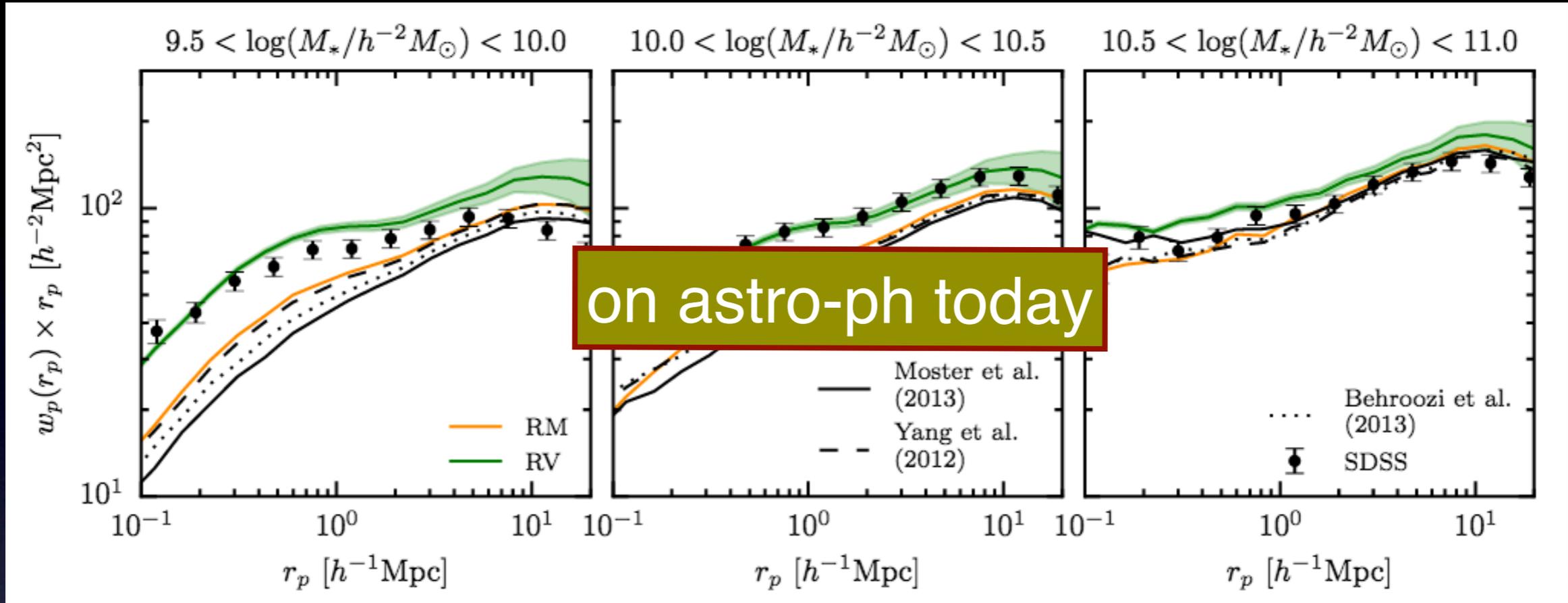
$M_h=10^3 m_s$



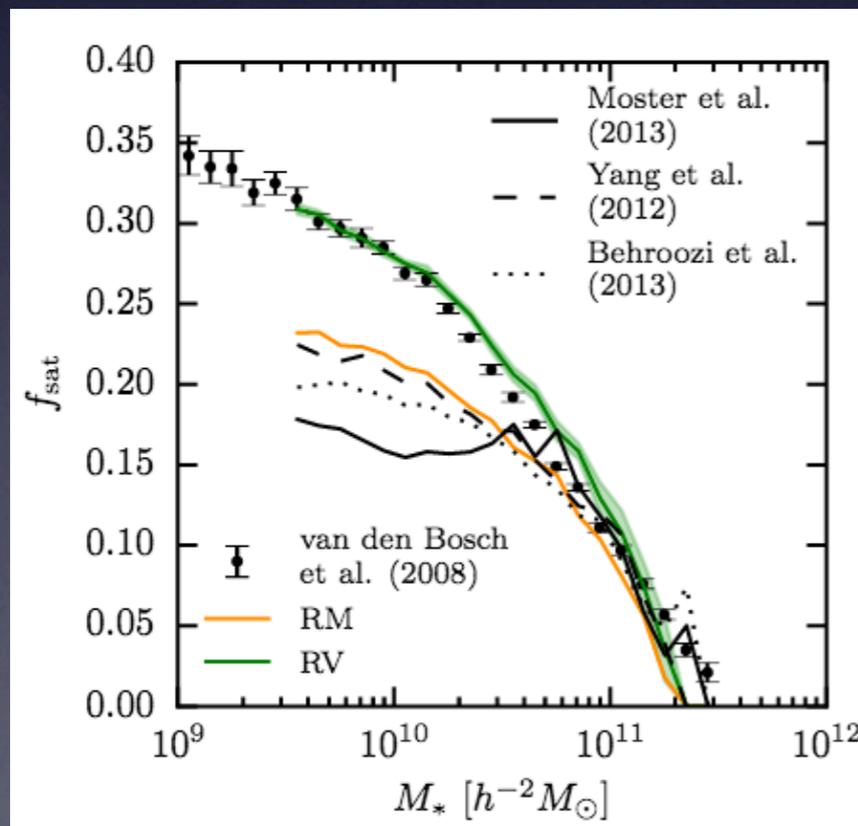
Lessons Learned

- We have explored vast areas of parameter space:
 - subhalo orbit: $E_{\text{orb}}, L_{\text{orb}}$
 - halo properties: $C_{\text{host}}, C_{\text{sub}}$
 - numerical params: $N, \varepsilon, \Delta t, \theta_{\text{tree}}$
- With sufficient numerical resolution (large N , small ε), we never find any *physical* disruption (in absence of baryonic effects & dynamical friction).
- Evolution of substructure in N-body simulations suffers from two problems:
 - Discreteness-driven run-away instability whenever $N < 1000$
 - Artificial disruption due to over-softening
- Properly resolving dynamics within inner region of host ($r \lesssim 0.2 r_{\text{vir}}$) requires very high mass resolution (large N) and force resolution (small ε).

The Galaxy Clustering Crisis

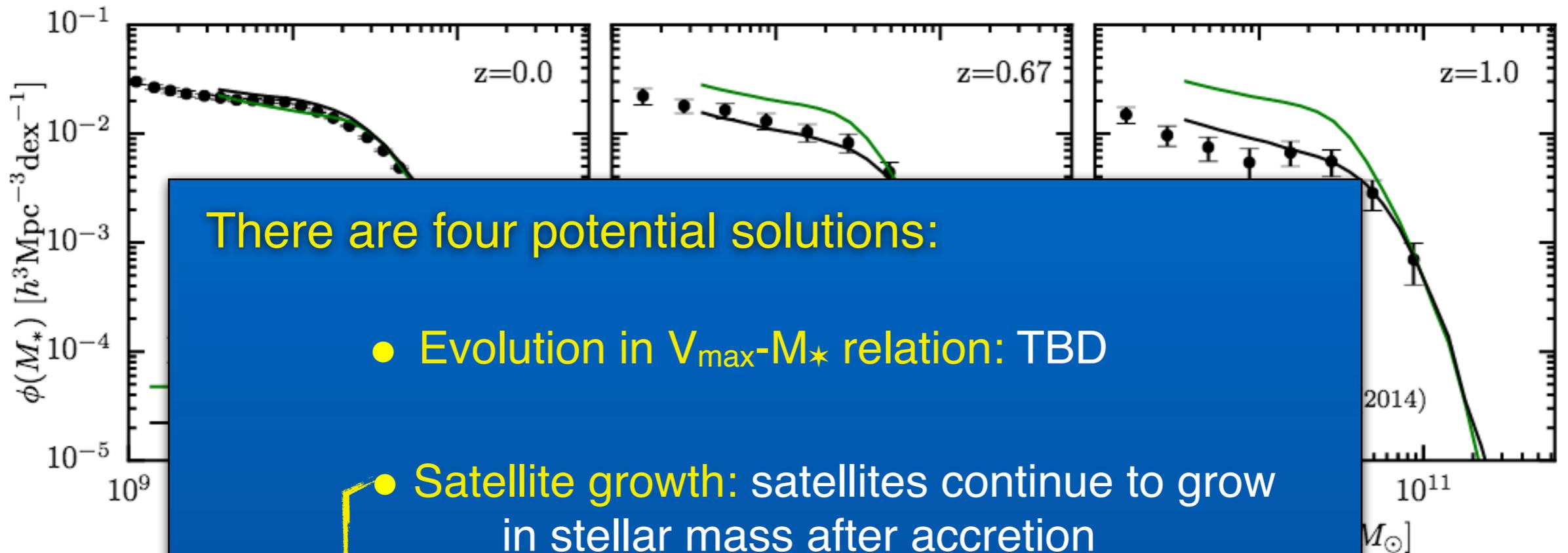


Campbell, vdB et al. 2017



- M_{peak} based SHAM (RM) dramatically underestimates clustering on small scales.
- Same is true for 'evolving SHAM' models (Yang+12, Moster+13, Behroozi+13)
- V_{peak} based SHAM (RV) fits clustering well (e.g. Conroy+06, Reddick+13)

