

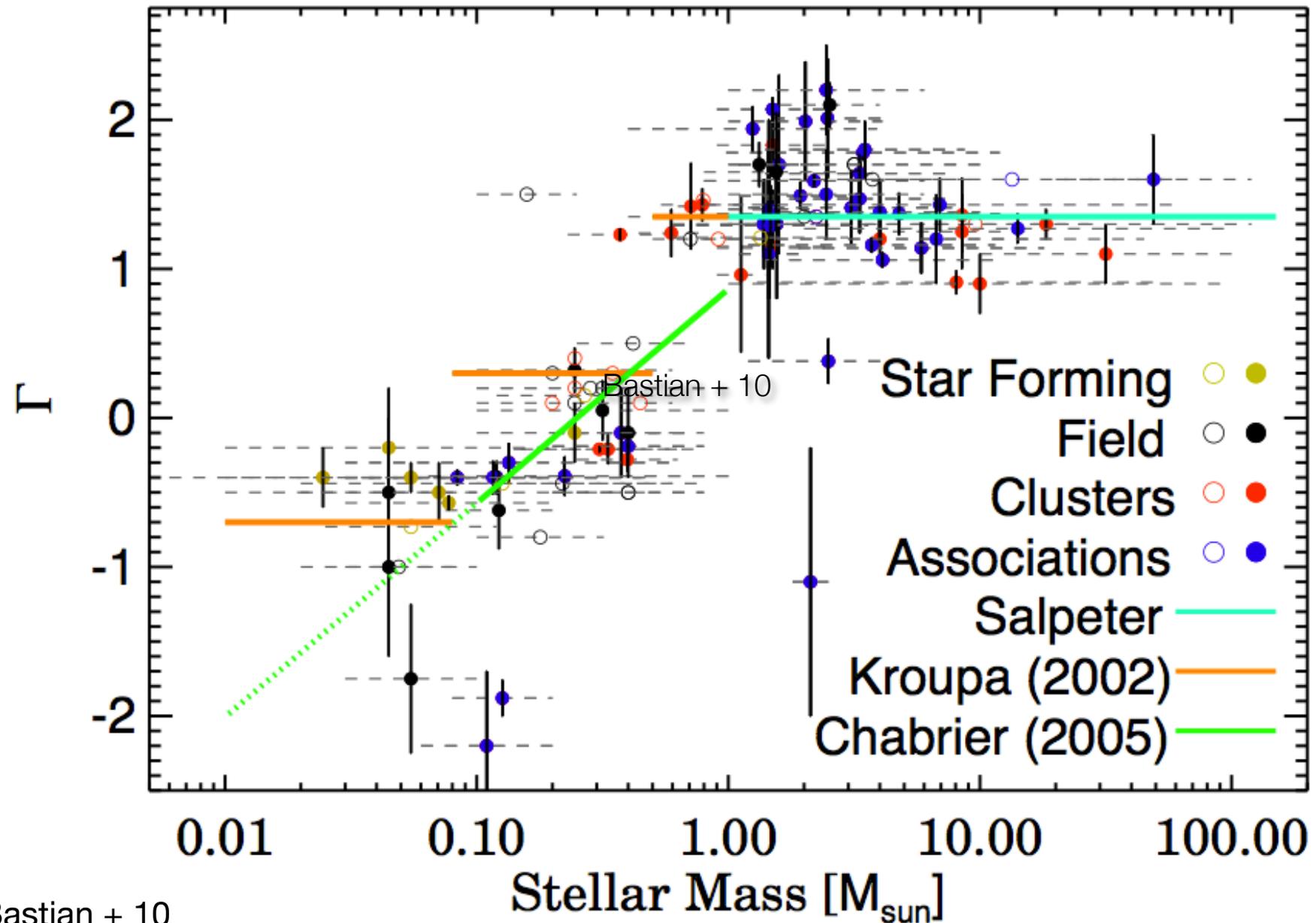
# Towards a Complete Picture of Galactic Star Formation: Theory

Mordecai-Mark Mac Low

Department of Astrophysics



# Initial Mass Function

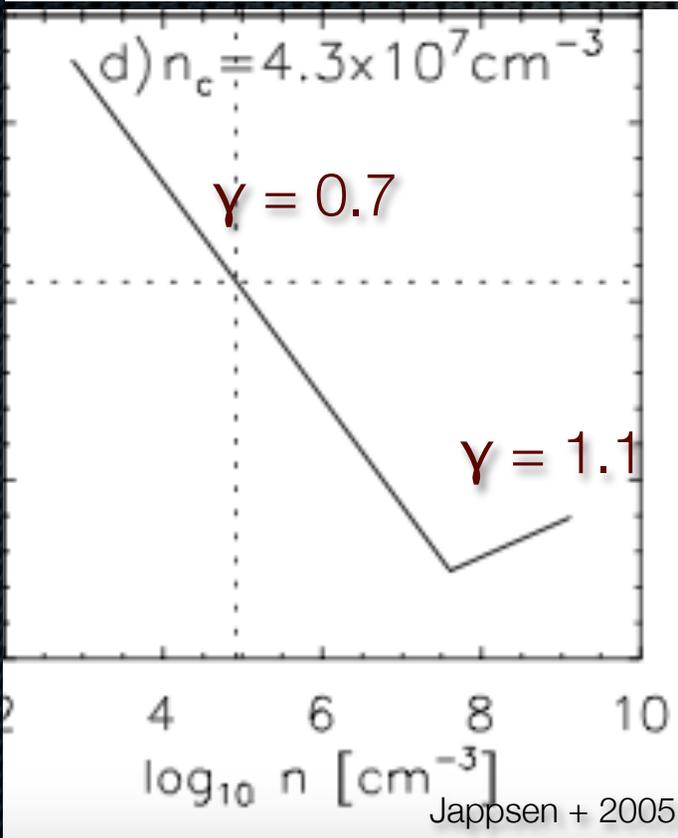


Thermodynamics can set  
the fragmentation scale.

Larson 2005

$$T = 4.4 \rho_{18}^{-0.27} \text{ K}, \quad \rho < 10^{-18} \text{ g cm}^{-3}$$

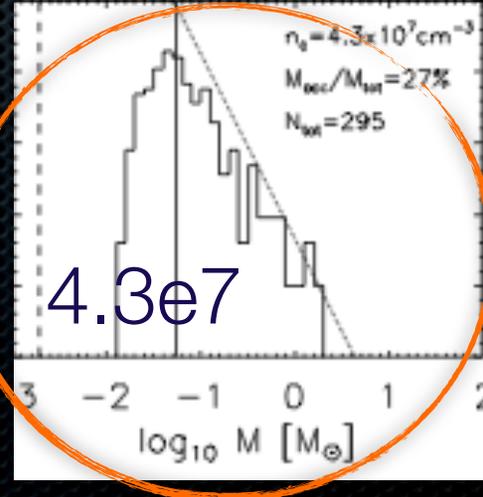
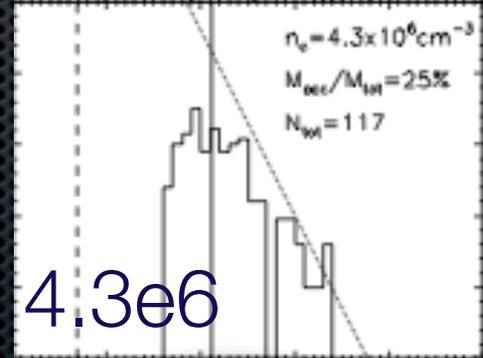
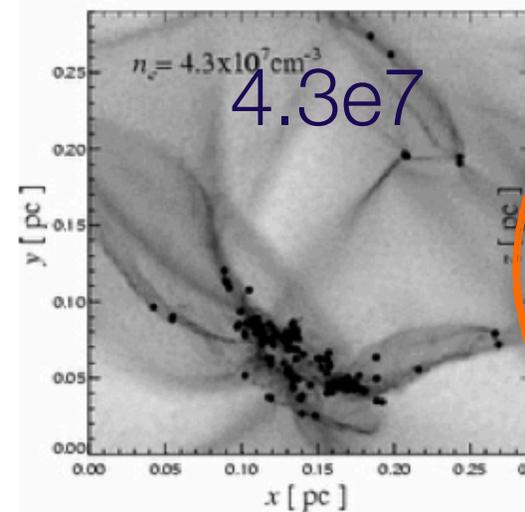
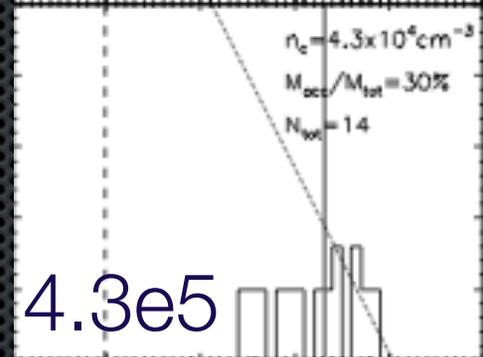
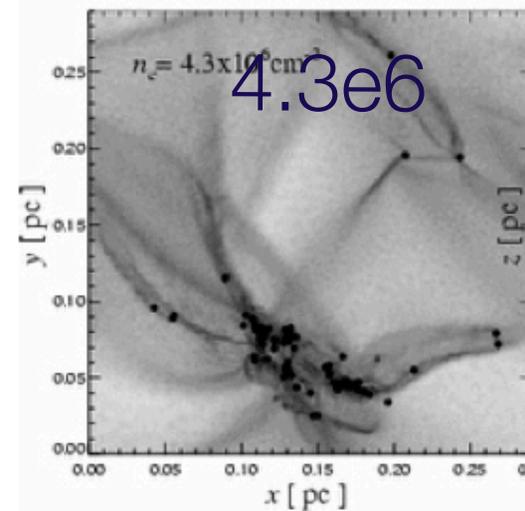
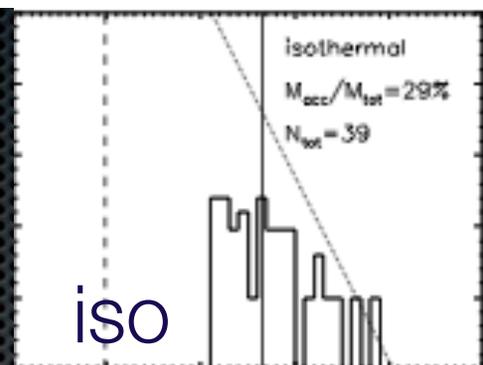
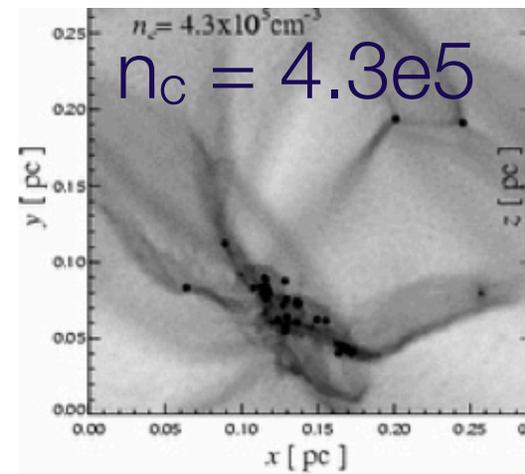
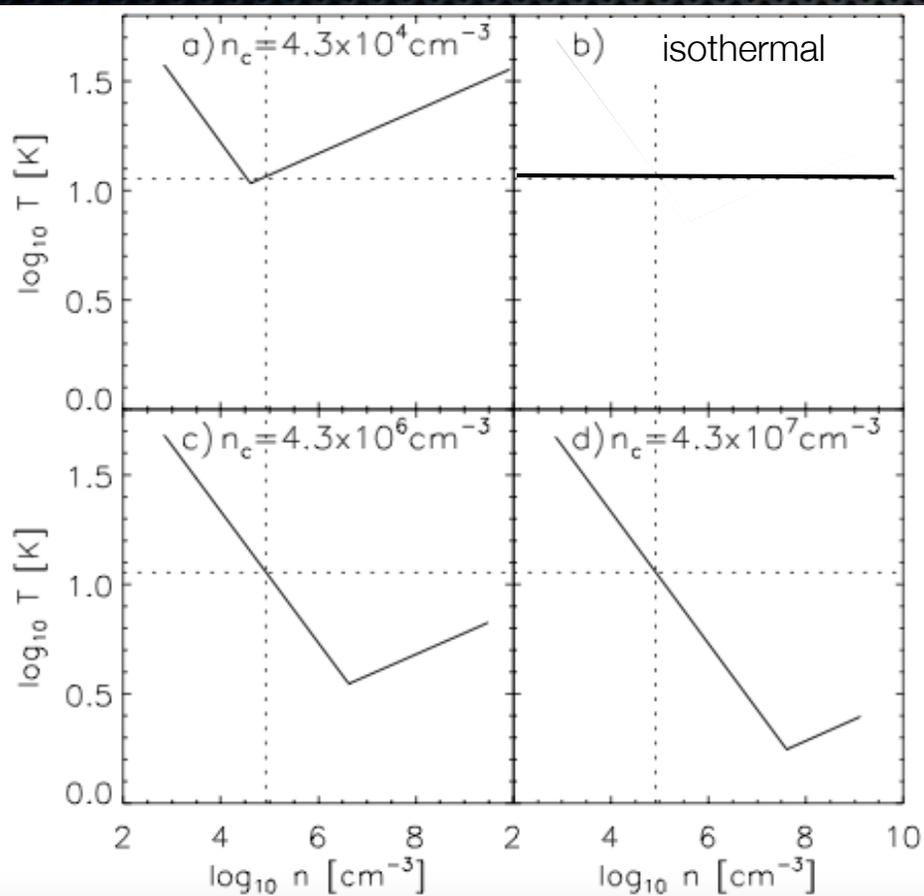
$$T = 4.4 \rho_{18}^{+0.07} \text{ K}, \quad \rho > 10^{-18} \text{ g cm}^{-3}$$



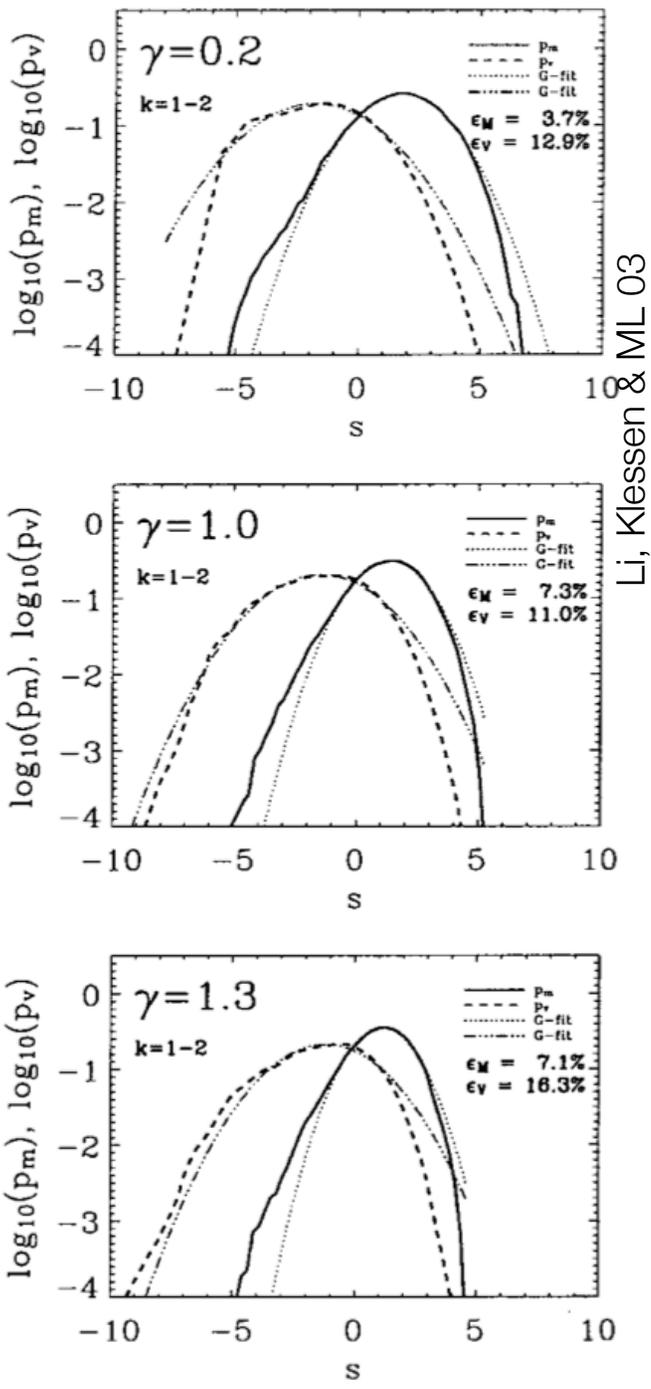
Cylinders  
fragment if  $\gamma < 1$   
(Mestel 1965a)



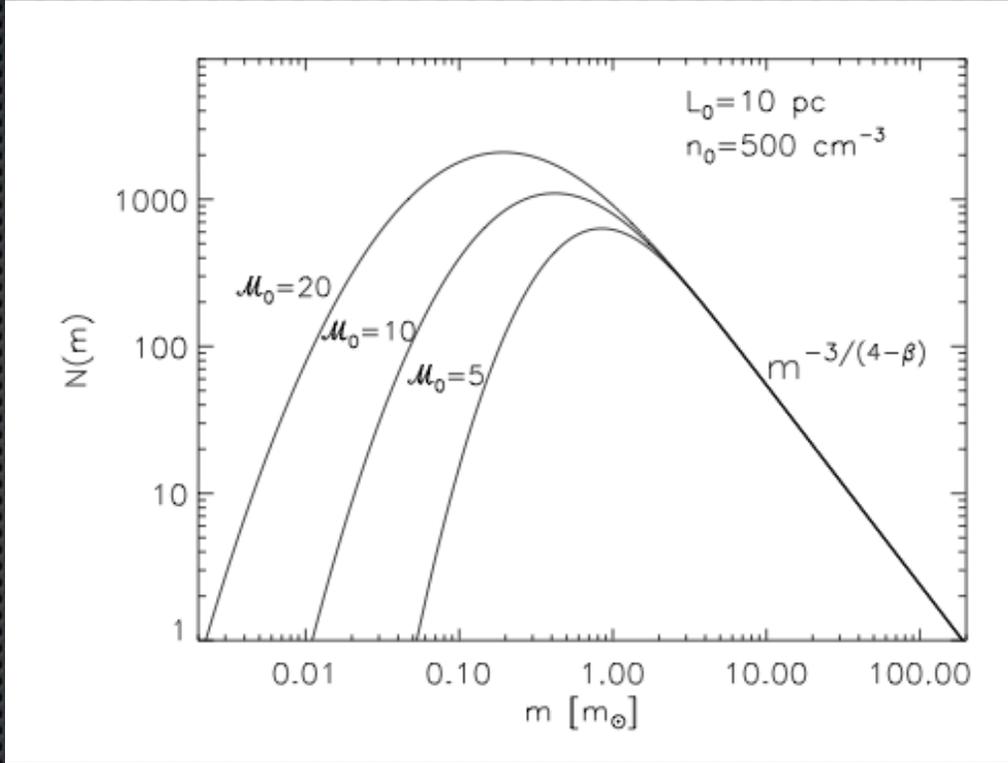
Increasing the transition density to realistic values predicts IMF that agrees with observations



Turbulent fragmentation also argued to determine IMF (Padoan 95).



Li, Klessen & ML 03



Padoan & Nordlund 02

PDF used to derive fraction of mass in clumps characterized by post-shock density

see Ostriker + 01 for isothermal model

However, it has become clear that a hierarchical description of collapse is required to capture the behavior of molecular clouds

## The Salpeter Slope of the IMF Explained

Sally Oey

University of Michigan, Department of Astronomy, Ann Arbor, MI 48109-1042

**Abstract:** If we accept a paradigm that star formation is a self-similar, hierarchical process, then the Salpeter slope of the IMF for high-mass stars can be simply and elegantly explained as follows. If the intrinsic IMF at the smallest scales follows a simple,  $-2$  power-law slope, then the steepening to the  $-2.35$  Salpeter value results when the most massive stars cannot form in the lowest-mass clumps of a cluster. This steepening must occur if clusters form hierarchically from clumps, and the lowest-mass clumps can form stars. This model is consistent with a variety of observations as well as theoretical simulations.

For more details, see:

Oey 2011, *ApJL* 739, L46

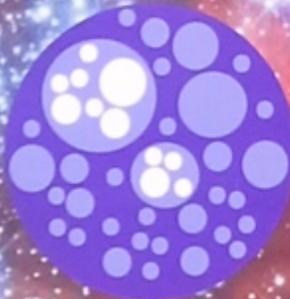


Figure 1: Clumps in a cloud follow a  $\Gamma = -1$  MF as the stars which form

The highest-mass stars cannot form in the lowest-mass clumps if their masses are

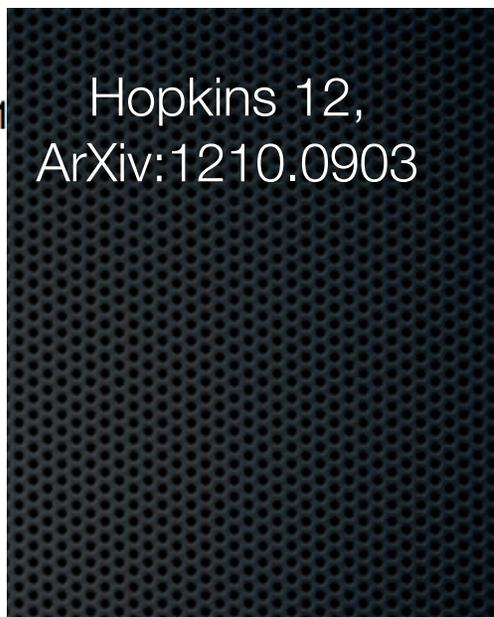
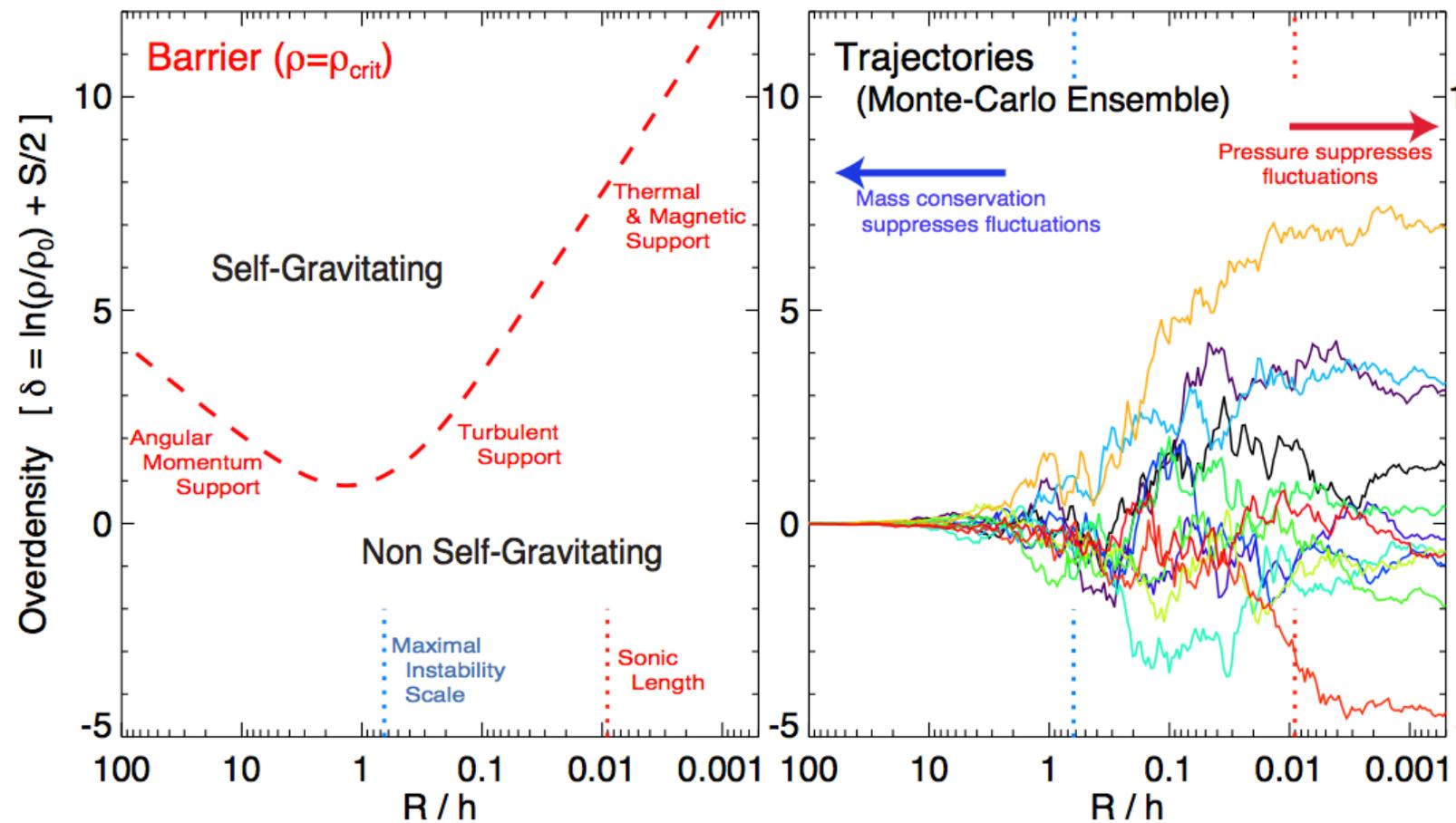
Box 1.  $\gamma = -2$  (linear) or  $\Gamma = -1$  (log):

However, it has become clear that a hierarchical description of collapse is required to capture the

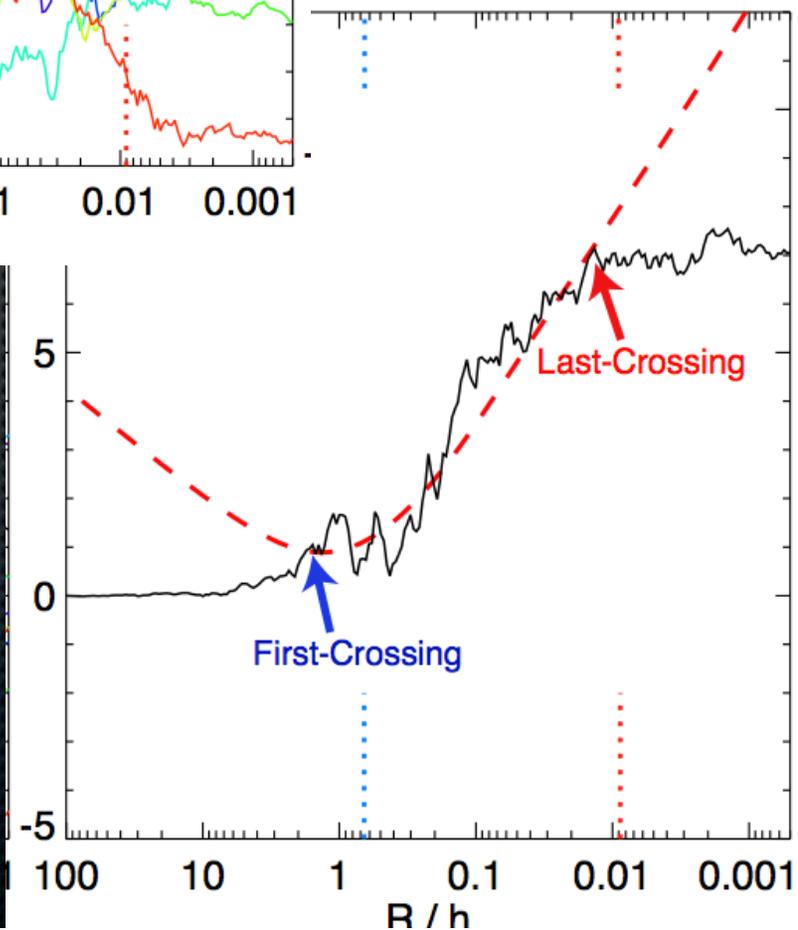
**Summary.** Data for many molecular clouds and condensations show that the internal velocity dispersion of each region is well correlated with its size and mass, and these correlations are approximately of power-law form. The dependence of velocity dispersion on region size is similar to the Kolmogoroff law for subsonic turbulence, suggesting that the observed motions are all part of a common hierarchy of interstellar turbulent motions. The regions studied are mostly gravitationally bound and in approximate virial equilibrium. However, they cannot have formed by simple gravitational collapse, and it appears likely that molecular clouds and their substructures have been created at least partly by processes of supersonic hydrodynamics. The hierarchy of subcondensations may terminate with objects so small that their internal motions are no longer supersonic; this predicts a minimum protostellar mass of the order of a few tenths of a solar mass. Massive ‘protostellar’



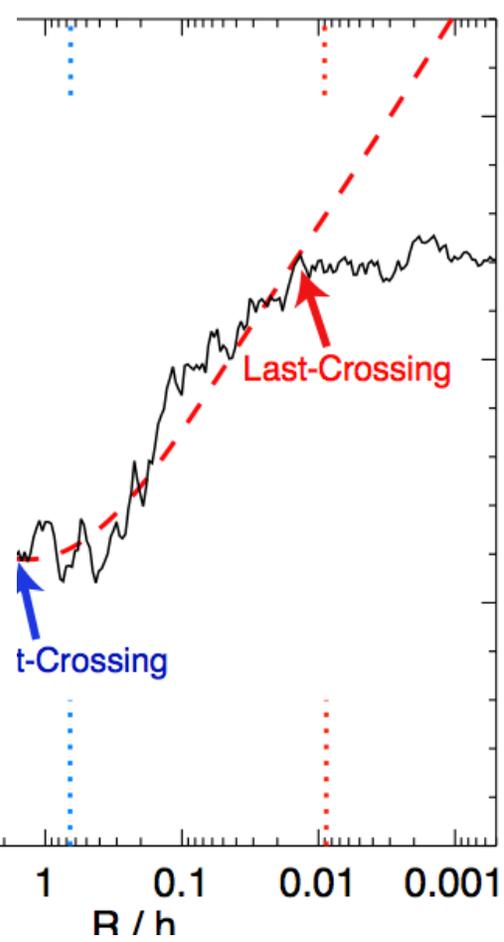
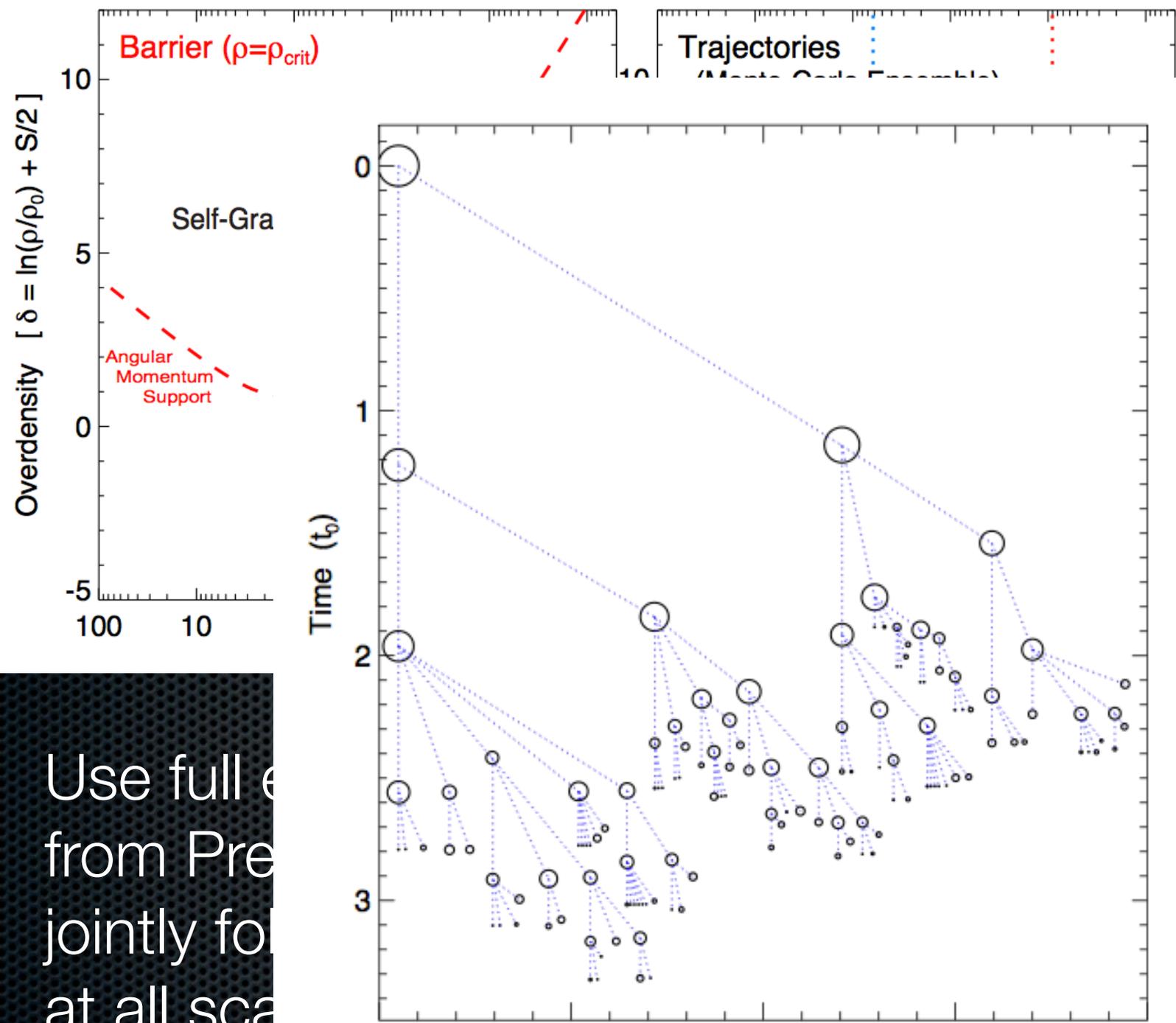
Larson 81



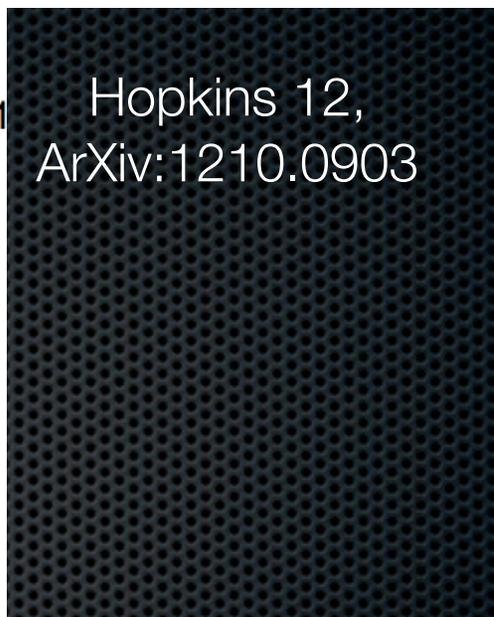
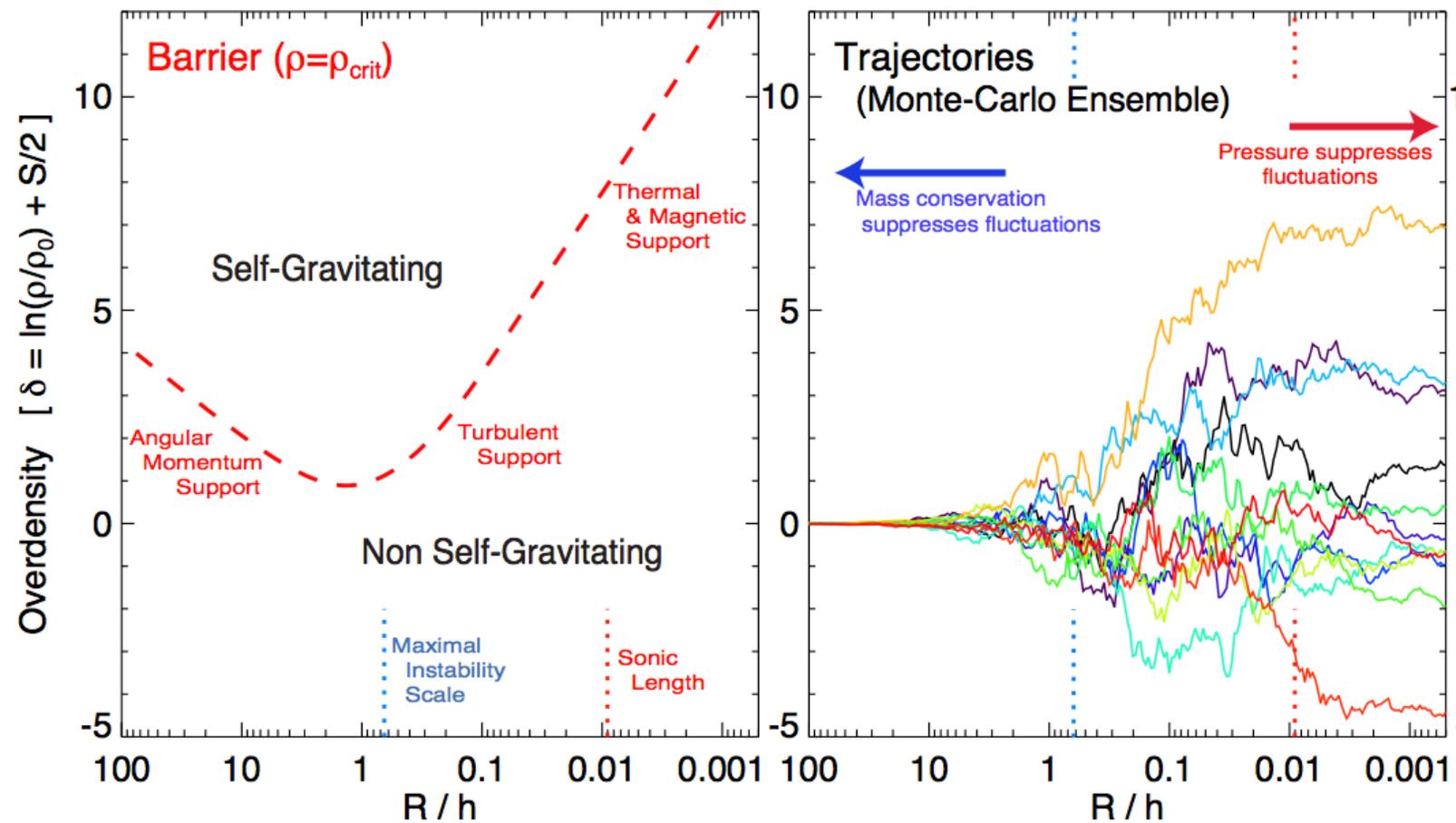
Use full excursion set formalism from Press-Schechter theory to jointly follow hierarchical collapse at all scales simultaneously