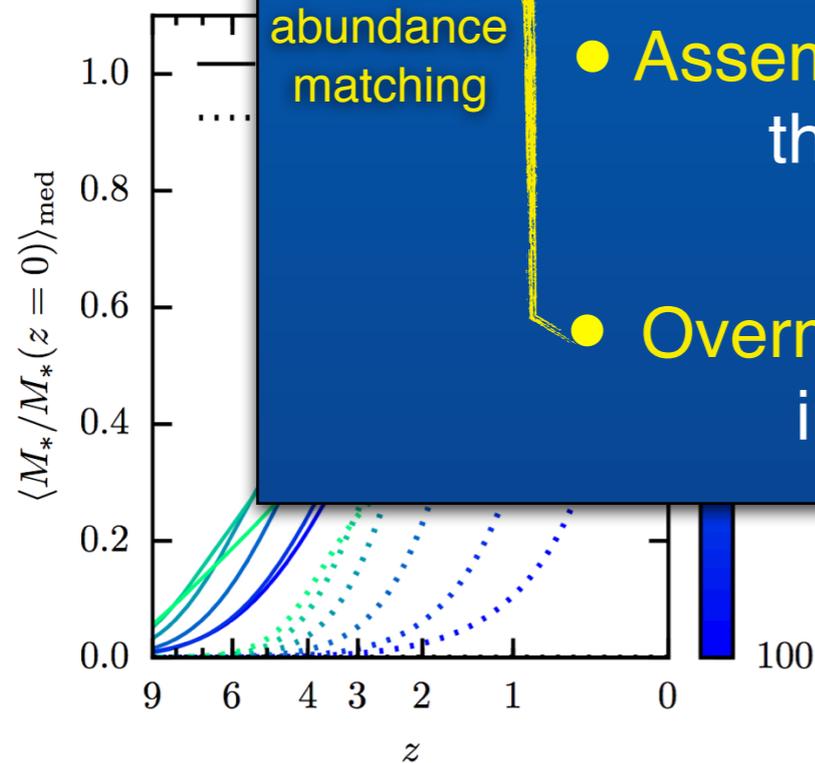
The Galaxy Clustering Crisis



There are four potential solutions:

- **Evolution in V_{\max} - M_* relation:** TBD
(Effect not strong enough)
- **Satellite growth:** satellites continue to grow in stellar mass after accretion
(No success yet...really difficult)
- **Assembly Bias:** satellites are more massive than centrals at moment of accretion
(requires orphans)
- **Overmerging:** artificial disruption of subhalos in numerical simulations

M-based abundance matching

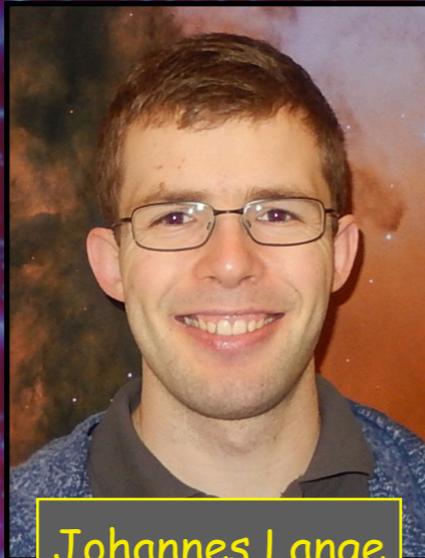


predict that stellar masses are assembled early, overpredicting $\Phi(M_*)$ at high z .

Conclusions

- ★ Current generation of cosmological simulations still suffers from severe **overmerging**.
(cf., Kazantzidis, Moore & Mayer 2004)
 - serious road-block for **small-scale cosmology program**
 - serious road-block for understanding **galaxy formation**
- ★ SHAM needs to include **'orphan galaxies'**, but without knowing how many orphans to add, and where, the information content of the 1-halo term is extremely limited.

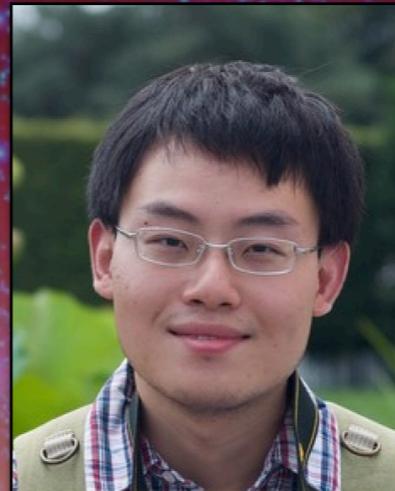
with thanks to my collaborators



Johannes Lange



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Go Ogiya



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