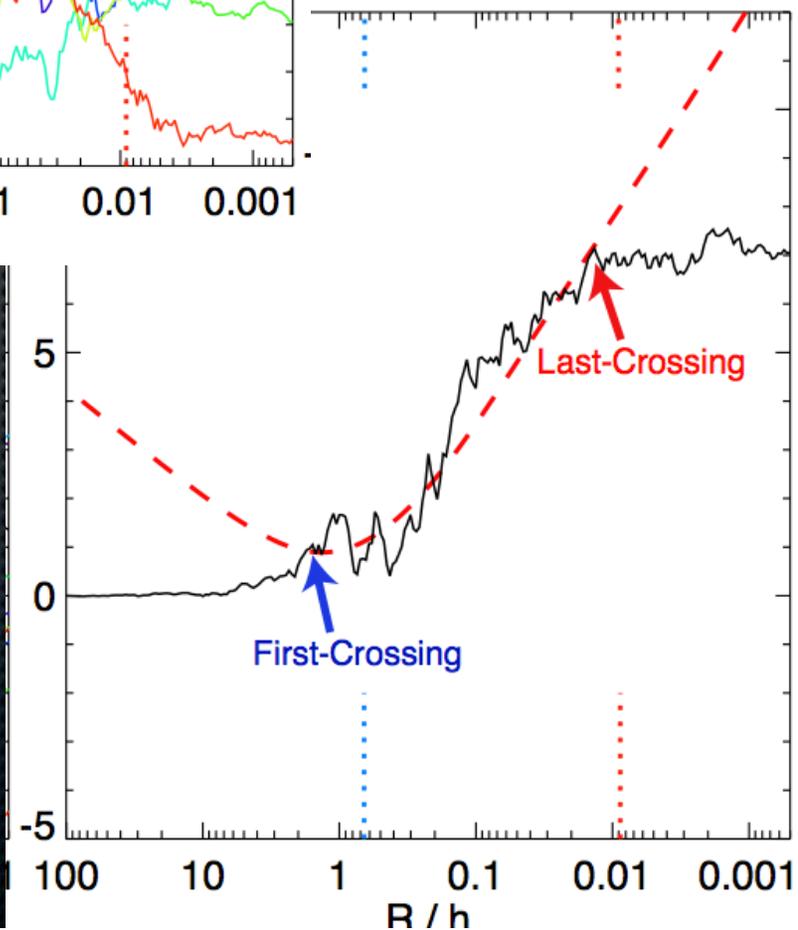
Hopkins 12,  
ArXiv:1210.0903

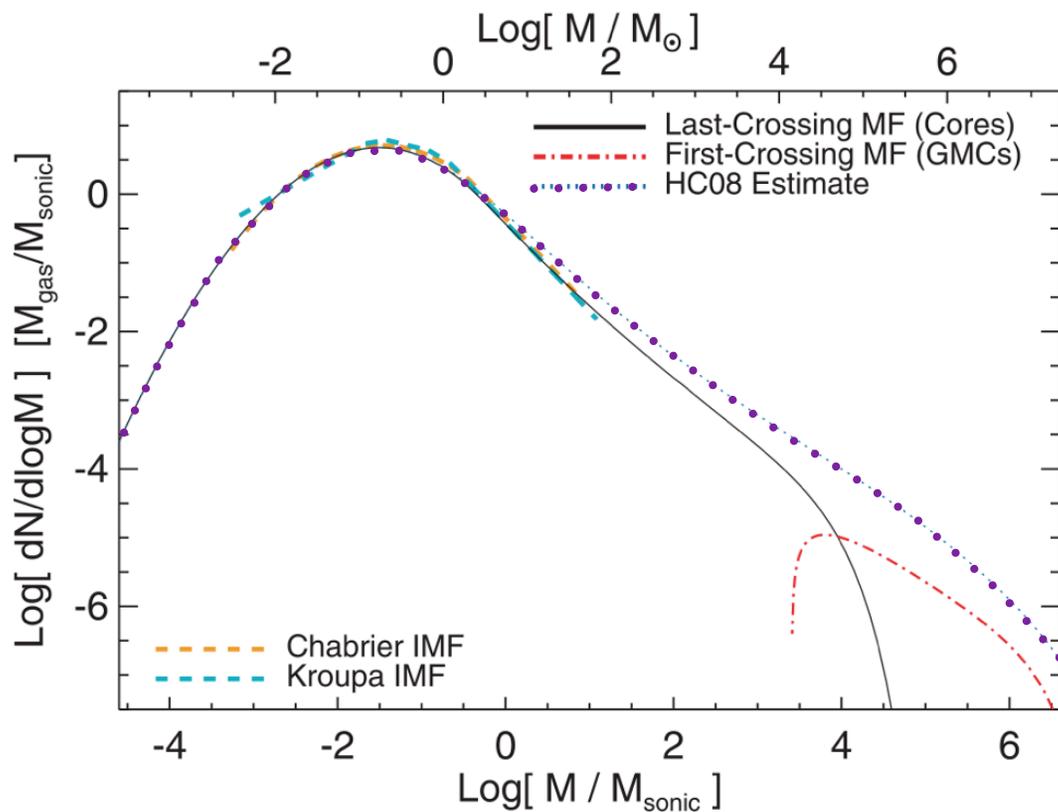


Hopkins 12, ArXiv 1210.0903

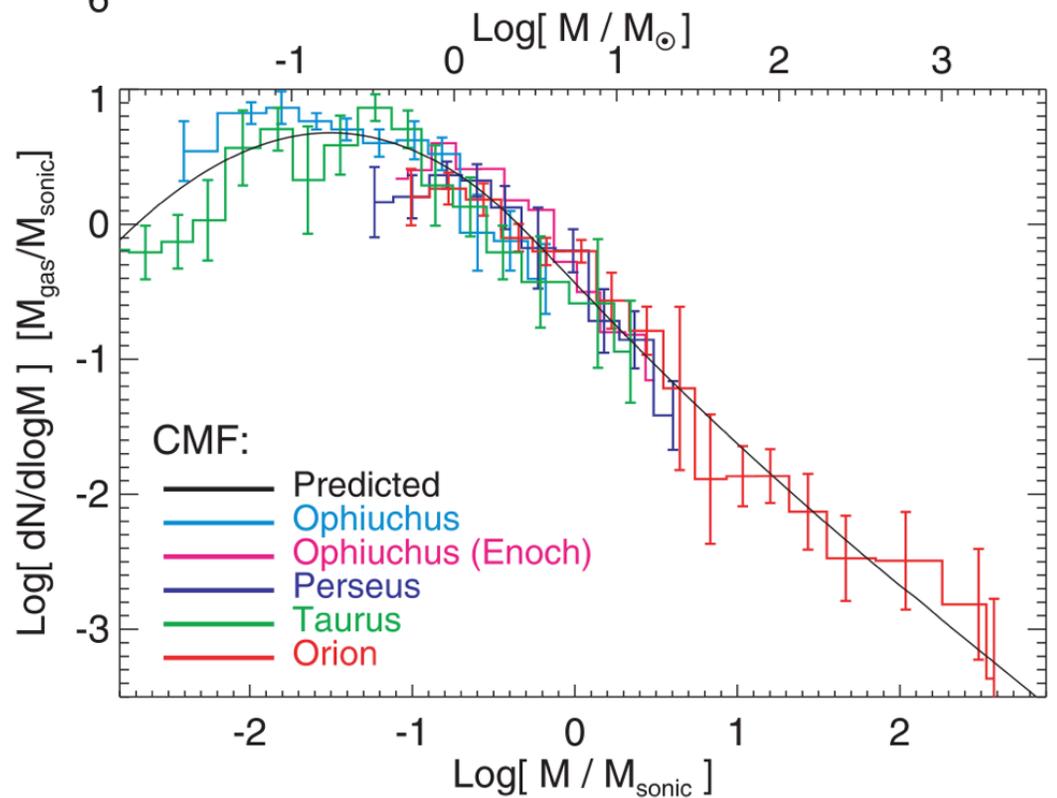


Use full excursion set formalism from Press-Schechter theory to jointly follow hierarchical collapse at all scales simultaneously



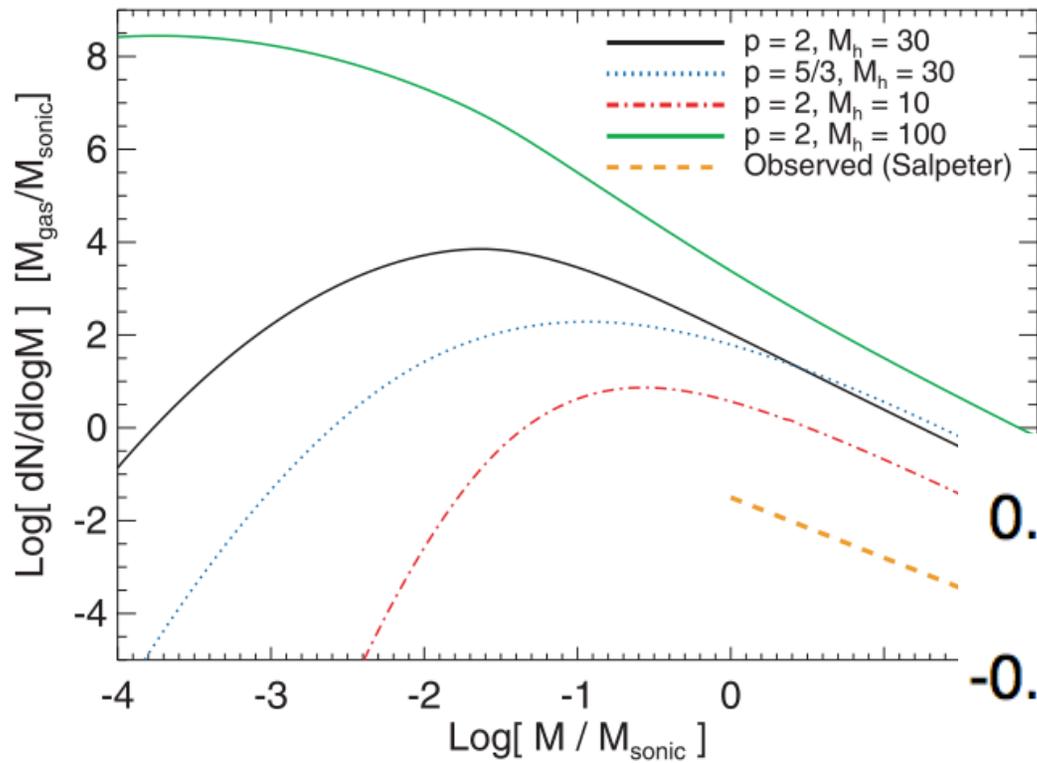


Hopkins 12  
(MNRAS 423, 2037)



Hennebelle & Chabrier 08  
derived mass distribution  
ignoring multiple barrier  
crossings.

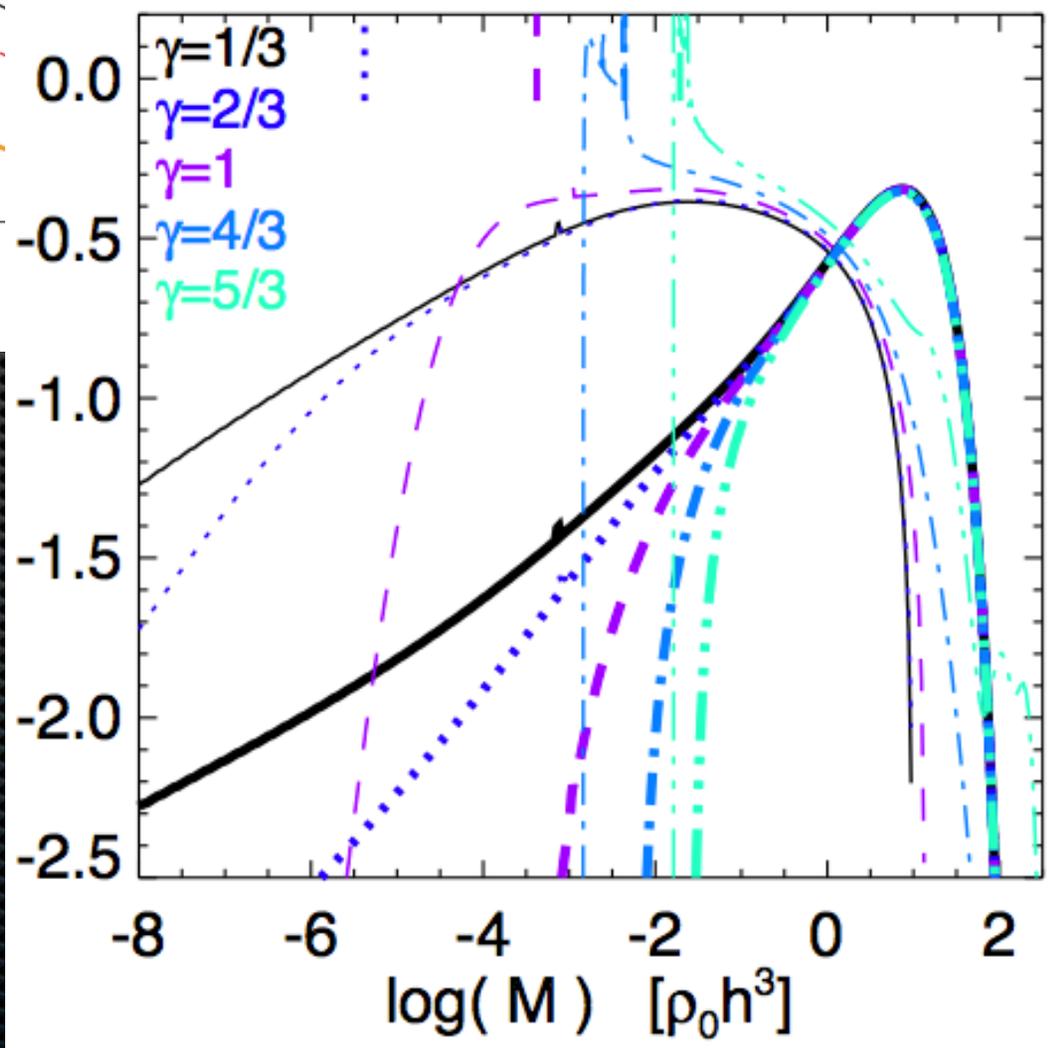
Hopkins 12 examines first, last  
crossings as separate  
distributions.

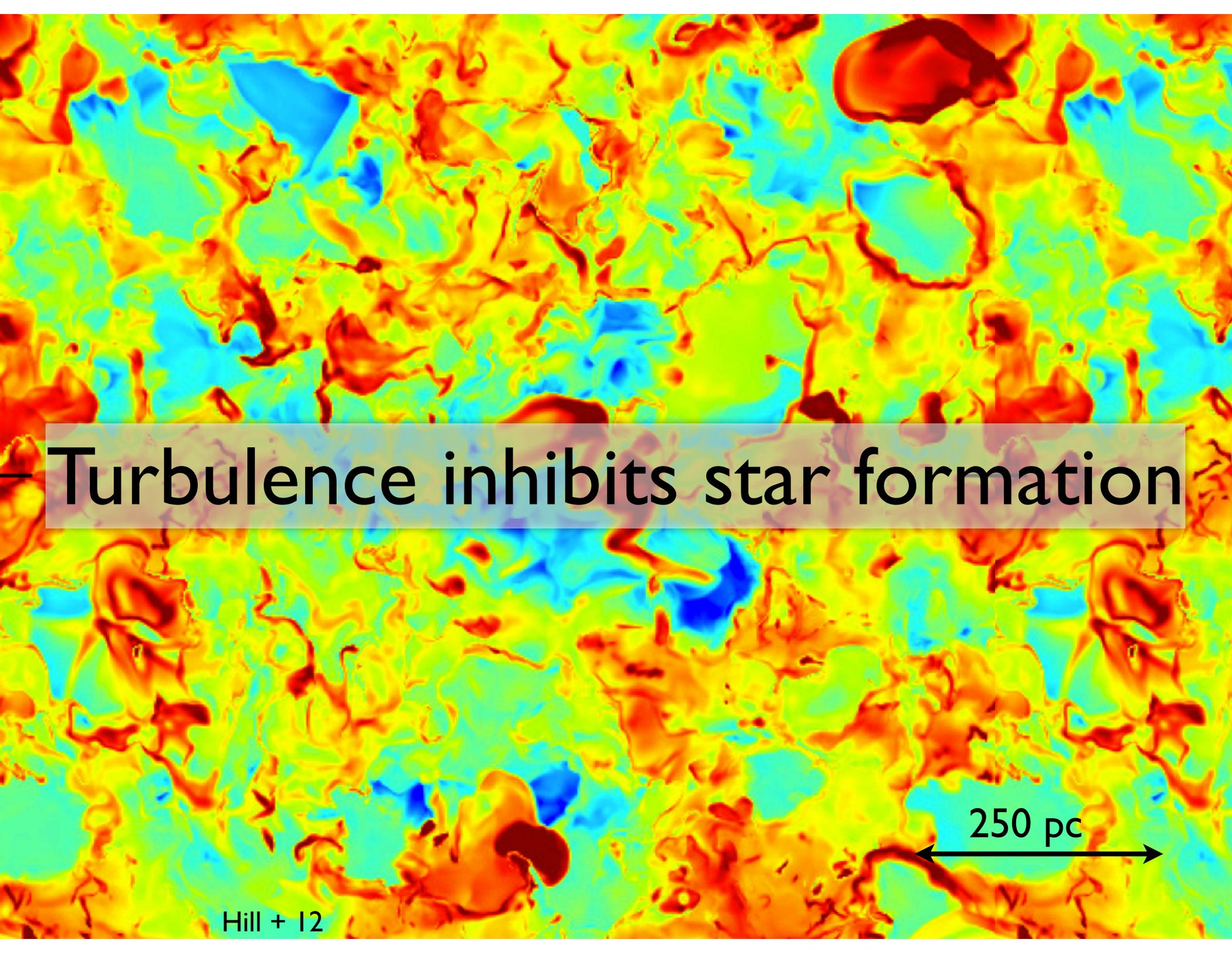


Hopkins 12, ArXiv 1210.0903

Hopkins 12 (MNRAS 423, 2037)

Formalism can be generalized to include non-isothermal EOS, varying turbulent properties, other physics.





Turbulence inhibits star formation

250 pc

# Turbulence *Prevents* Collapse

- Turbulent motions can be treated as an additional pressure (Chandrasekhar 1951, von Weizsäcker 1951)

$$c_{s,\text{eff}}^2 = c_s^2 + \frac{\langle v^2 \rangle}{3}$$

- Supersonic turbulence increases the mass supported against collapse

$$M_J = \left( \frac{\pi}{G} \right)^{3/2} \rho^{-1/2} c_{s,\text{eff}}^3$$

# Turbulence *Promotes* Collapse

- Supersonic turbulence drives shock waves that produce density enhancements.
- In isothermal gas, the postshock density increases with the Mach number  $M$  as

$$\rho_s = \rho M^2$$

- Supersonic turbulence decreases the mass supported against collapse

$$M_J = \left( \frac{\pi}{G} \right)^{3/2} \rho_s^{-1/2} c_{s,\text{eff}}^3$$

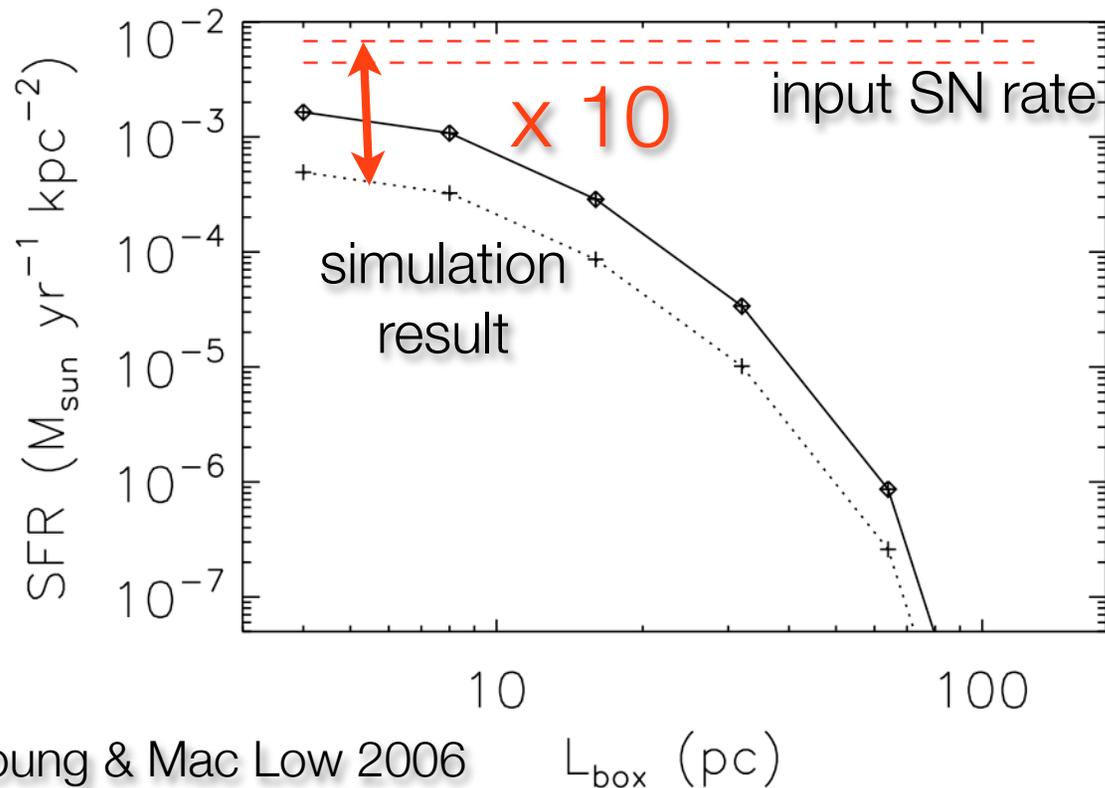
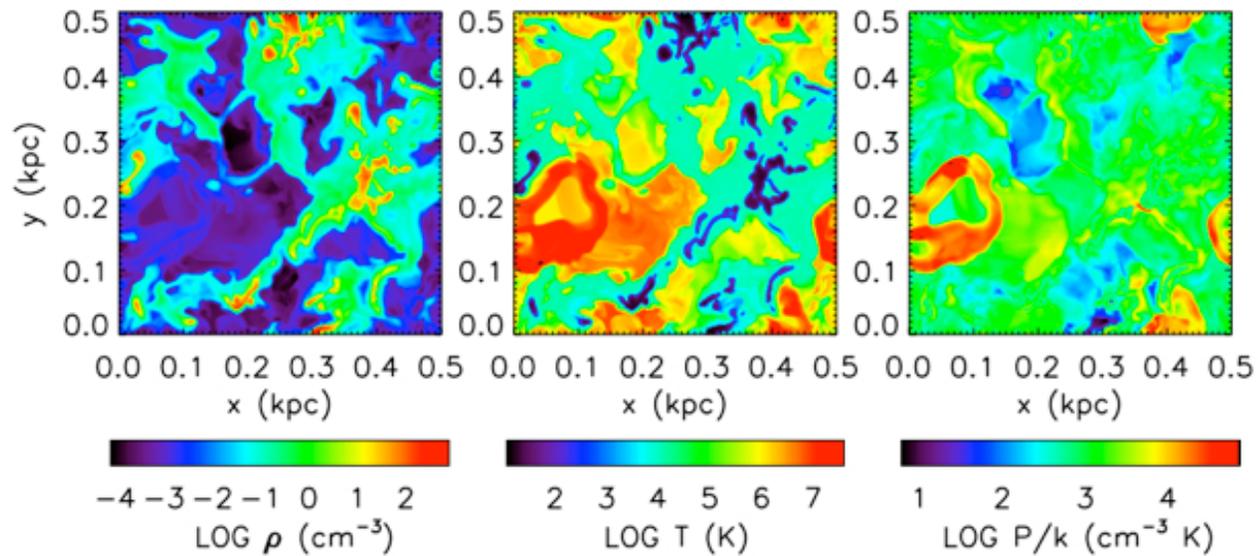
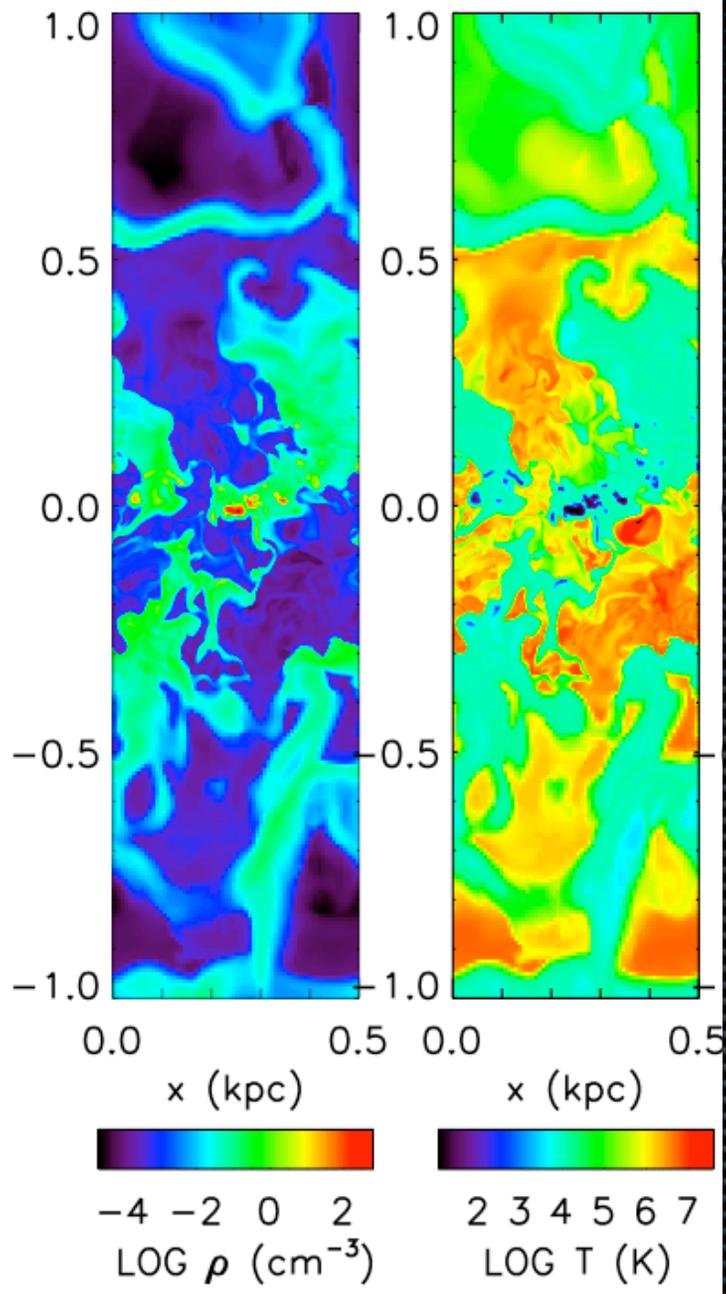
# Turbulence *Inhibits* Collapse

$$M_J = \left( \frac{\pi}{G} \right)^{3/2} \rho_s^{-1/2} c_{s,\text{eff}}^3 \propto$$

$$\propto \frac{c_s}{v} \left( c_s^2 + \frac{\langle v^2 \rangle}{3} \right)^{3/2} \sim v^2$$

- Turbulence is intermittent, so uniform pressure does not represent it well.
- On average, increasing velocity increases Jeans mass, but locally, compressions can decrease it

z (kpc)



Flash (Fryxell + 00) models of stratified, SN-driven ISM

Joung & Mac Low 2006

$L_{\text{box}}$  (pc)

40x40pc

Control run

Induced run

Time = 0.54t<sub>1</sub>



-4

-3

-2

-1

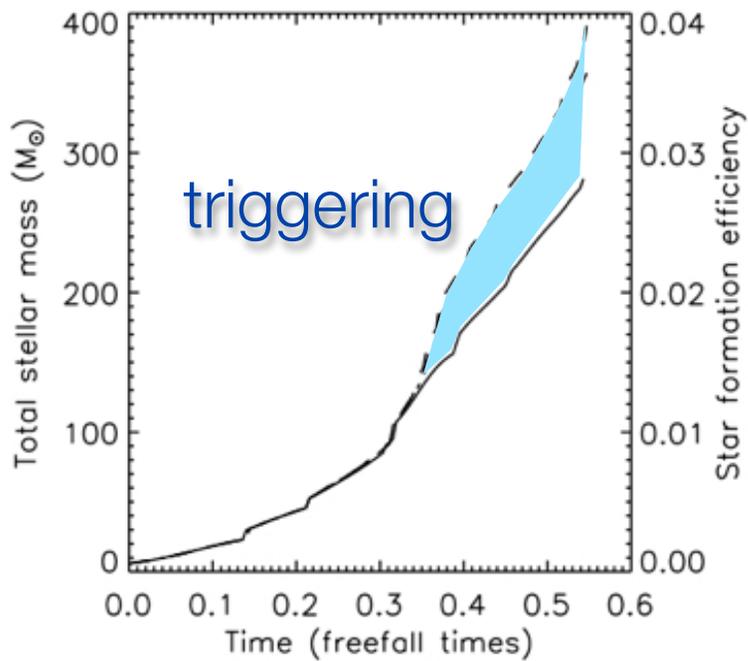
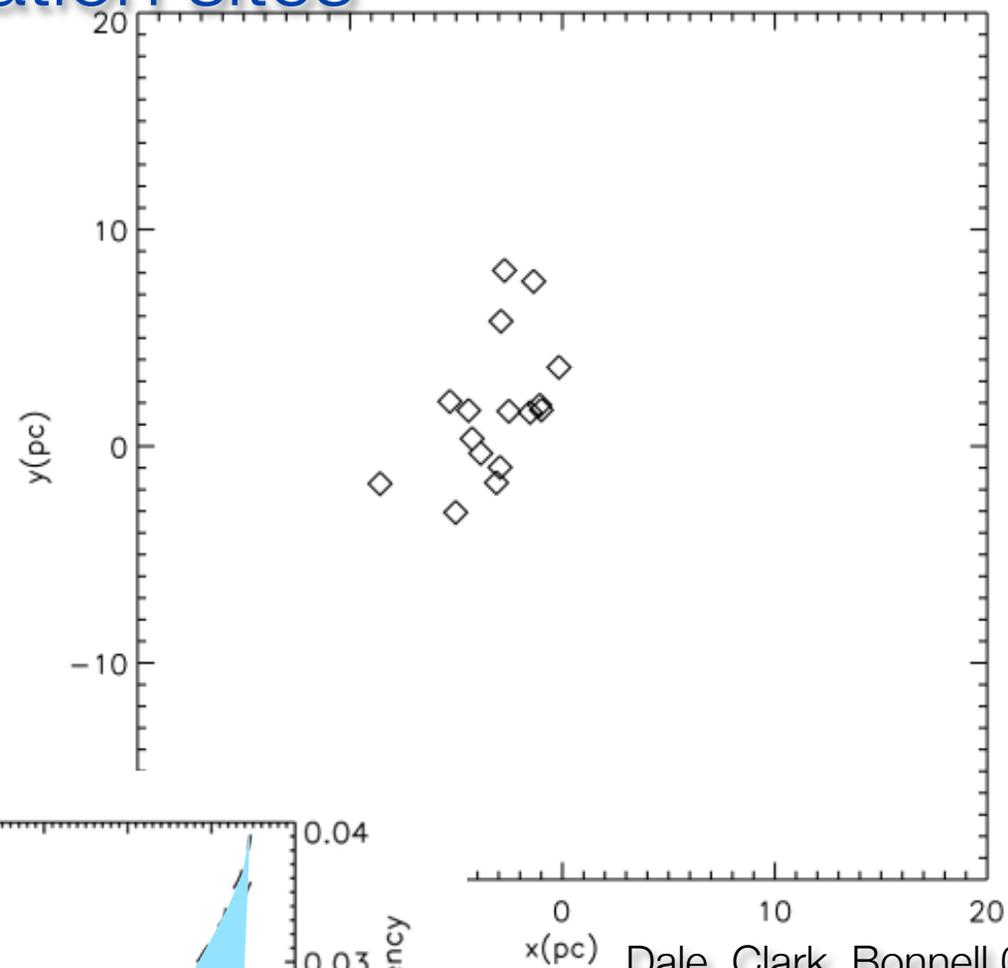
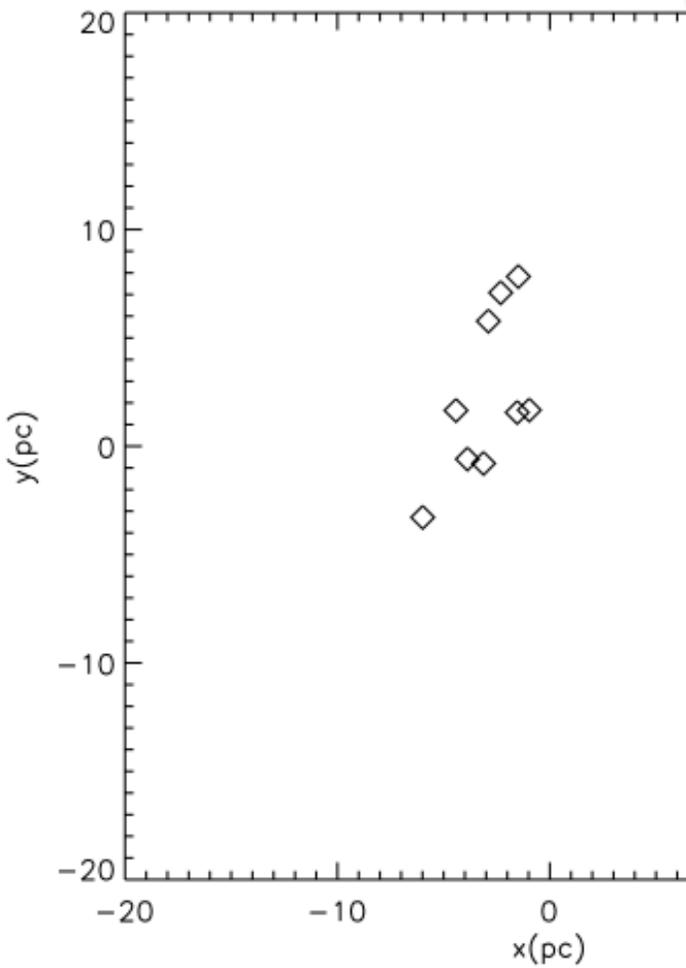
Log Column Density [g/cm<sup>2</sup>]

Dale, Clark, Bonnell 07

# Star formation sites

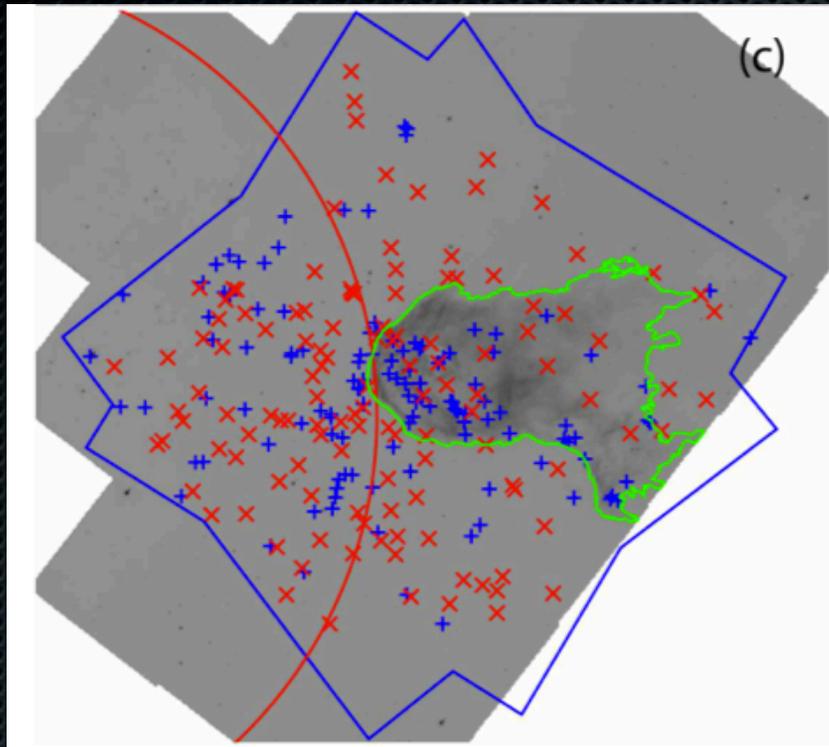
Control run

Feedback run

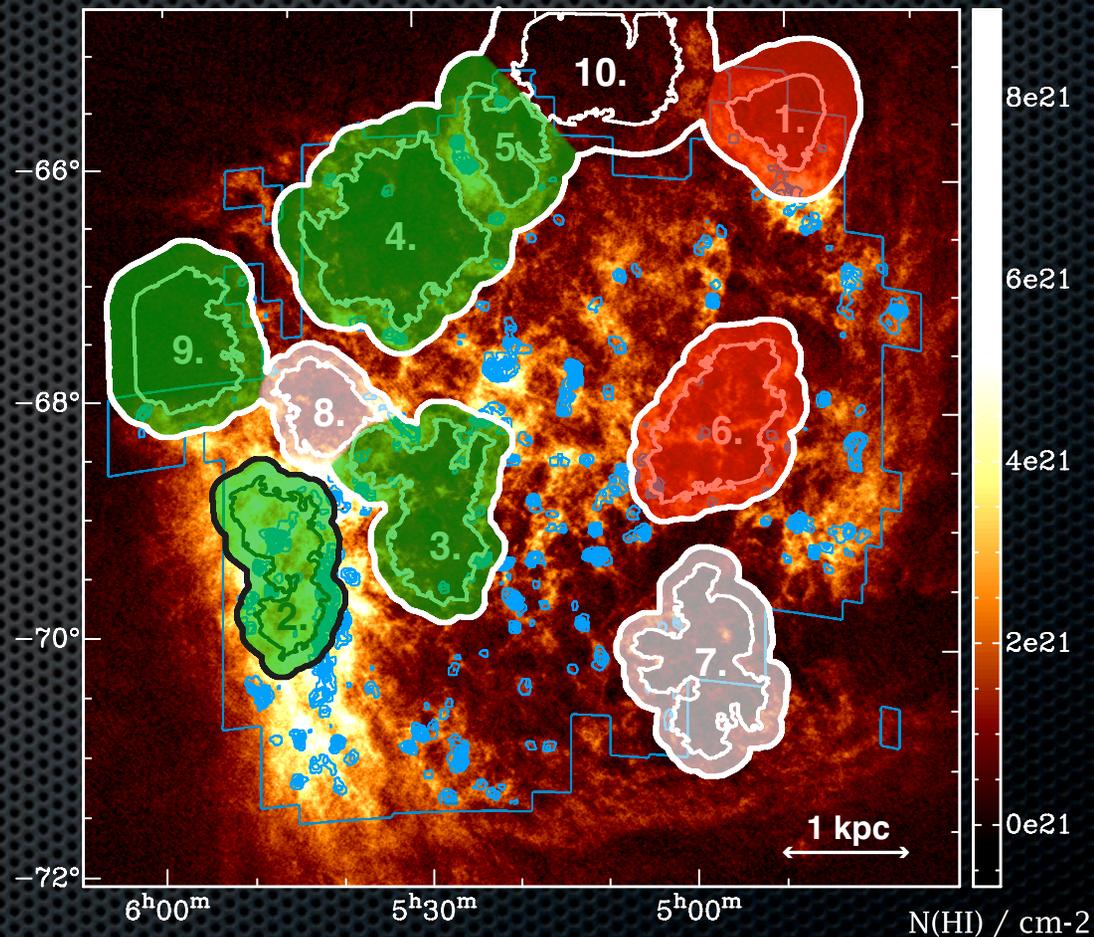


Dale, Clark, Bonnell 07

# Quantitative observations of triggering show that most star formation is untriggered

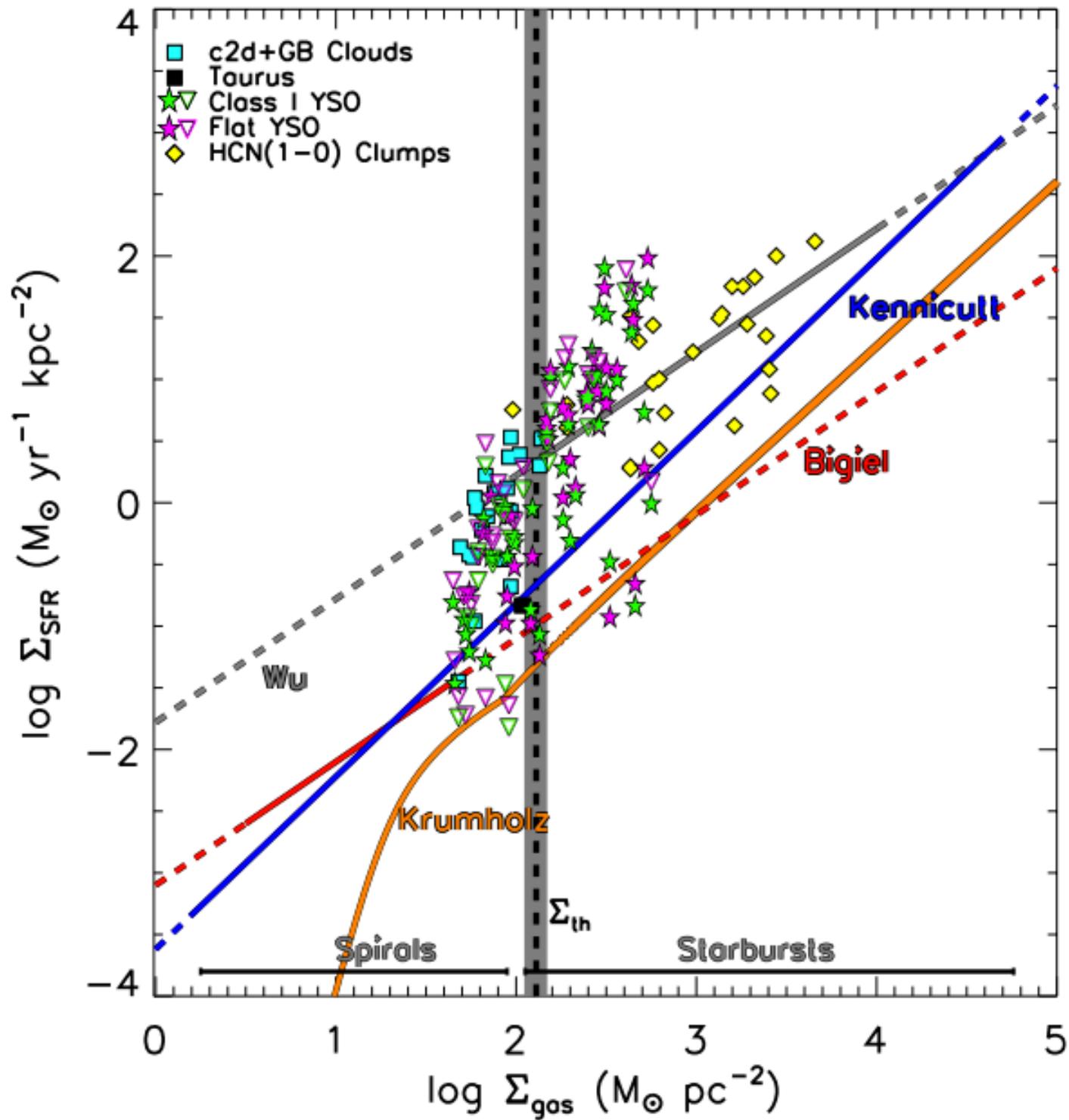


only 14 - 25% of stars in  
Elephant Trunk Nebula triggered  
(Getman + 12)

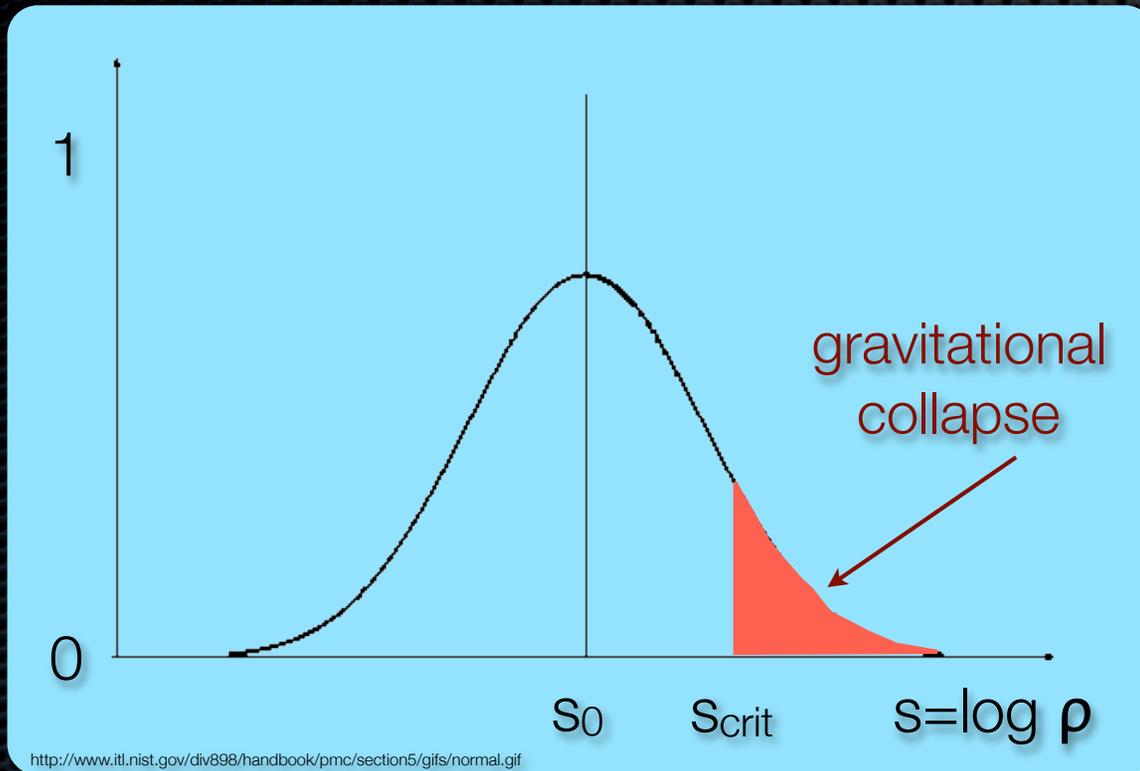


~ 5% of CO in whole  
LMC formed in shell  
walls (Dawson + 12)

# Star Formation Rate



The log-normal density distribution of isothermal turbulence argued to predict the star formation rate.



Krumholz & McKee 05  
Hennebelle & Chabrier 11  
Padoan & Nordlund 11  
Padoan + 12  
Federrath & Klessen 12

derived from similar arguments for the IMF starting with Padoan 95 and Padoan & Nordlund 02.

Models differ on how to choose the relevant density and free-fall time:

$$\text{SFR} \equiv \frac{M_c}{t_{\text{ff}}(\rho_0)} \text{SFR}_{\text{ff}}$$

$$\text{SFR}_{\text{ff}} = \frac{\epsilon}{\phi_t} \int_{s_{\text{crit}}}^{\infty} \frac{t_{\text{ff}}(\rho_0)}{t_{\text{ff}}(\rho)} \frac{\rho}{\rho_0} p(s) ds$$

Federrath & Klessen 12

Analytic Model	Freefall-time Factor	Critical Density $\rho_{\text{crit}}/\rho_0 = \exp(s_{\text{crit}})$
KM	1	$(\pi^2/5) \phi_x^2 \times \alpha_{\text{vir}} \mathcal{M}^2 (1 + \beta^{-1})^{-1}$
PN	$t_{\text{ff}}(\rho_0)/t_{\text{ff}}(\rho_{\text{crit}})$	$(0.067) \theta^{-2} \times \alpha_{\text{vir}} \mathcal{M}^2 f(\beta)$
HC	$t_{\text{ff}}(\rho_0)/t_{\text{ff}}(\rho)$	$(\pi^2/5) y_{\text{cut}}^{-2} \times \alpha_{\text{vir}} \mathcal{M}^{-2} (1 + \beta^{-1}) + \tilde{\rho}_{\text{crit,turb}}$
multi-ff KM	$t_{\text{ff}}(\rho_0)/t_{\text{ff}}(\rho)$	$(\pi^2/5) \phi_x^2 \times \alpha_{\text{vir}} \mathcal{M}^2 (1 + \beta^{-1})^{-1}$
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original

(Krumholz & McKee 05)

magnetic fields

(Padoan & Nordlund 11)

multi-density

(Hennebelle & Chabrier 11)

combined

(Federrath & Klessen 12)

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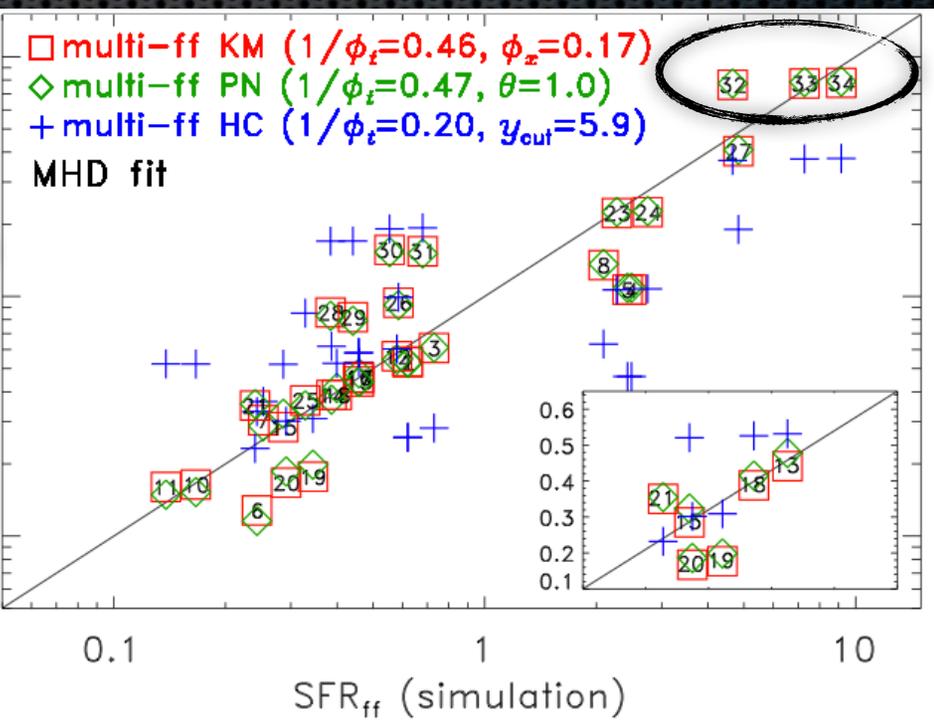
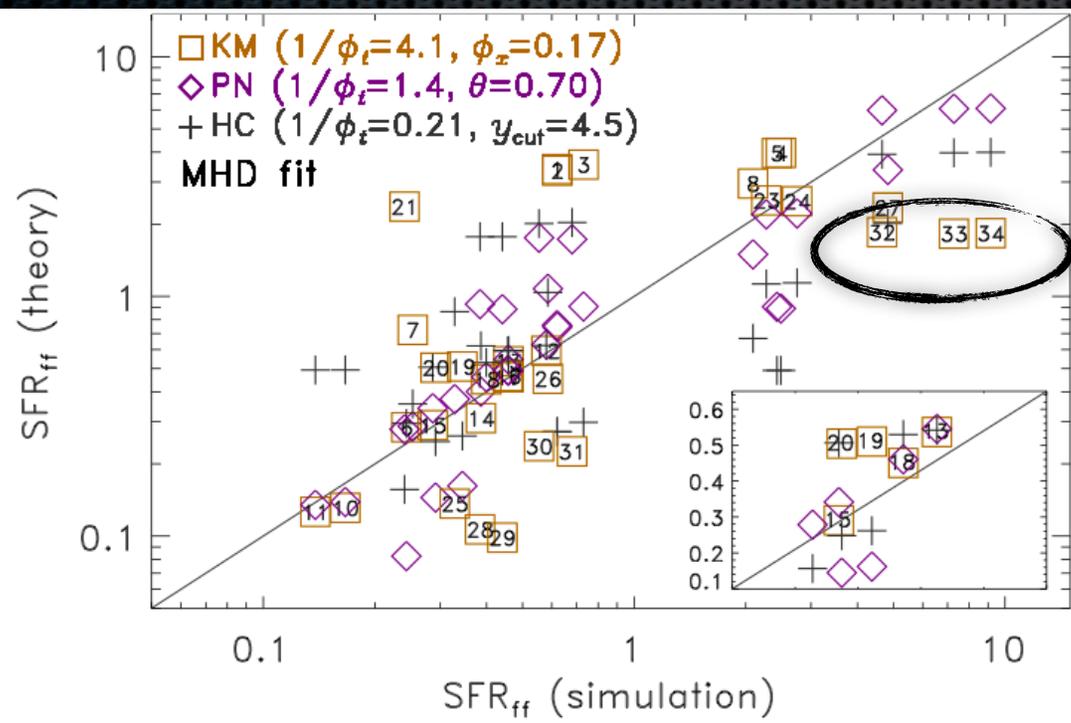
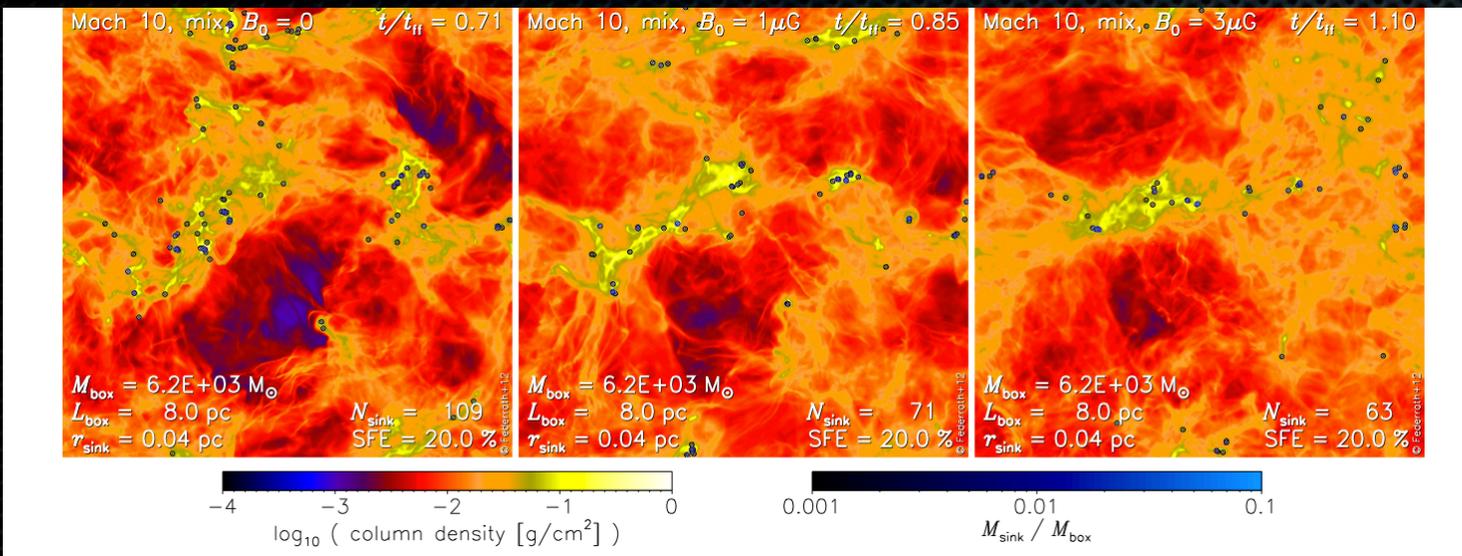
(Padoan & Nordlund 11)

multi-density

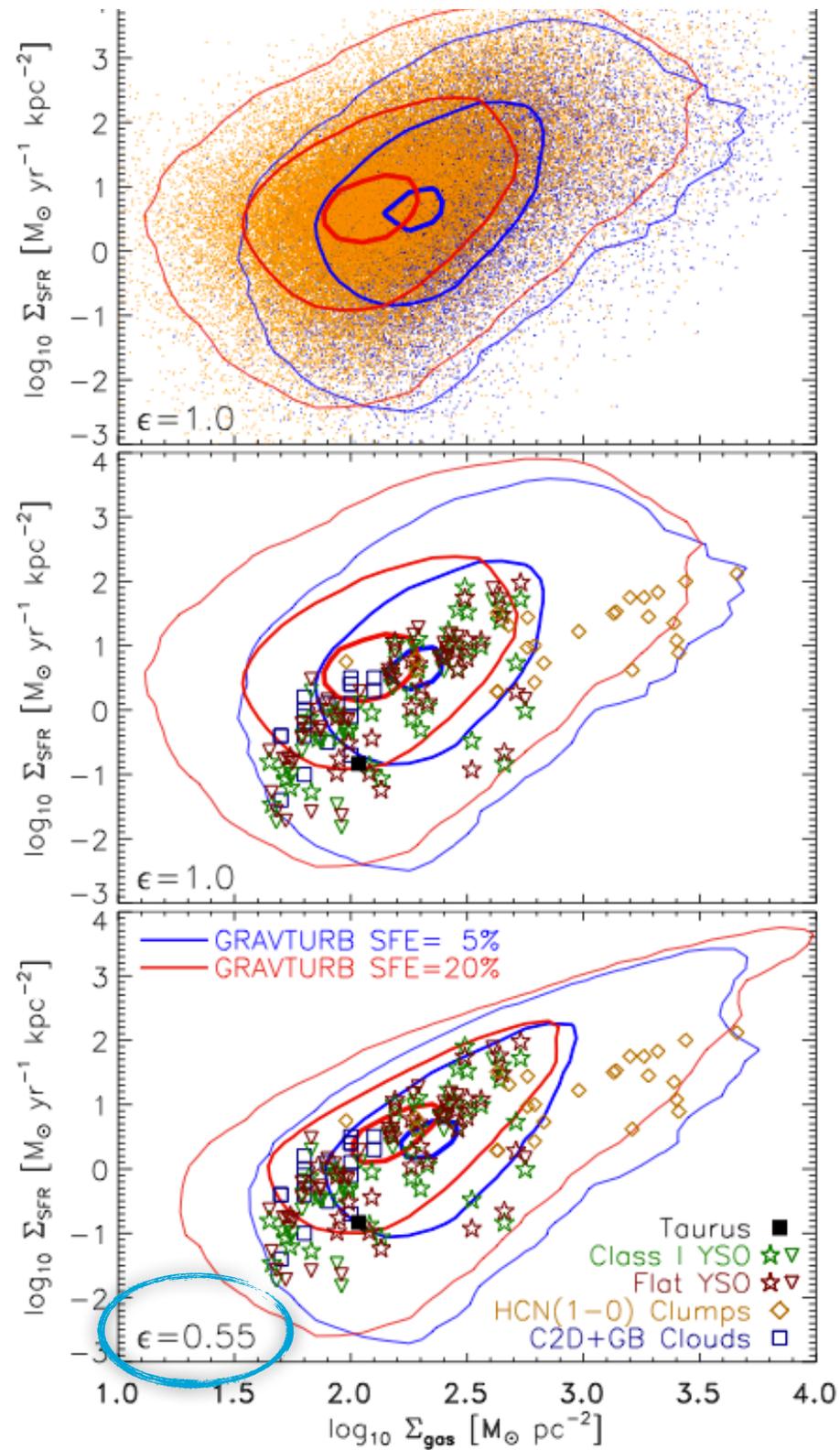
(Hennebelle & Chabrier 11)

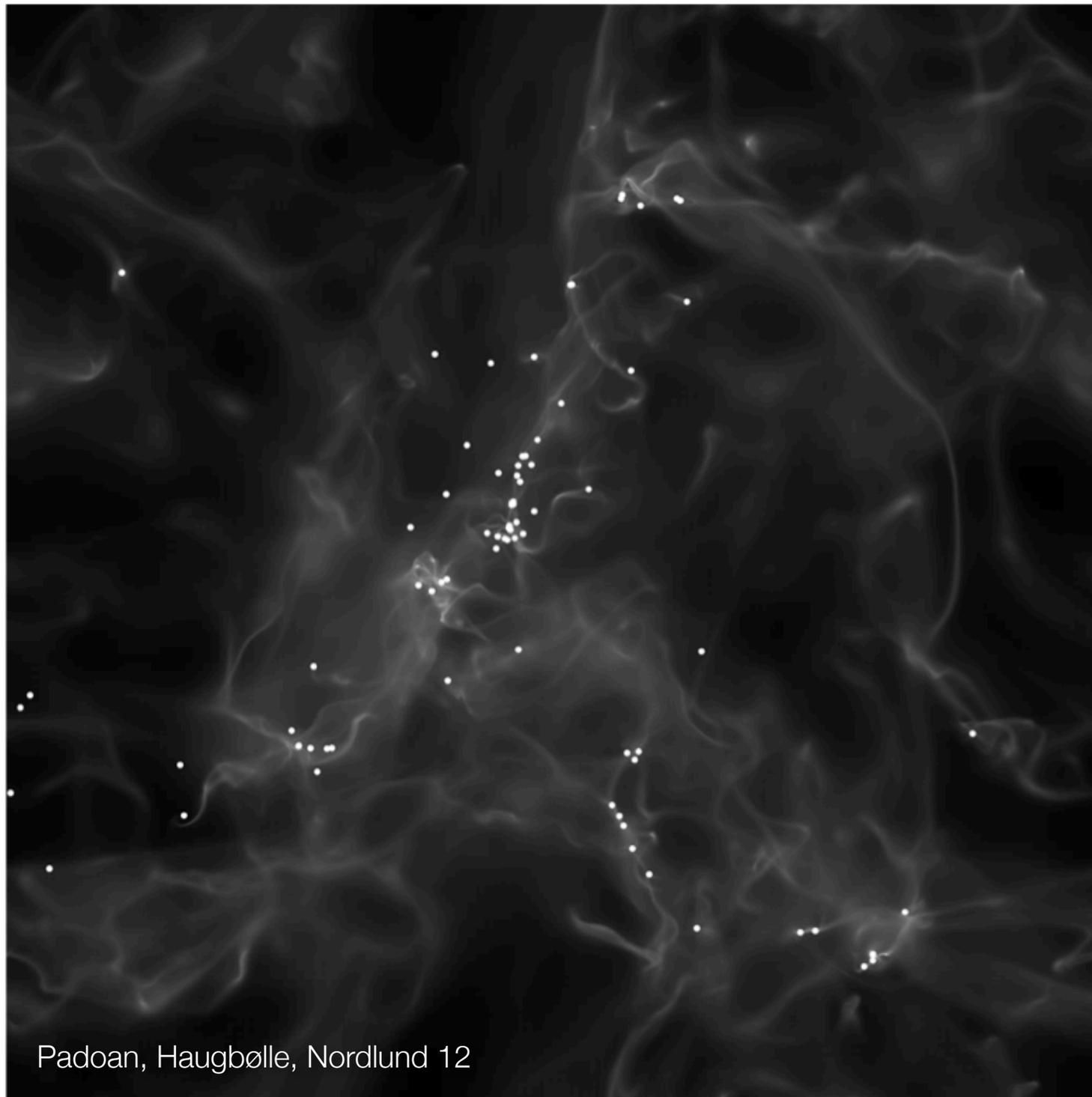
combined

(Federrath & Klessen 12)

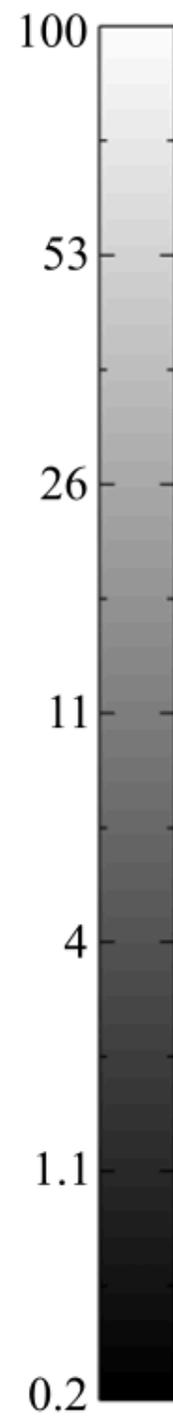


Federrath & Klessen 12  
data from Heiderman + 10

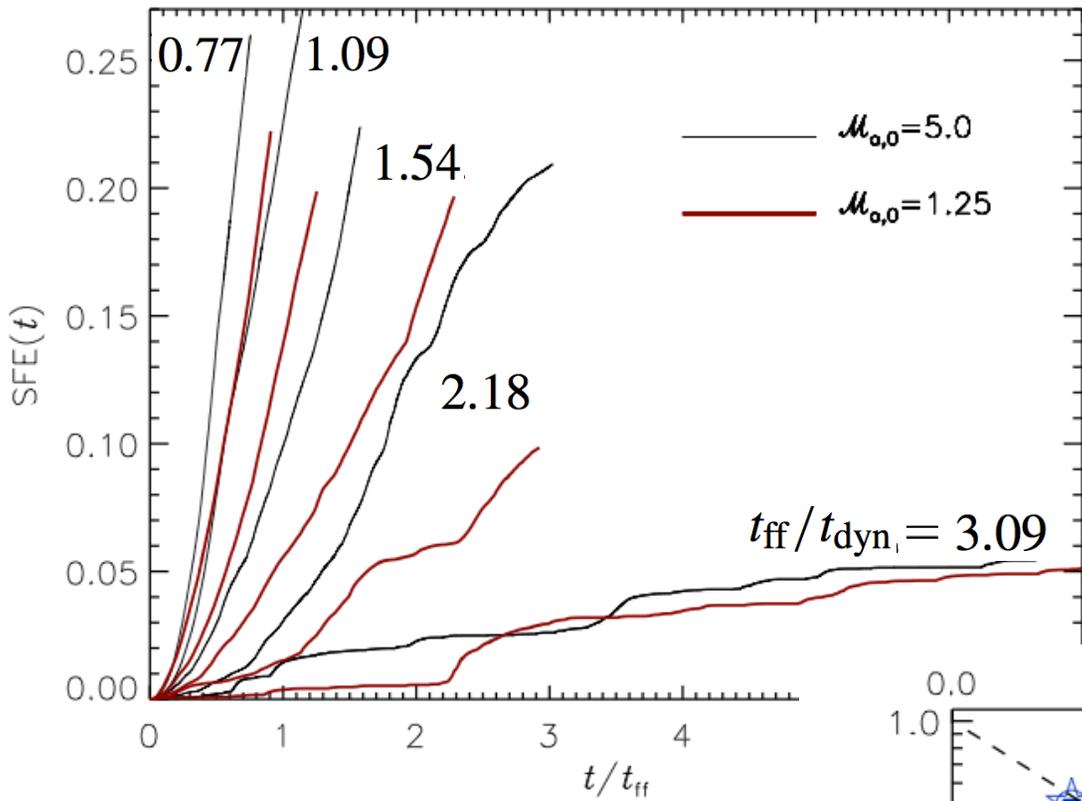




Padoan, Haugbølle, Nordlund 12



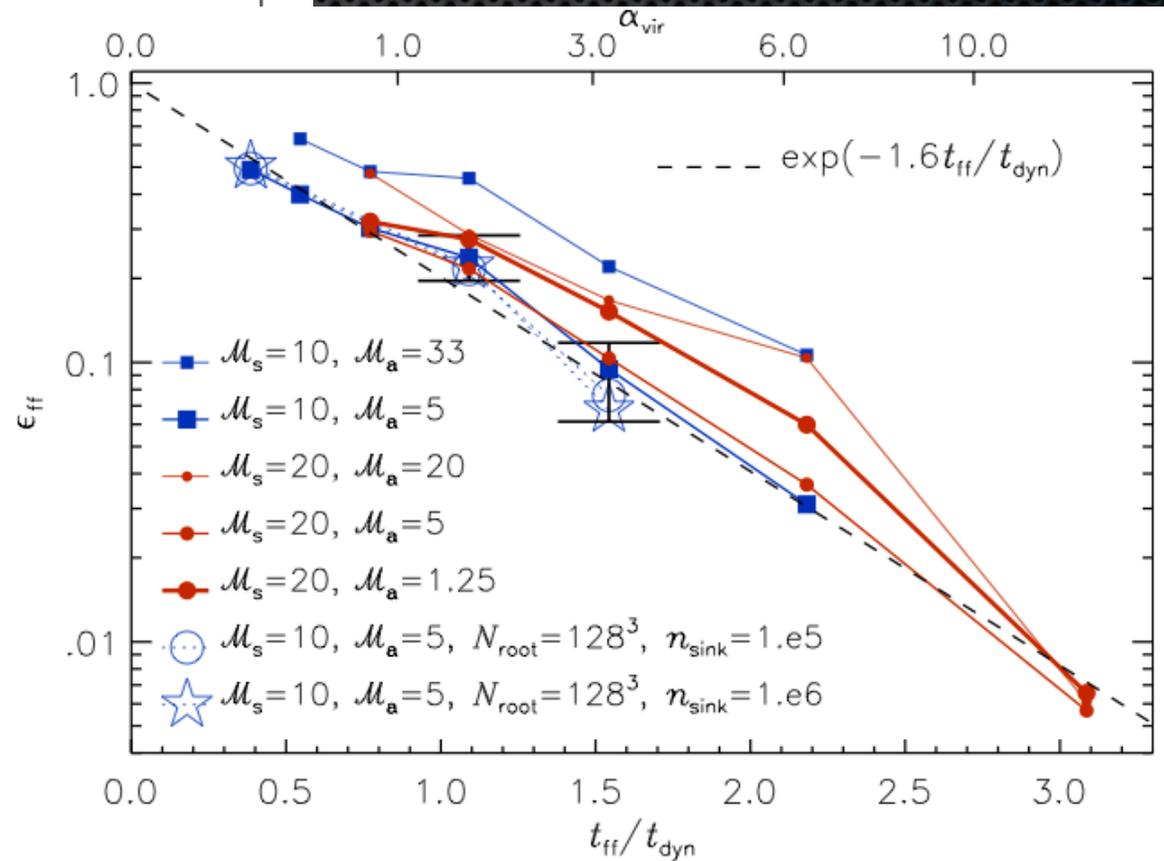
RAMSES AMR  
128<sup>3</sup> base grid  
+ 5 levels of  
refinement  
(4096<sup>3</sup> equiv.)



$$\epsilon_{ff} \approx \epsilon_{wind} \exp(-1.6 t_{ff}/t_{dyn})$$

$$\epsilon_{wind} \approx 0.5$$

Padoan, Haugbølle, Nordlund 12

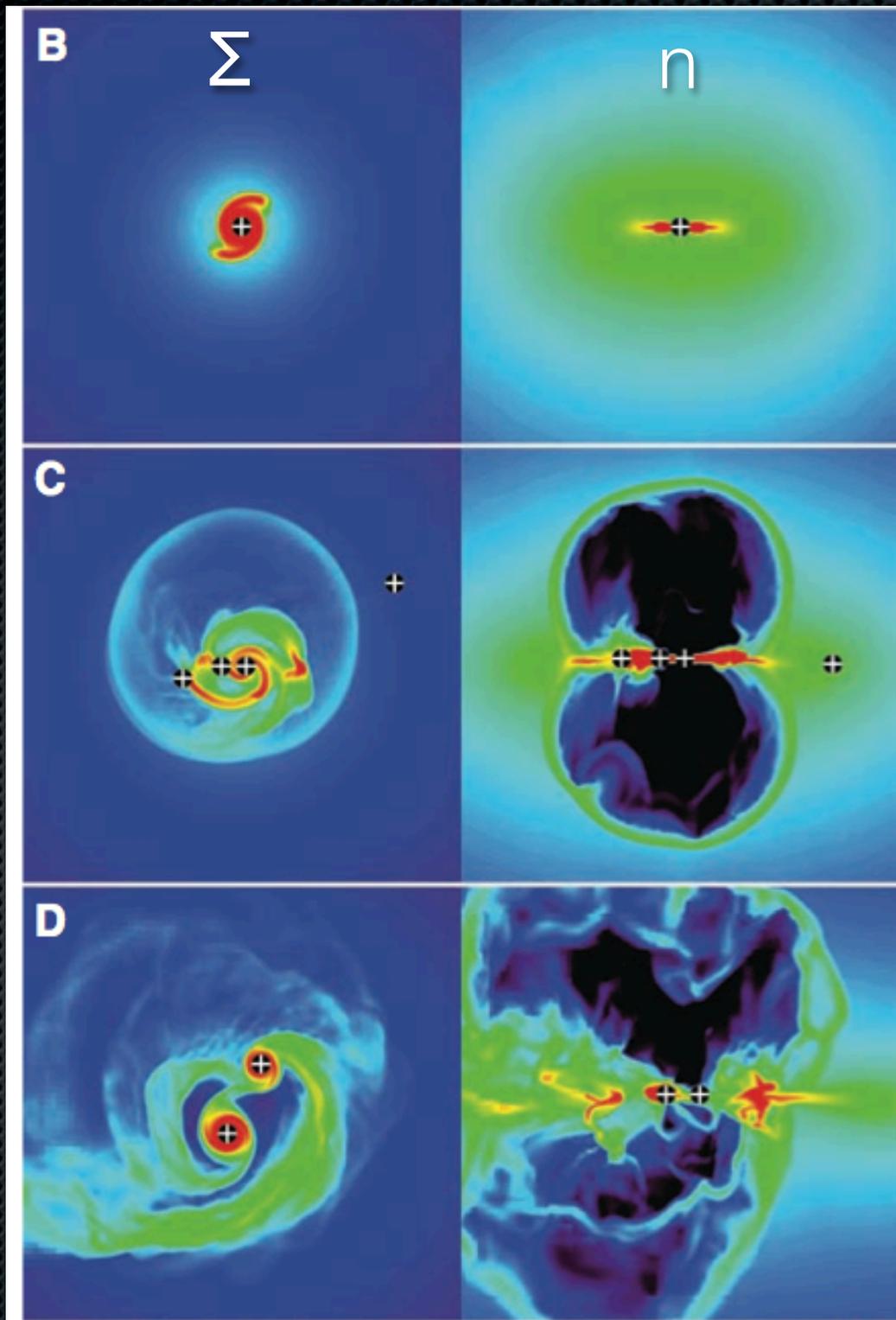


# Massive star formation

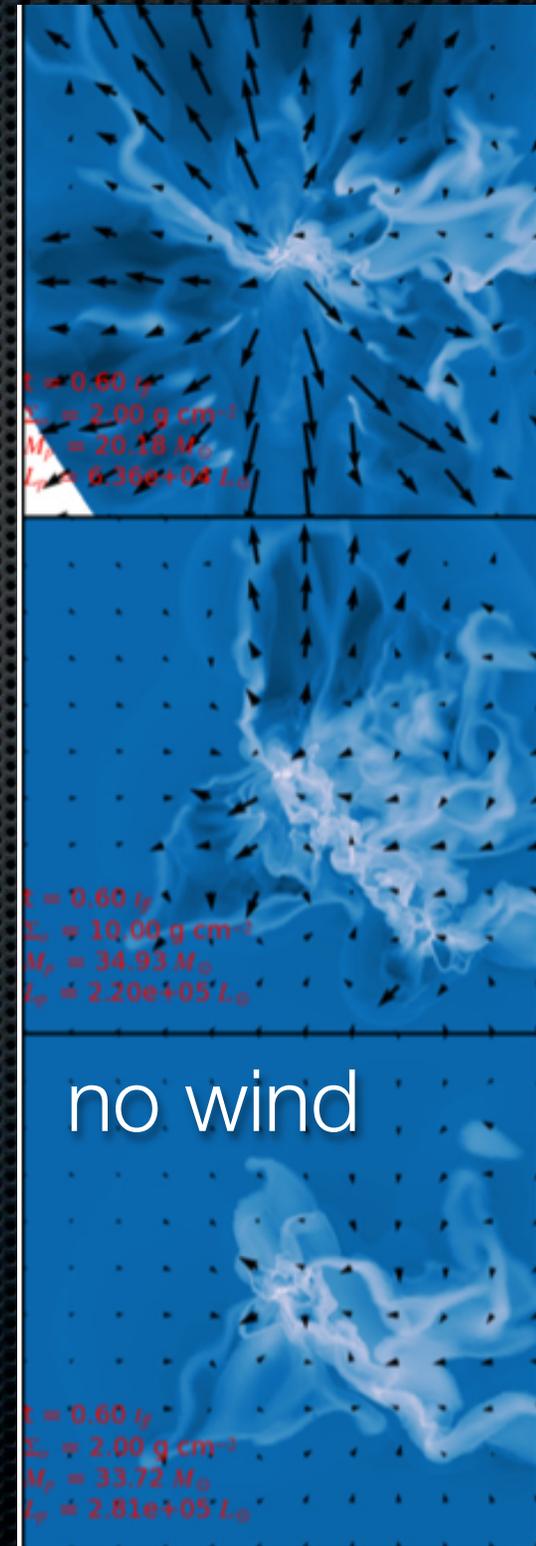
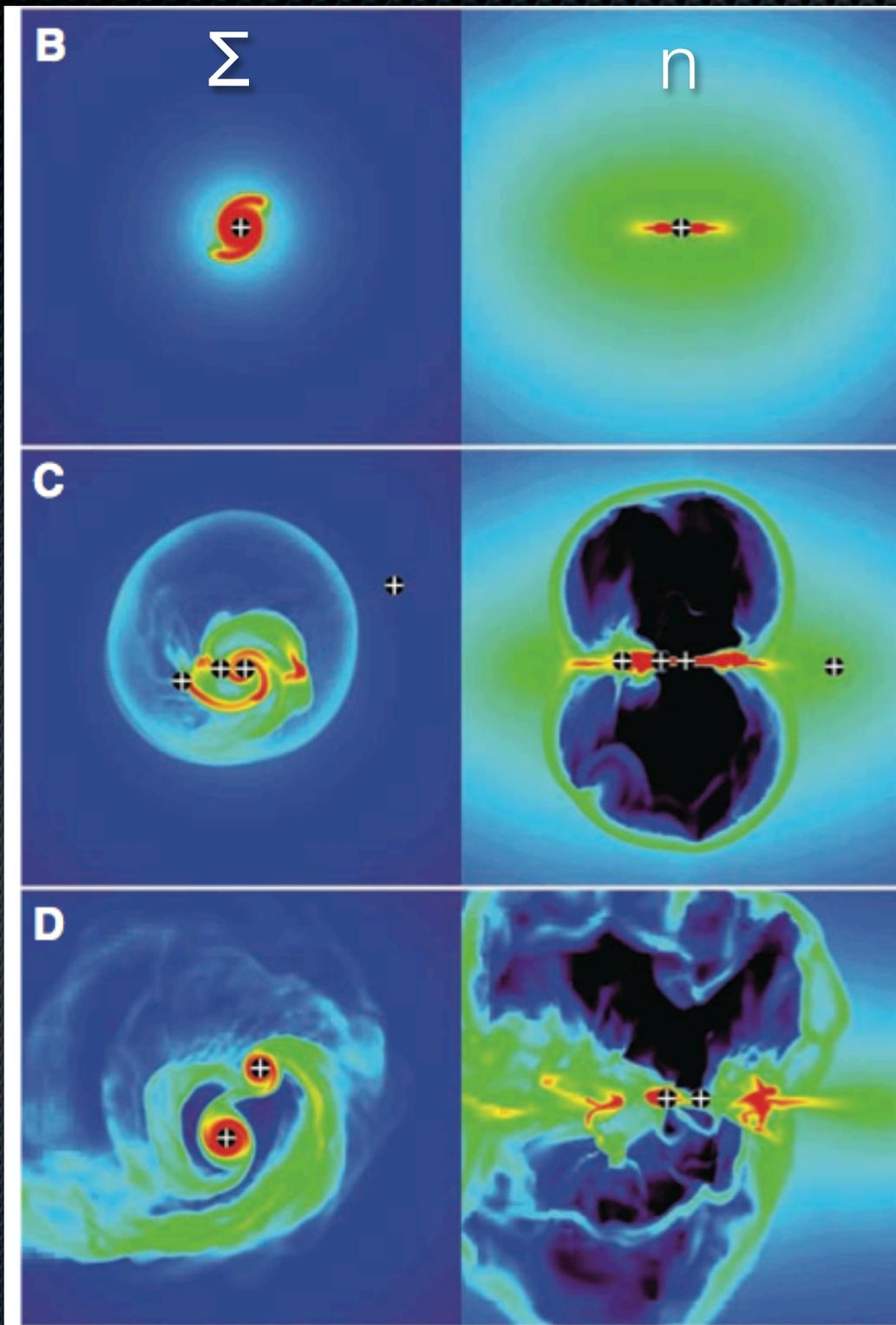


NGC 3603, [NASA](#), [ESA](#), R. O'Connell (University of Virginia), F. Paresce (National Institute for Astrophysics, Bologna, Italy), E. Young (Universities Space Research Association/Ames Research Center), the WFC3 Science Oversight Committee, and the [Hubble Heritage Team \(STScI/AURA\)](#)

Radiation pressure can't prevent massive star formation (Krumholz + 09)...



Radiation pressure  
can't  
prevent  
massive  
star  
formation  
(Krumholz + 09)...



...particularly  
not when  
outflows  
occur  
(Cunningham + 11)

# What about ionization?

$$\left(\frac{\rho}{\rho_0}\right) = \left(\frac{r}{0.5 \text{ pc}}\right)^{-3/2}$$

$$\rho_0 = 3 \times 10^{-20} \text{ g cm}^{-3}$$

$$M = 1000 M_{\odot}$$

$$\beta \equiv E_{rot}/E_{grav} = 0.05$$

solid body rotation

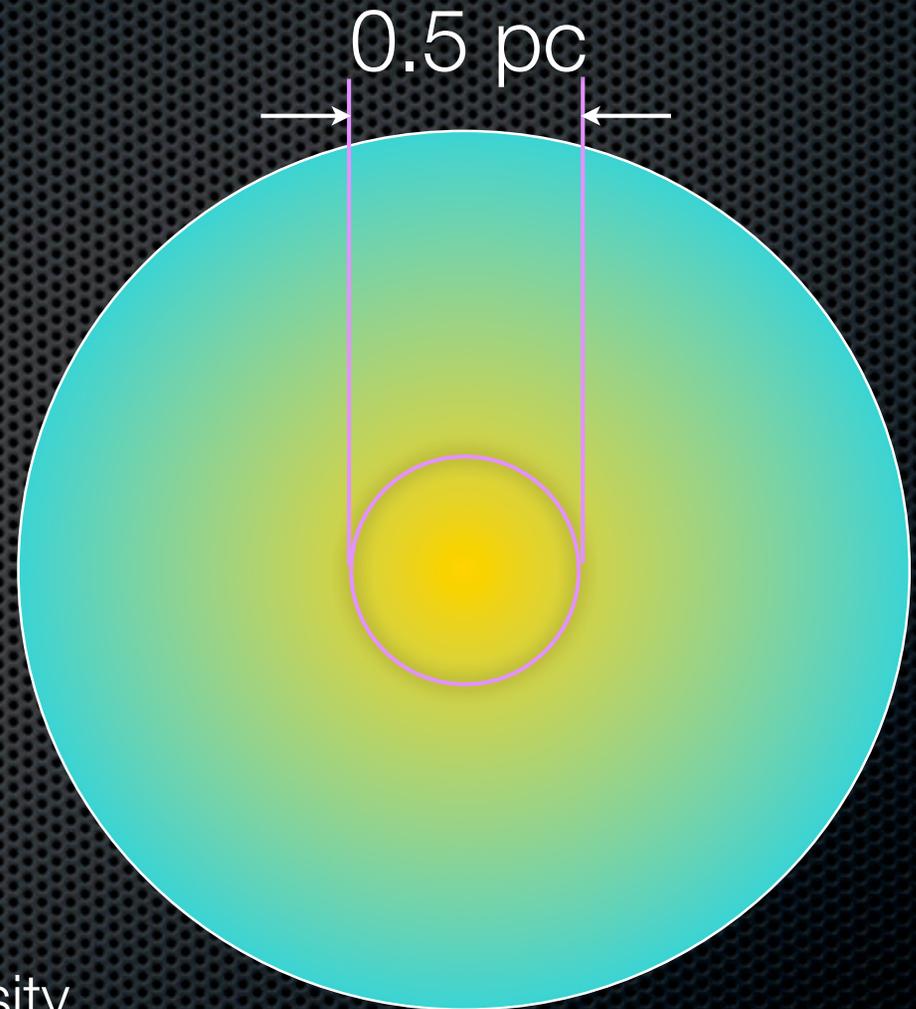
$$\Gamma = \Gamma_{ph} + \Gamma_{st} + \Gamma_{acc}$$

photoionization      accretion luminosity  
dust heating by star

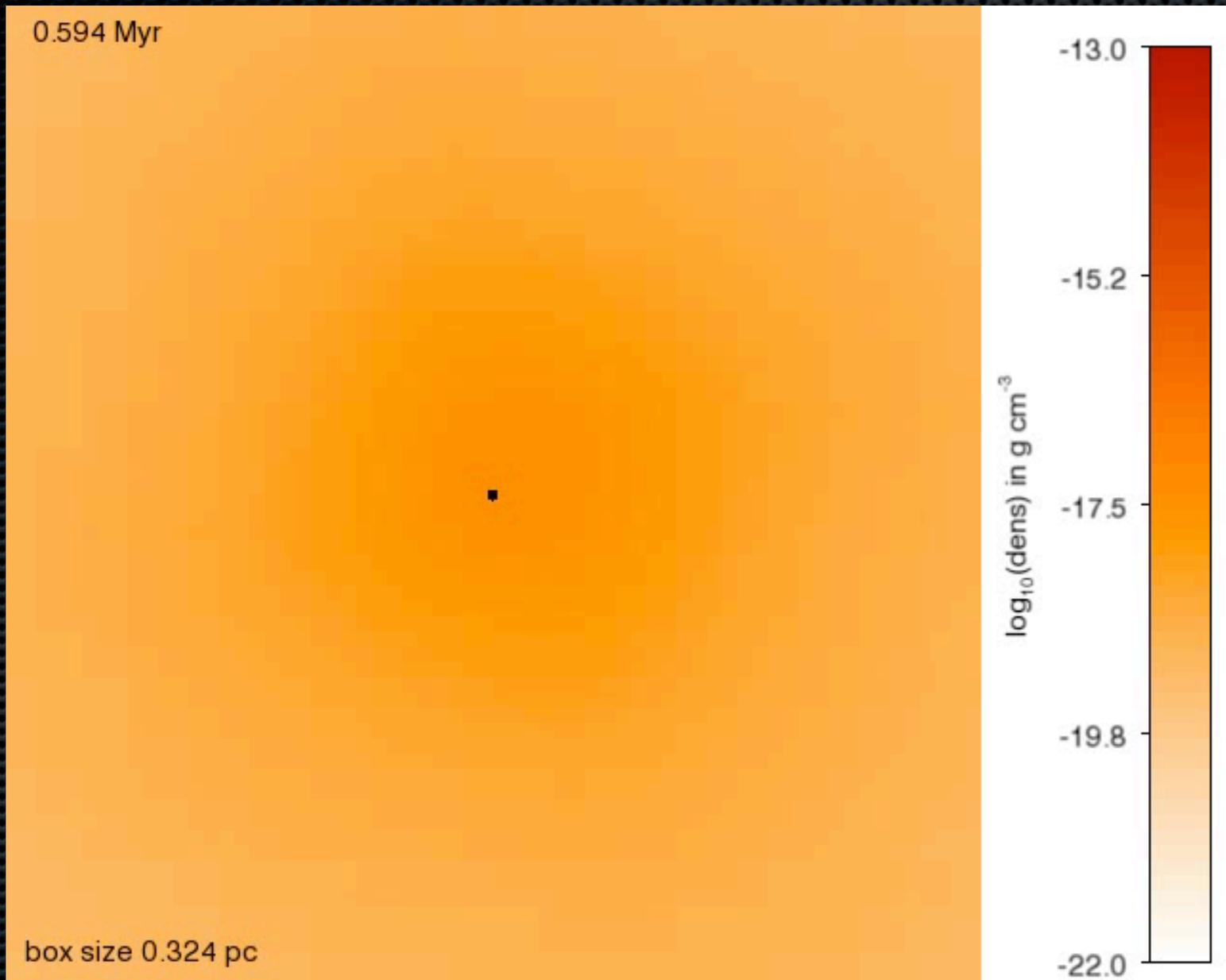
$$\Lambda = \Lambda_{ml} + \Lambda_{mol} + \Lambda_{gd}$$

metal lines      molecules      dust

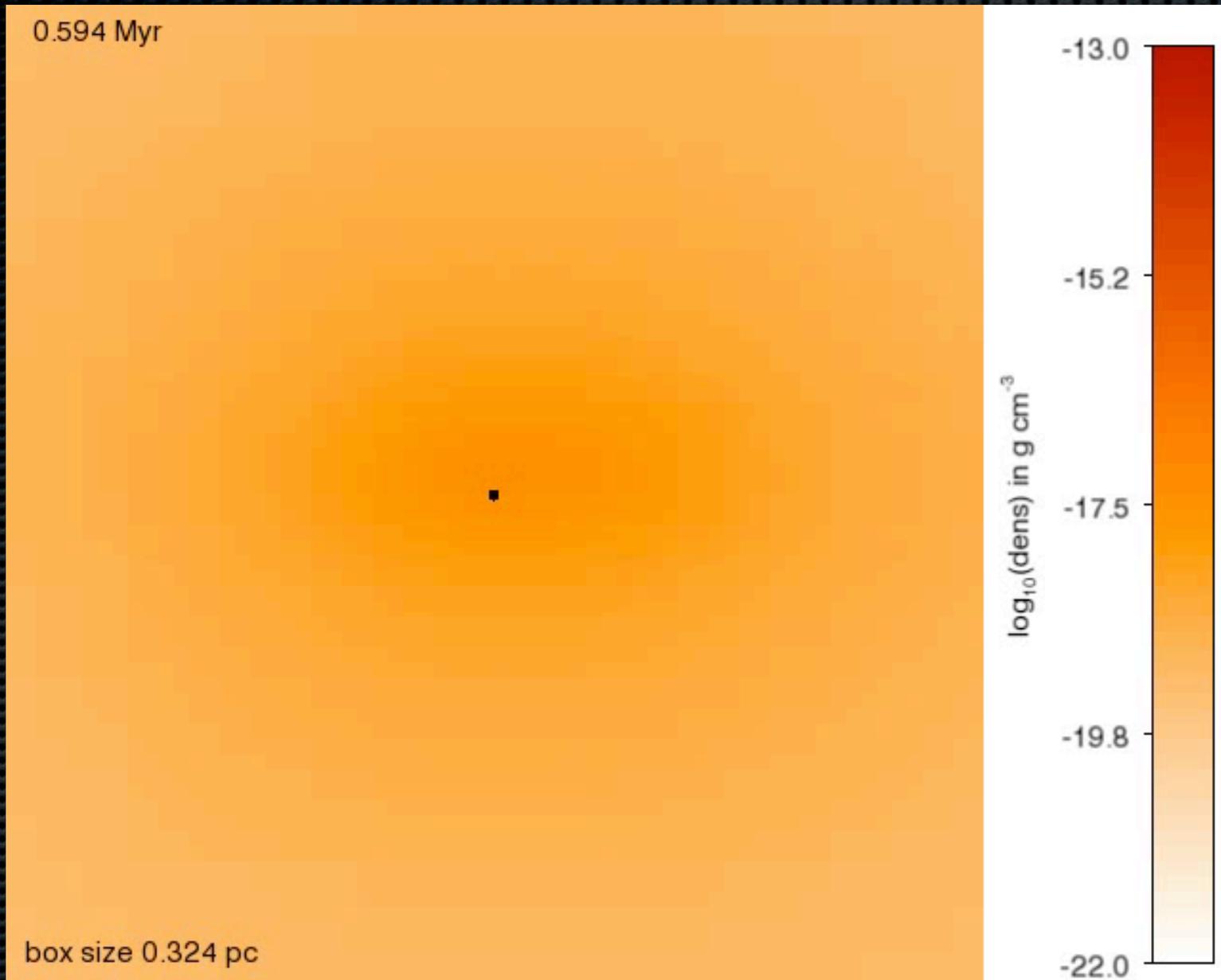
FLASH adaptive mesh refinement (Fryxell 00)



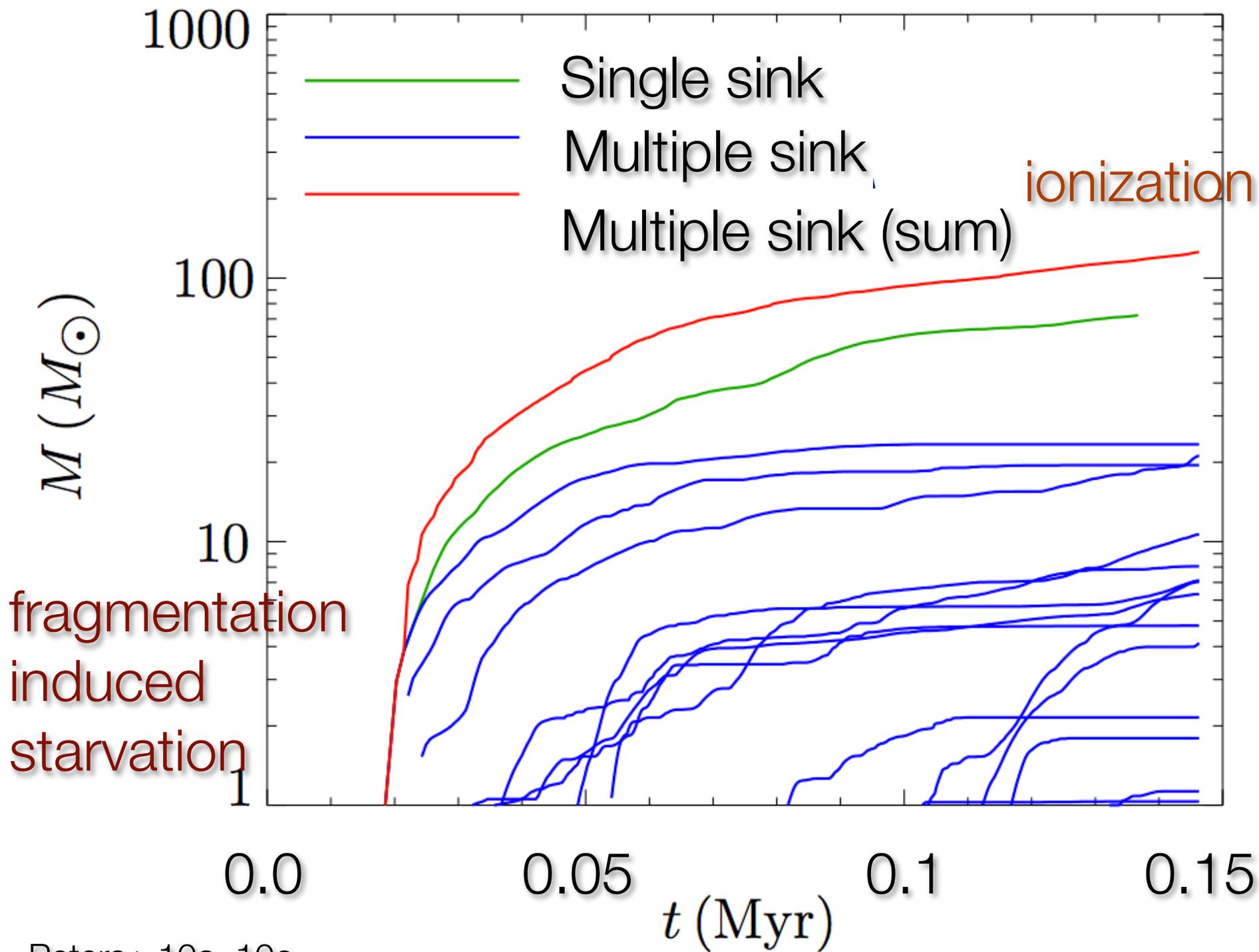
sink particles, ionization  
 along hybrid characteristics  
 (Rijkhorst + 06; modified to allow  
 higher resolution)



Peters+ 10a

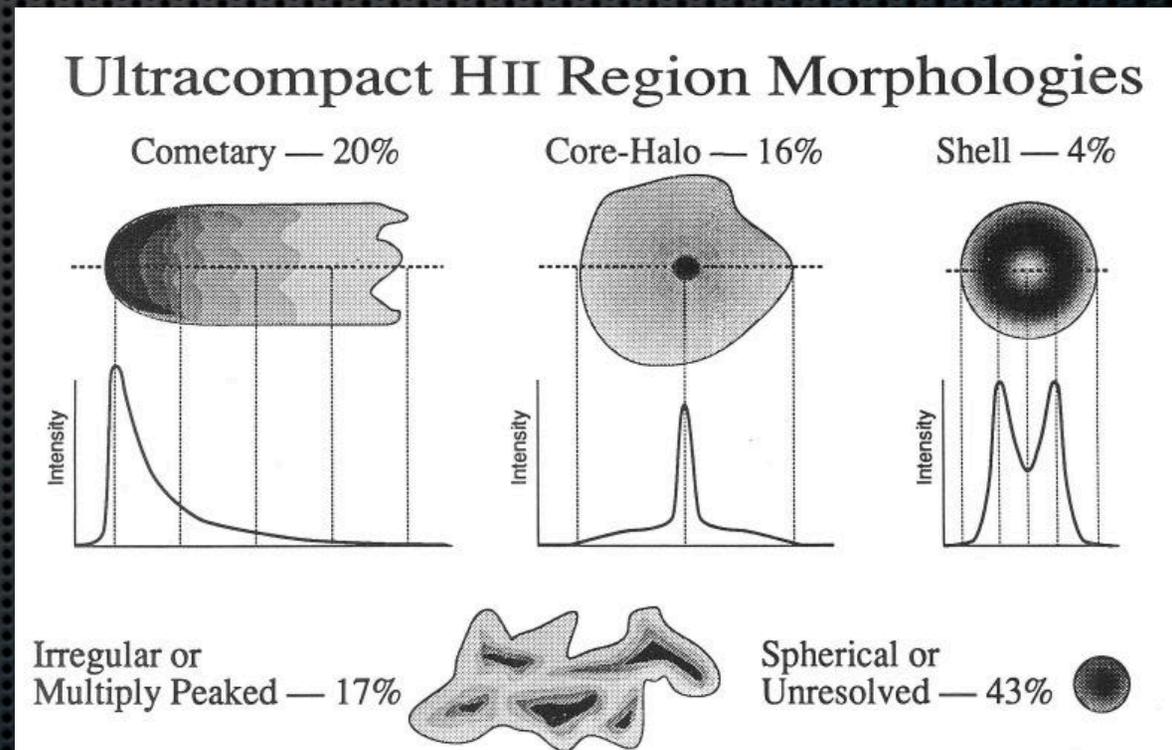


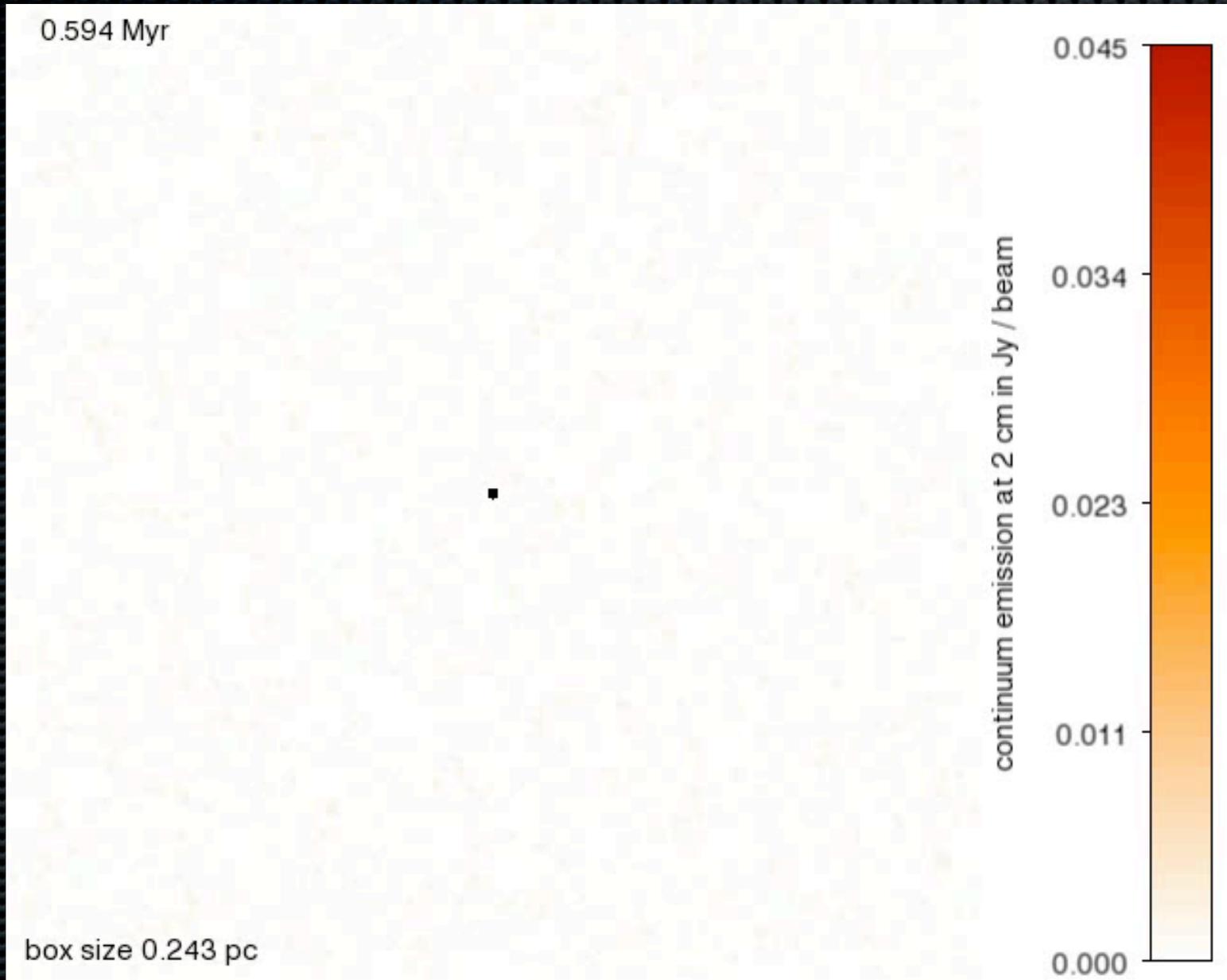
Peters+ 10a



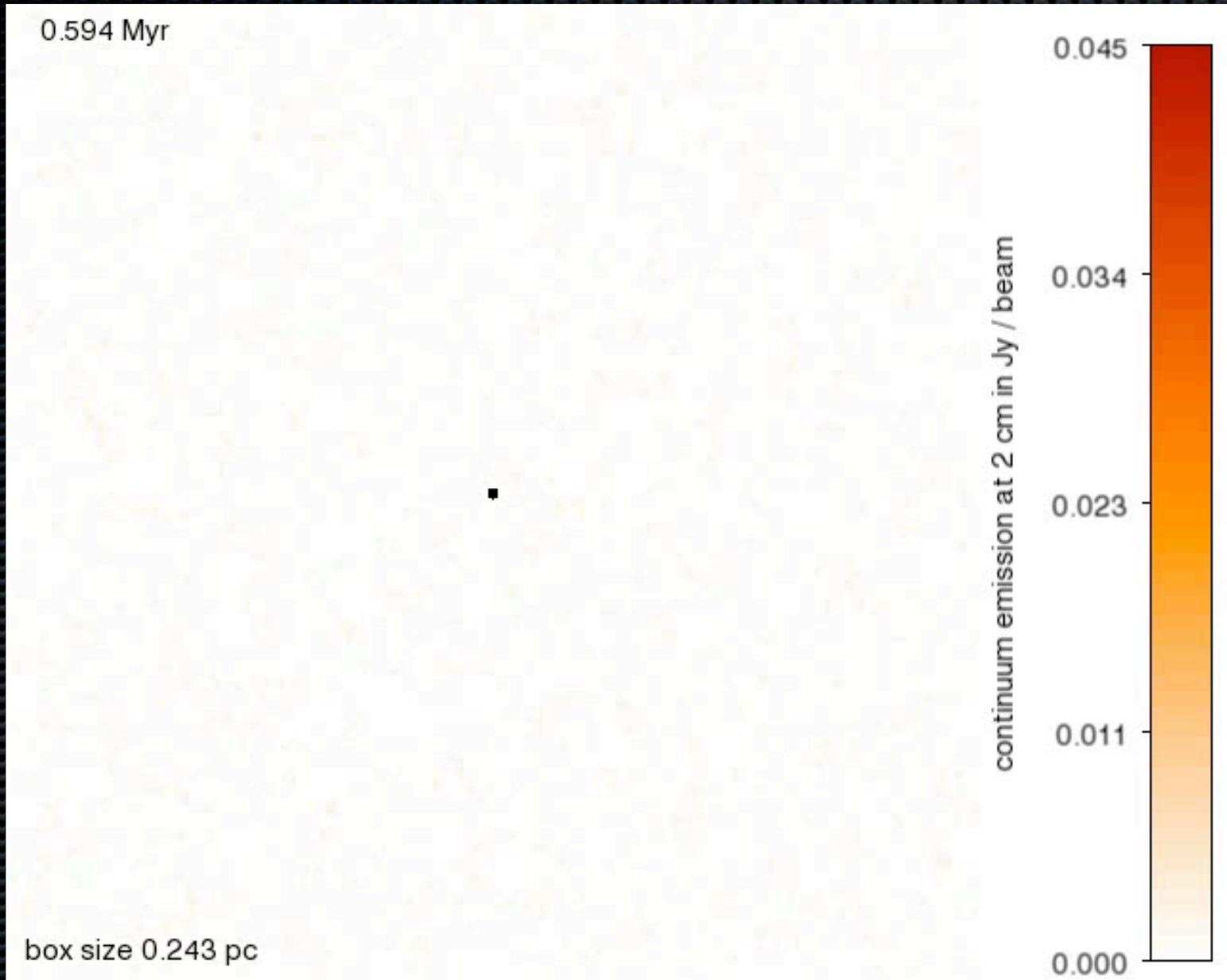
# Ultracompact H II Region live too long and have a variety of shapes.

- Radio continuum sources with  $r < 0.1$  pc,  $EM > 10^7$  pc  $\text{cm}^{-6}$ . Probably most easily observed consequences of ionizing feedback.
- Lifetime problem: if every UC H II region seen surrounds an OB star, UC H II lifetime is  $10^5$  yr, but dynamical ages are only  $10^4$  yr
- Morphologies: What explains wide variety of observed morphologies?





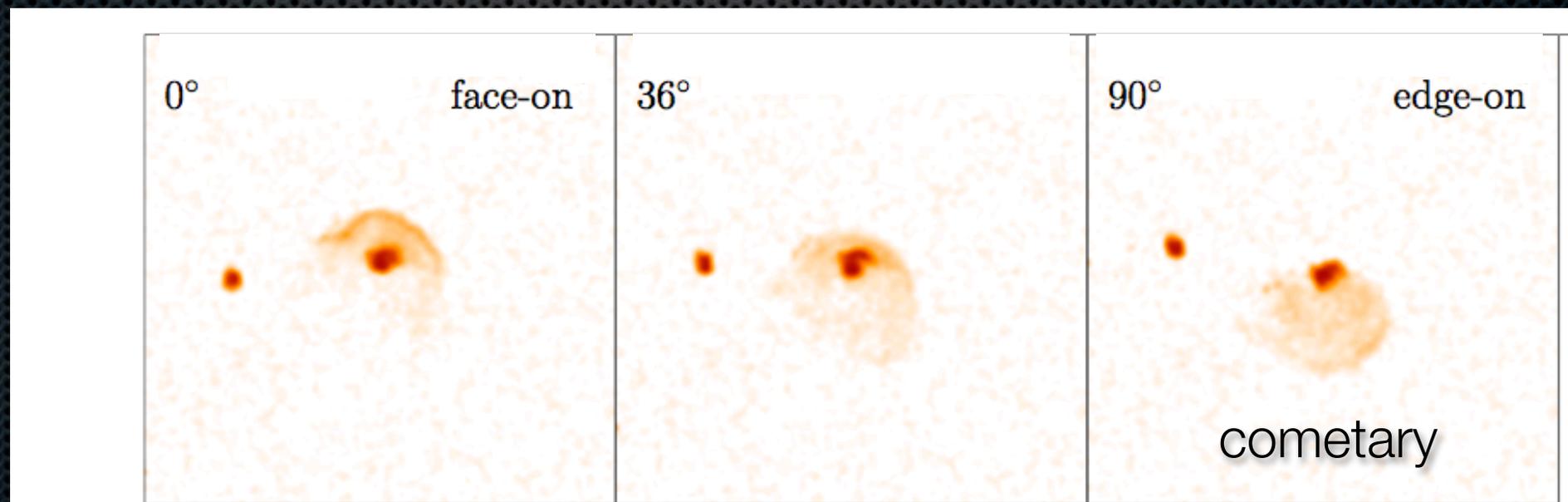
Peters+ 10a, 10b,  
Galván-Madrid 11

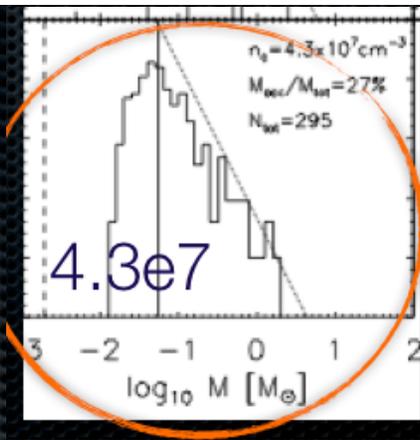
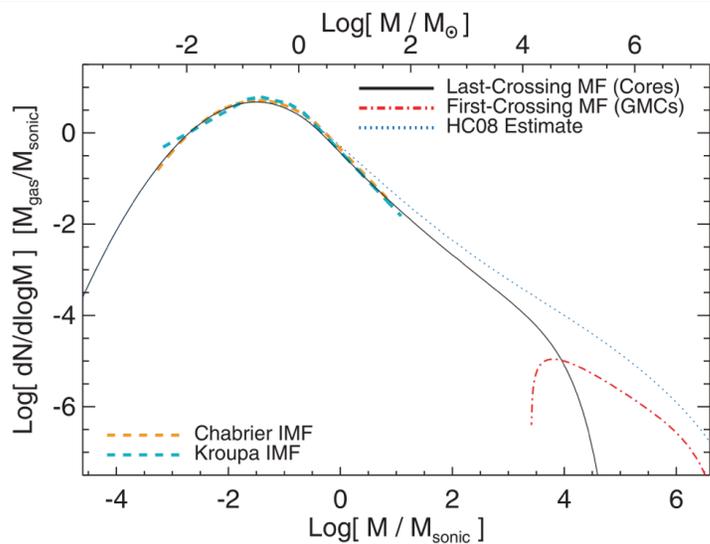


Peters+ 10a, 10b,  
Galván-Madrid 11

**Table 3**  
Percentage Frequency Distribution of Morphologies

Type	WC89	K94	Single sink	Multiple sink
Spherical/Unresolved	43	55	19	60 $\pm$ 5
Cometary	20	16	7	10 $\pm$ 5
Core-halo	16	9	15	4 $\pm$ 2
Shell-like	4	1	3	5 $\pm$ 1
Irregular	17	19	57	21 $\pm$ 5





# Conclusions

